

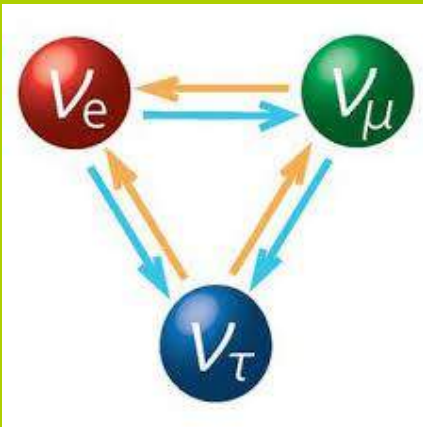
neutrino pathways to new physics

José W F Valle



<https://www.facebook.com/ific.ahep/>

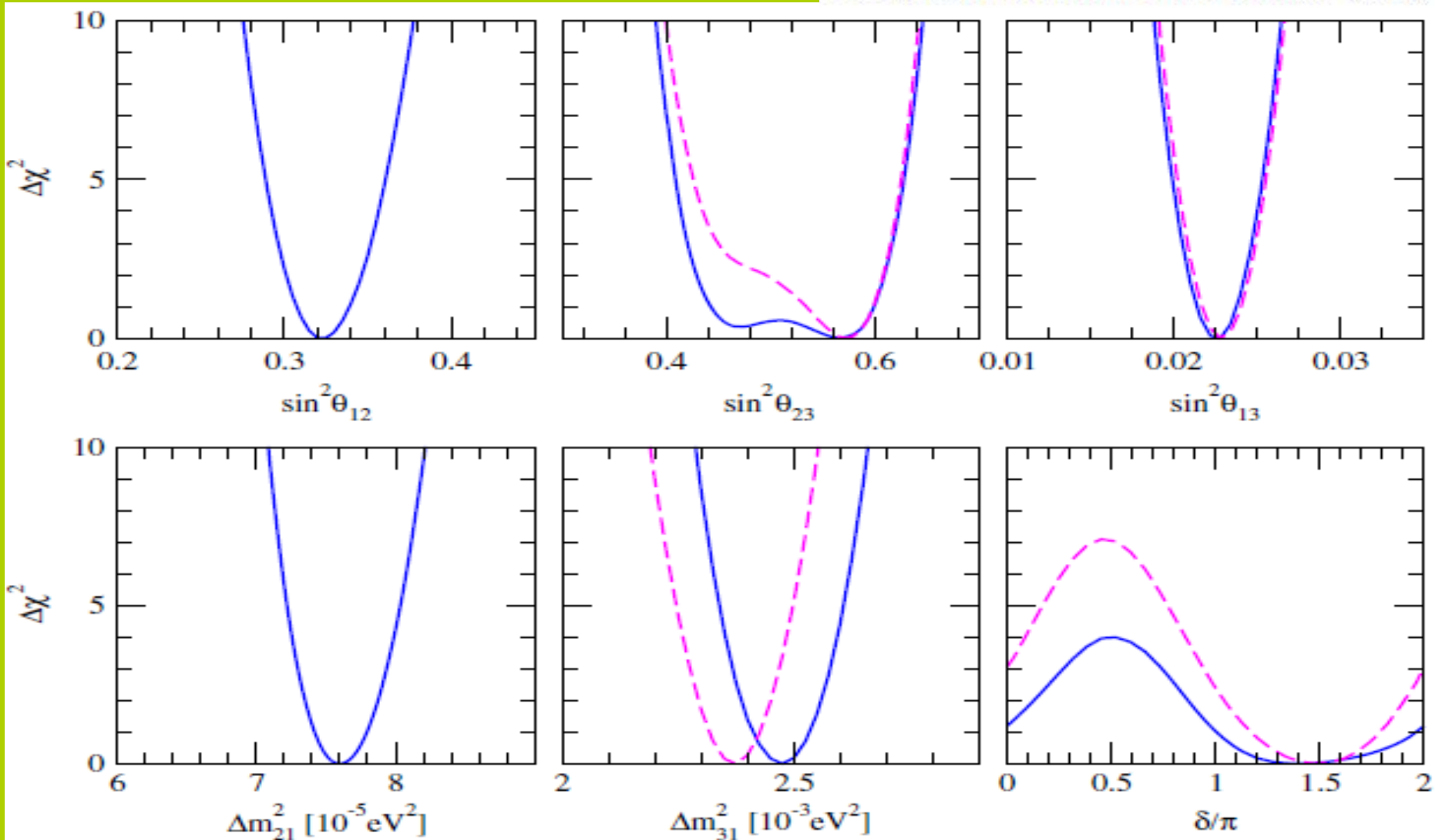
NCTS Annual Theory Meeting. 2016 Hsinchu, Taiwan



