# Entanglement renormalization and integral geometry

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### Emergent geometry in AdS/CFT Our framework

- $AdS_{d+1} \leftrightarrow CFT_d$
- large N, strong coupling  $\Rightarrow$  classical gravity in the bulk
- entanglement = geometry

Ryu-Takayanagi formula



$$S_{\rm EE} = rac{A(\gamma)}{4G}$$

 $\gamma$ : minimal surface ending on  $\partial A$ 

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# Crofton's formula



figures taken from Czech etal. arXiv:1505.05515

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Kinematic space



 $\mathrm{d}s^2 = \frac{1}{z^2} \left( \mathrm{d}z^2 + \mathrm{d}x^2 \right) \qquad \qquad \mathrm{d}s^2 = \frac{1}{\alpha^2} \left( -\mathrm{d}\alpha^2 + \mathrm{d}x^2 \right)$ 

Measure in the kinematic is given by the second derivative of entanglement entropy

$$\epsilon_{\mathcal{K}} \bigl( u, v \bigr) = \frac{\partial^2 S(u,v)}{\partial u \partial v} \mathrm{d} u \wedge \mathrm{d} v$$

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# Point curve and distance



A (blue): (x,z), B (red): (x,z+ $\delta$ ) A (blue): (x,z), B (red): (x- $\delta$ ,z)

The distance is given by integration over the region between the point curves

$$\frac{\mathrm{d}(A,B)}{4G} = \frac{1}{4} \int_{\left(\tilde{p}_A \cup \tilde{p}_B\right) - \left(\tilde{p}_A \cap \tilde{p}_B\right)} \epsilon_{\mathcal{K}}$$

## Kinematic space of geodesics in general dimensions

Crofton's formula in Higher dimensions:

$$\int_{M^q \cap L_r \neq \varnothing} \sigma_{q+r-d} (M^q \cap L_r) \epsilon_{\mathcal{K}} = \frac{O_d \dots O_{d-r} O_{q+r-d}}{O_r \dots O_1 O_0 O_q} \sigma_q(M^q)$$

$$\epsilon_{\mathcal{K}} = \frac{1}{4G} \det \left[ \frac{\partial^2 S(\vec{x}_1, \vec{x}_2)}{\partial \vec{x}_1 \partial \vec{x}_2} \right] \prod_{i=1}^{d-1} \mathrm{d} x_2^i \wedge \mathrm{d} x_1^i$$

- The measure follows from second derivative of *S*
- S is no longer related to entropy even though it can be computed from field theory

### Line element

We work in a space with translation symmetry

$$\mathrm{d}s^2 = -\tilde{f}(z)\mathrm{d}t^2 + \sum_{i=1}^{d-1}g(z)(\mathrm{d}x^i)^2 + g(z)f(z)\mathrm{d}z^2$$

The area of the strip can be reproduced from the flux of geodesic



$$\begin{aligned} \operatorname{vol}_{S_{d-2} \times ds} &= \sqrt{f(z)g(z)} dz \times g(z)^{\frac{d-2}{2}} \delta_{d-2}(\delta) \\ \Rightarrow ds &= \sqrt{f(z)g(z)} dz \end{aligned}$$

We can read off the line element from an infinitesimal strip

# Reproducing RT-formula



- Count all the pairs with only one point in A
- The result is divergent

$$\frac{\sigma(\gamma_D)}{4G} = \frac{a_{d-2}}{\delta^{d-2}} \cdots + \frac{a_{d-2-2k}}{\delta^{d-2-2k}} \cdots + \begin{cases} a_0 \log \delta & , \text{d even} \\ a_0 & , \text{d odd} \end{cases}, \quad k \in \text{integers} \end{cases}$$

$$a_{0} = \frac{L^{d-1}}{4G} \times \begin{cases} \frac{(-1)^{\frac{d}{2}} O_{d-2}(d-3)!!}{(d-2)!!} & , d \text{ even} \\ \pi^{\frac{d}{2}-1} \Gamma\left(1-\frac{d}{2}\right) & , d \text{ odd} \end{cases}$$

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### Entanglement contour

Entanglement contour  $s_A(i)$  (Chen&Vidal 2015)

$$\sum_{i \in A} s_A(i) = S_A$$

The entanglement entropy of a point can be expressed in the form above

$$S_{EE}(\mathrm{d} x; \tilde{z}) \coloneqq \mathrm{d} x \left( \int_{-\infty}^{\tilde{x}} \frac{\partial^2 S(u, v; \tilde{z})}{\partial u \partial v} \Big|_{v = \tilde{x}} \mathrm{d} u + \int_{\tilde{x}}^{\infty} \frac{\partial^2 S(u, v; \tilde{z})}{\partial u \partial v} \Big|_{u = \tilde{x}} \mathrm{d} v \right)$$

and hence  $\partial^2 S(u, v)$  can be identified as entanglement contour

# Multi-scale entanglement renormalization ansatz



LUs (disentangler, isometry) modify the entanglement structure

MERA version of RT-formula

 $S_A \leq \#(\text{LUs cut})$ 

 $\begin{aligned} \mathsf{MERA} &= \mathsf{Discretized} \ \mathsf{AdS} \\ (\mathsf{Swingle} \ 2009) \end{aligned}$ 

figures taken from Swingle arXiv:1209.3304

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# Surface/state correspondence



Pure State  $|\Phi(\Sigma)\rangle$ 



Trivial State  $|\Omega\rangle$ 

- AdS/CFT applies to Φ(Σ), an excited state
- RT remains valid for  $\Phi(\Sigma)$



Mixed State  $\rho(\Sigma)$ 



Mixed State  $\rho(\Sigma)$ 



figures taken from Miyaji and Takayanagi arXiv:1503.03542

### Renormalization of entanglement contour



The short-distance entanglement (red) is removed while the long-distance (blue) entanglement is invariant under RG flow but it is reshuffled to shorter scale.

# Summary

- The punch line is that we have a duality between the kinematic space ↔ real space. We can reconstruct the kinematic space from field theory.
- The number density of geodesics can be understood as a measure of two-point entanglement (entanglement contour).
- Entanglement contour can be defined on arbitrary surface and hence we can study the RG behavior

# Future Problems

- ▶ ER=EPR (eternal AdS BH) and quench
- Co-dimension two surface higher dimensions
- Construction of bulk operator using Radon transform (Lin etal. 2014)
- DS/CFT and dynamics in the kinematic space (de Boer etal. 2015)