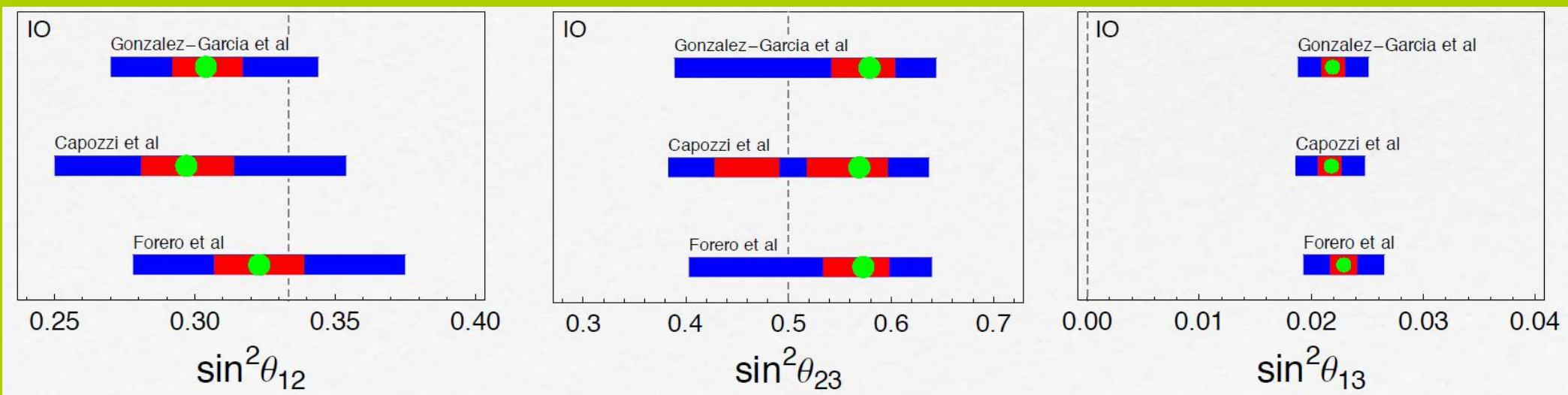
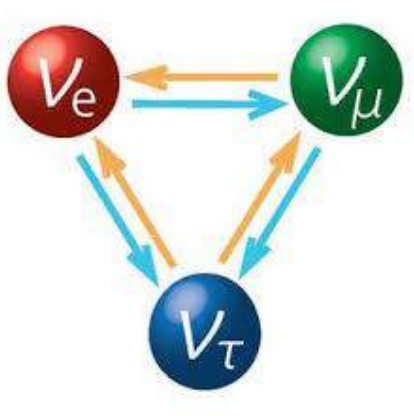
Oscillation parameters

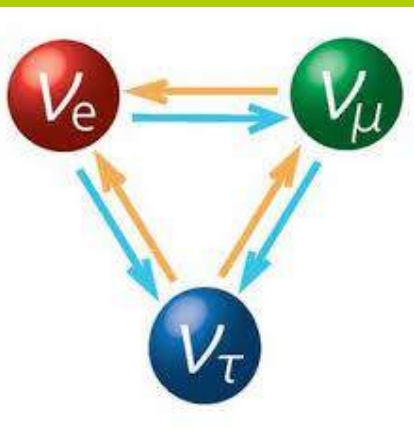
Precision era starts

PHYSICAL REVIEW D 90, 093006 (2014)



neutrino oscillation parameters

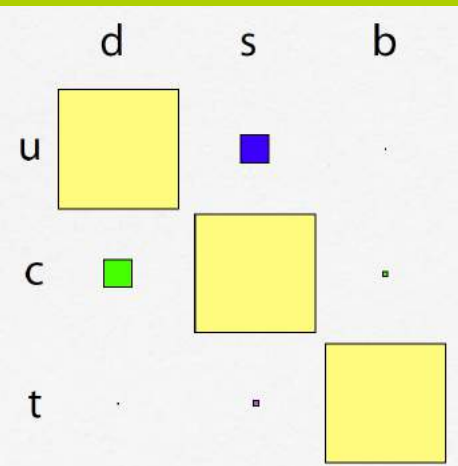




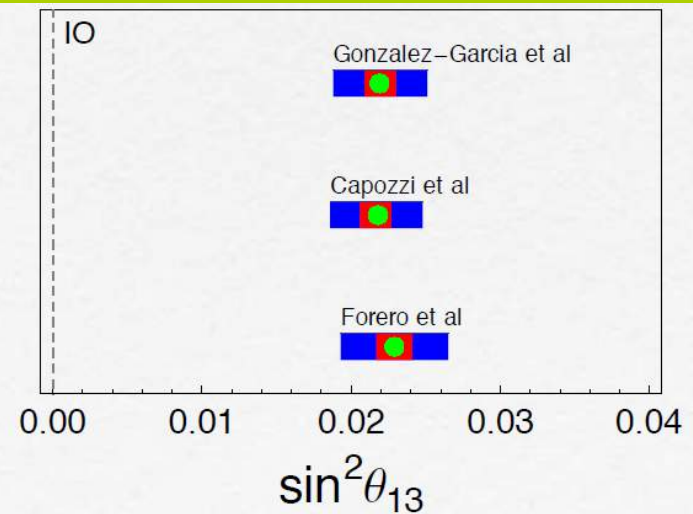
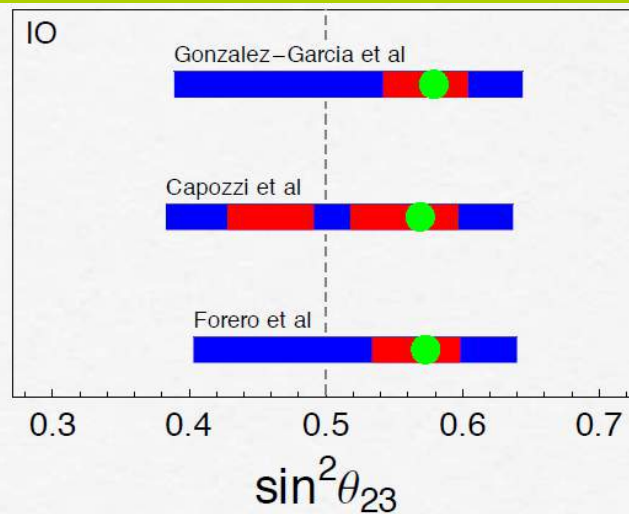
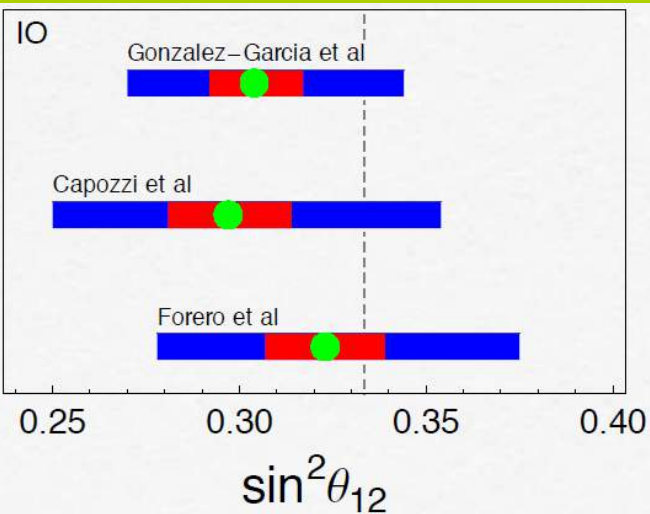
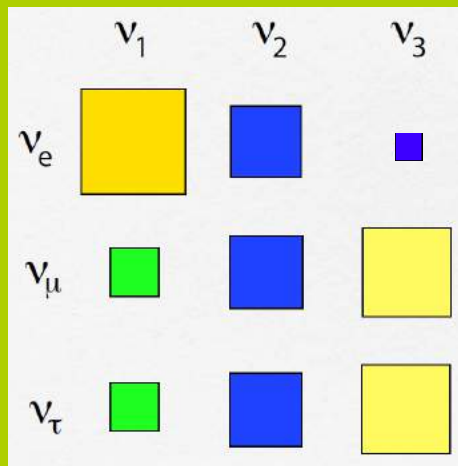
neutrino oscillation parameters

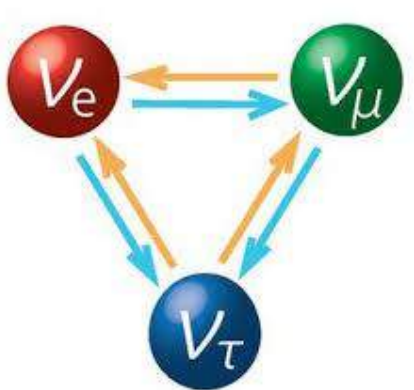
Phys.Lett. B748 (2015) 1-4

Phys.Rev. D86 (2012) 051301



VS

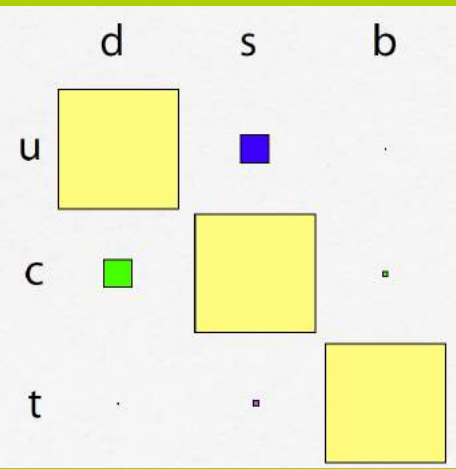




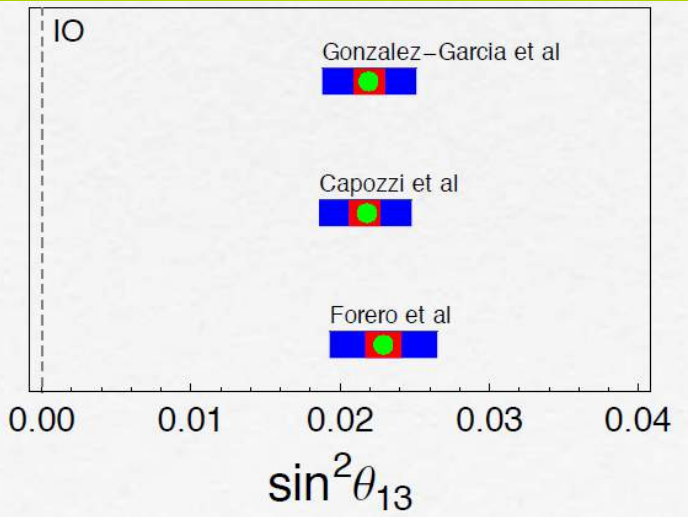
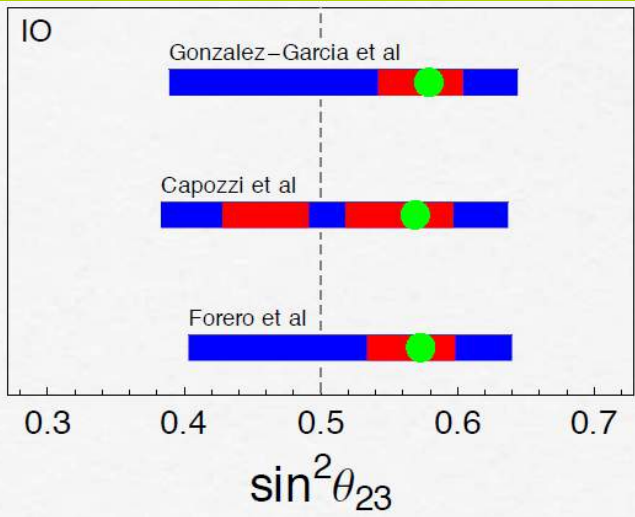
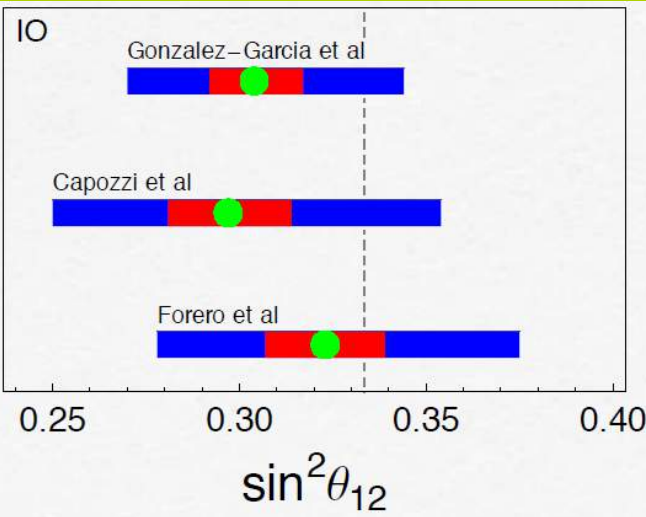
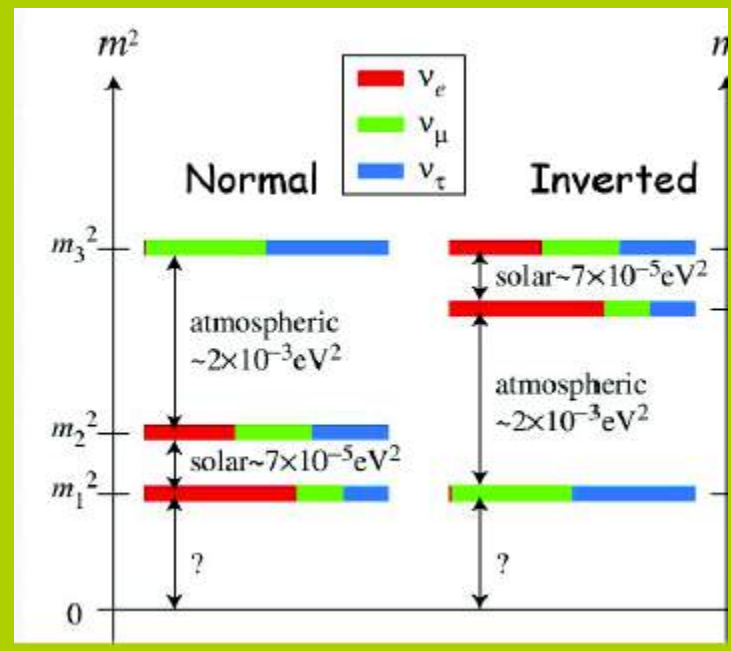
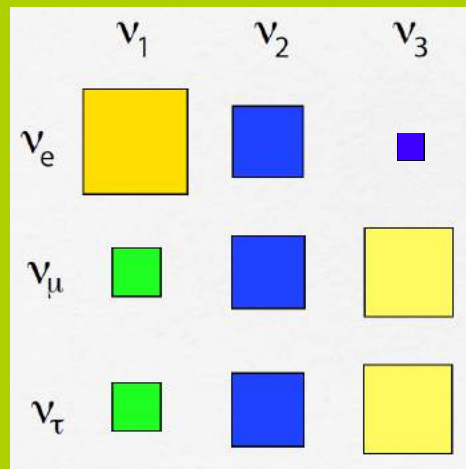
neutrino oscillation parameters

Phys.Lett. B748 (2015) 1-4

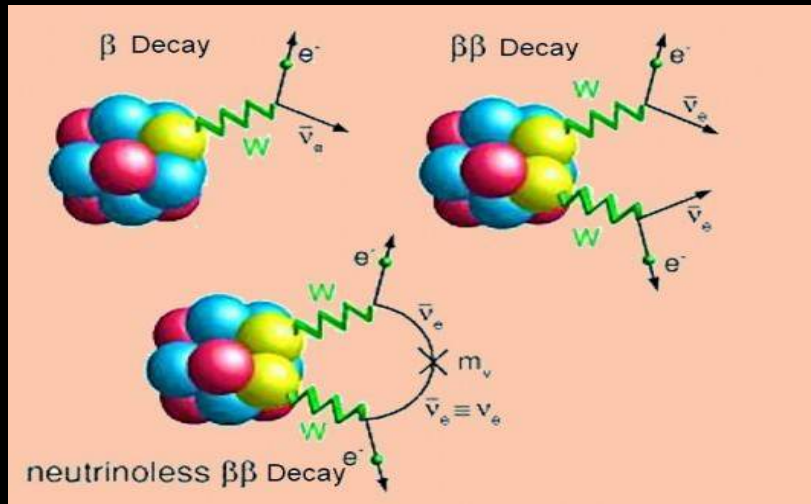
Phys.Rev. D86 (2012) 051301



ν_s

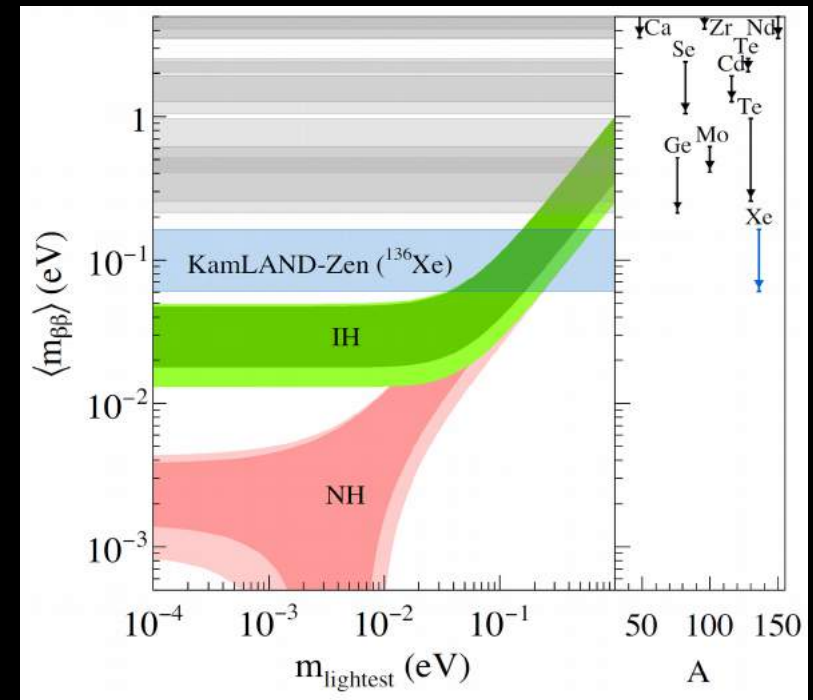
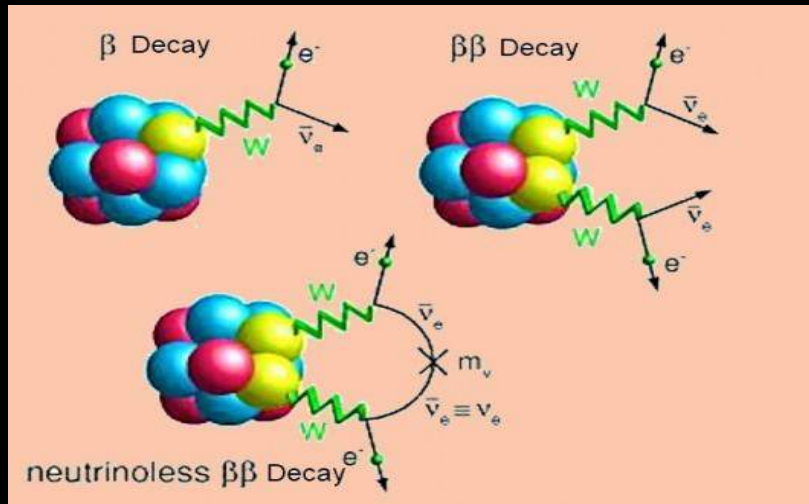


beta & double beta decay



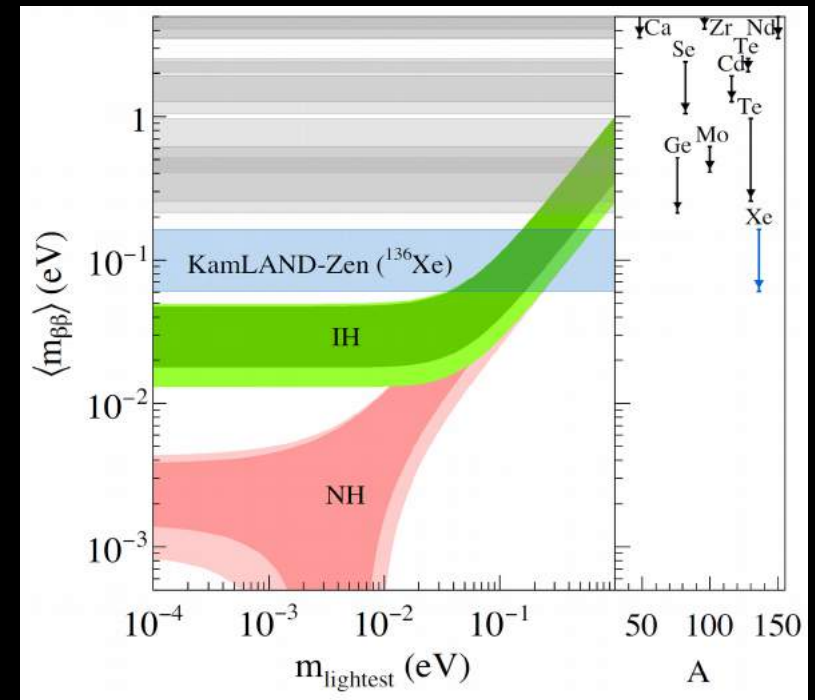
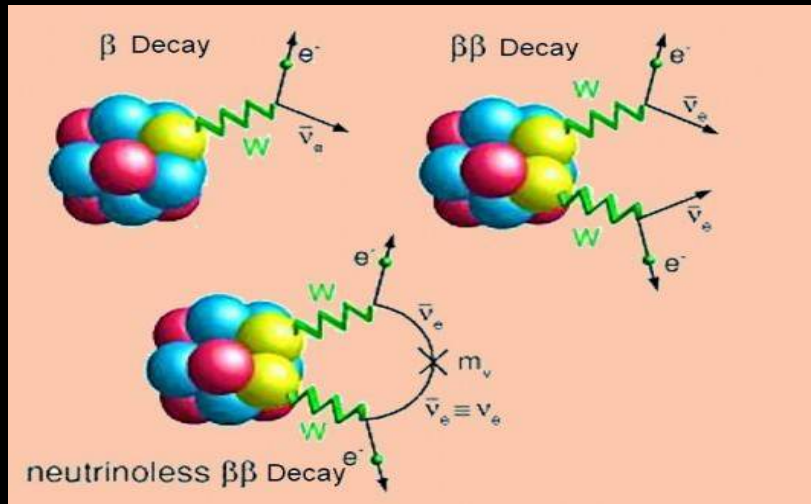
A.S. Barabash arXiv:1104.2714

beta & double beta decay

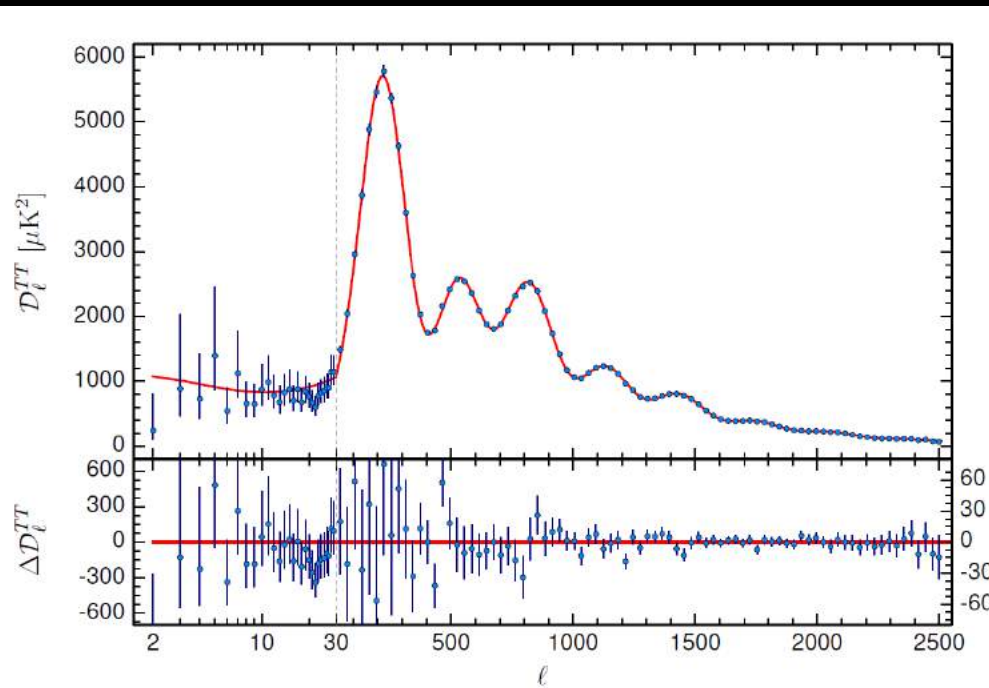


A.S. Barabash arXiv:1104.2714

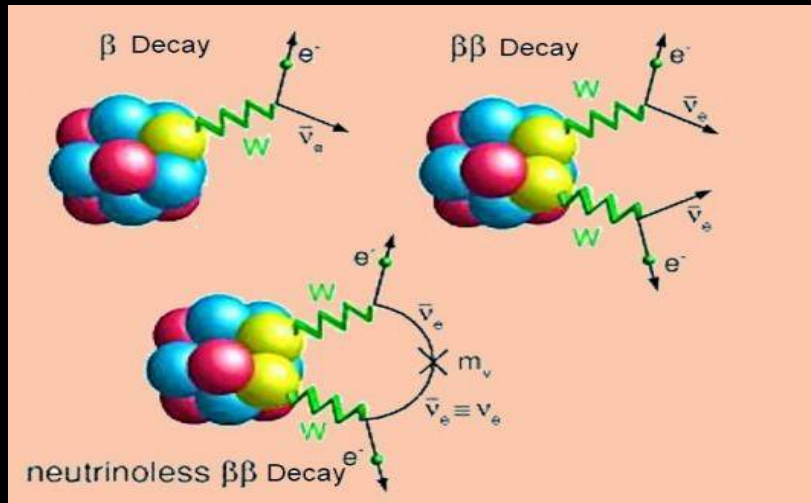
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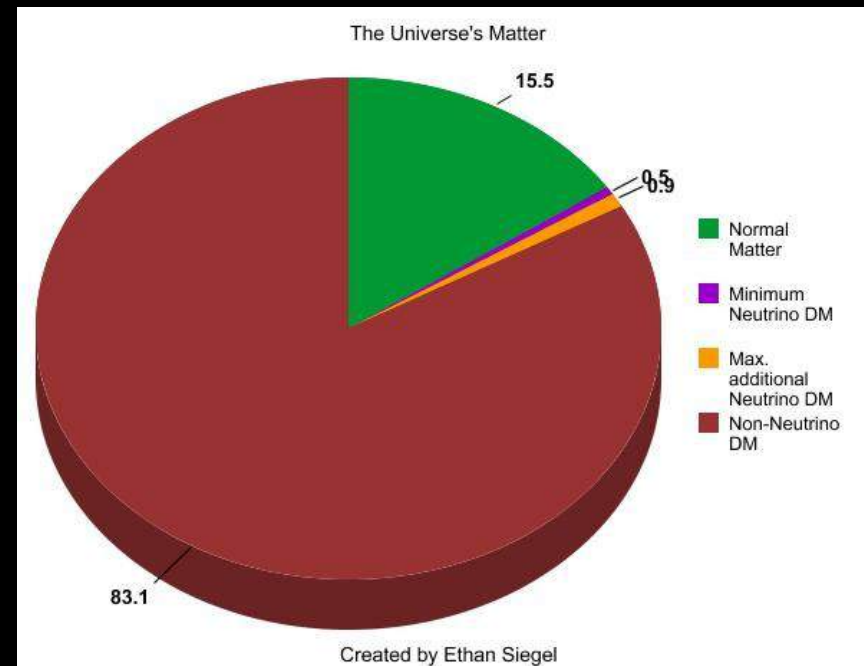
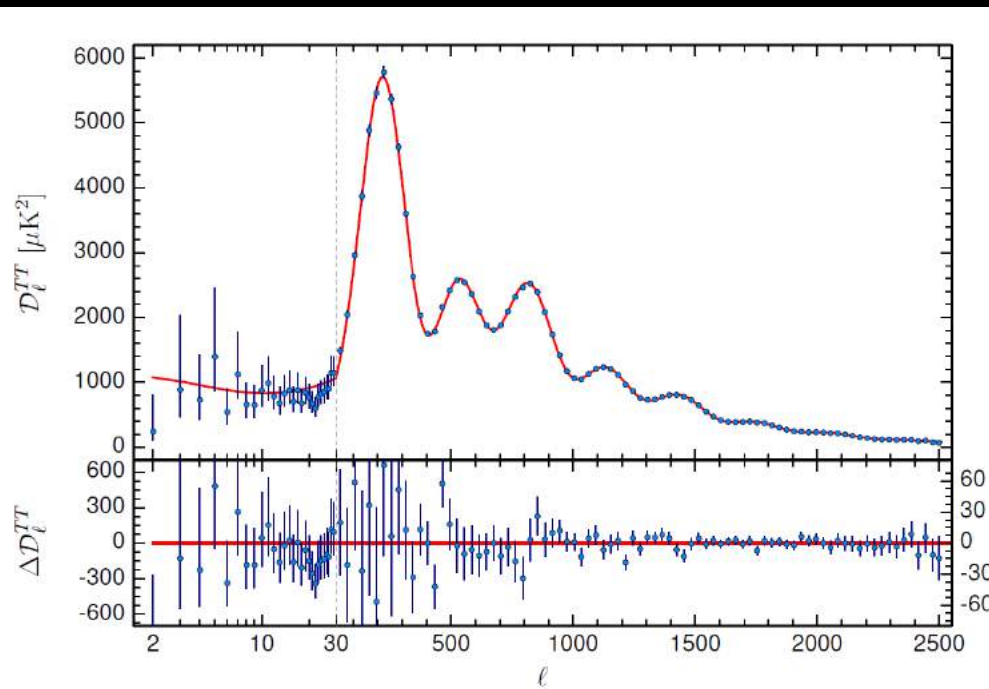
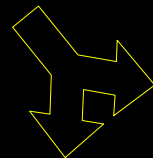
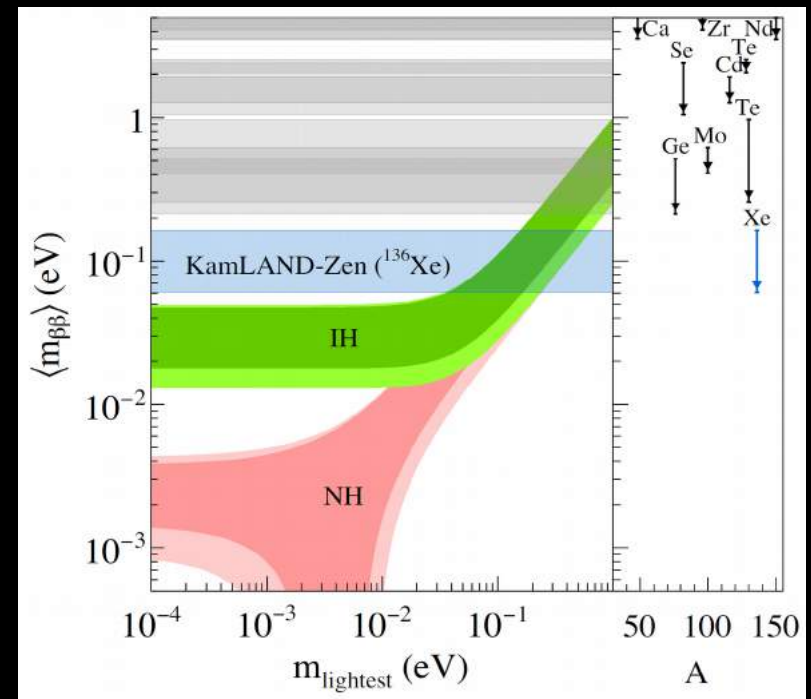
A.S. Barabash arXiv:1104.2714



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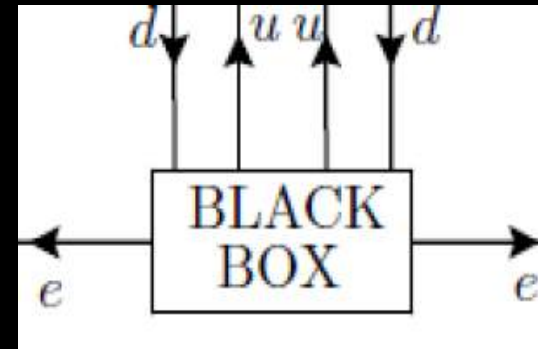


A.S. Barabash arXiv:1104.2714





The Majorana connection



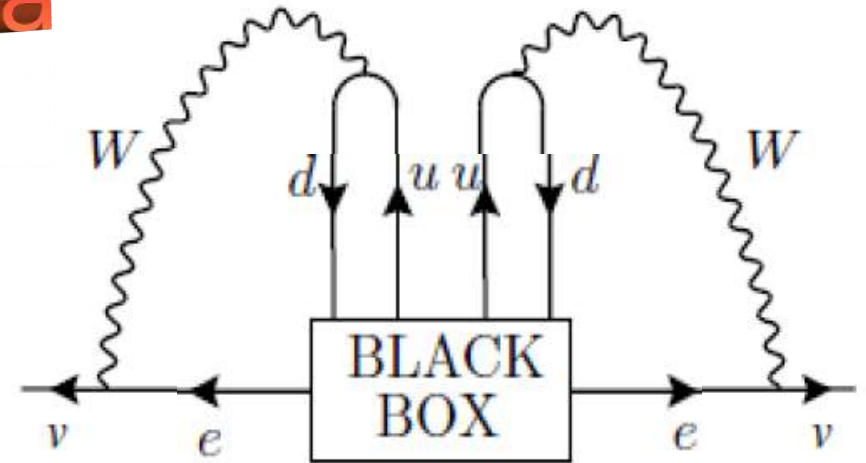
Schechter, JWFV 82

Lindner et al JHEP 1106 (2011) 091



The Majorana connection

*Even if mediated by
short-range mechanism ...
Heavy mediators*



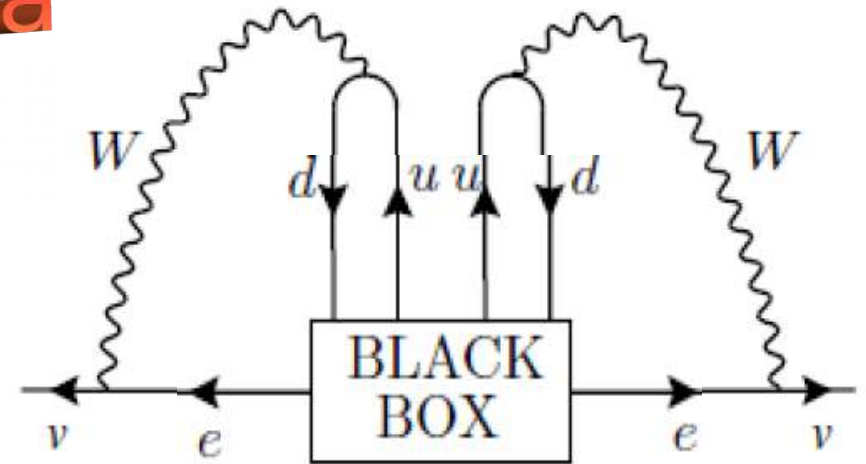
Schechter, JWFV 82

Lindner et al JHEP 1106 (2011) 091

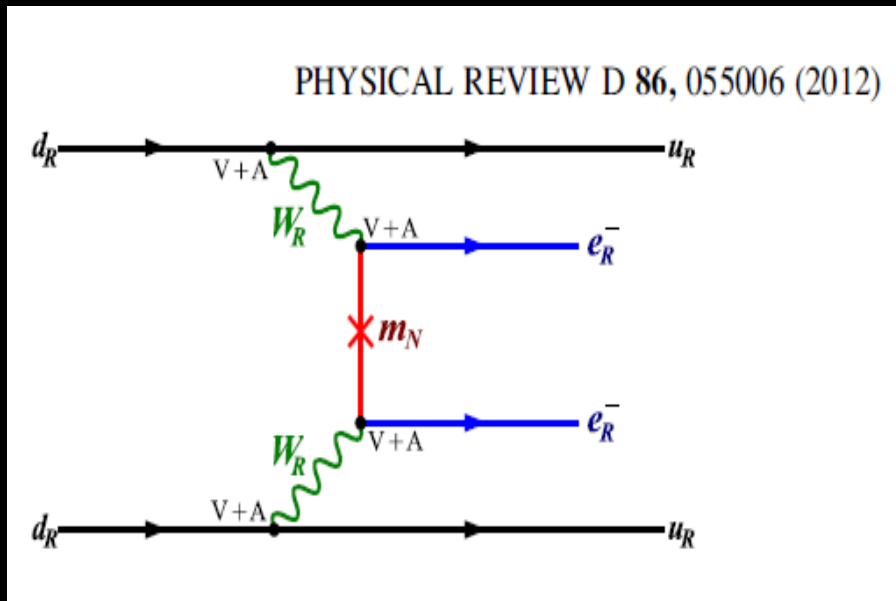


The Majorana connection

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Schechter, JWFV 82
Lindner et al JHEP 1106 (2011) 091





Standard model

THE STANDARD MODEL

FERMIONS (matter)

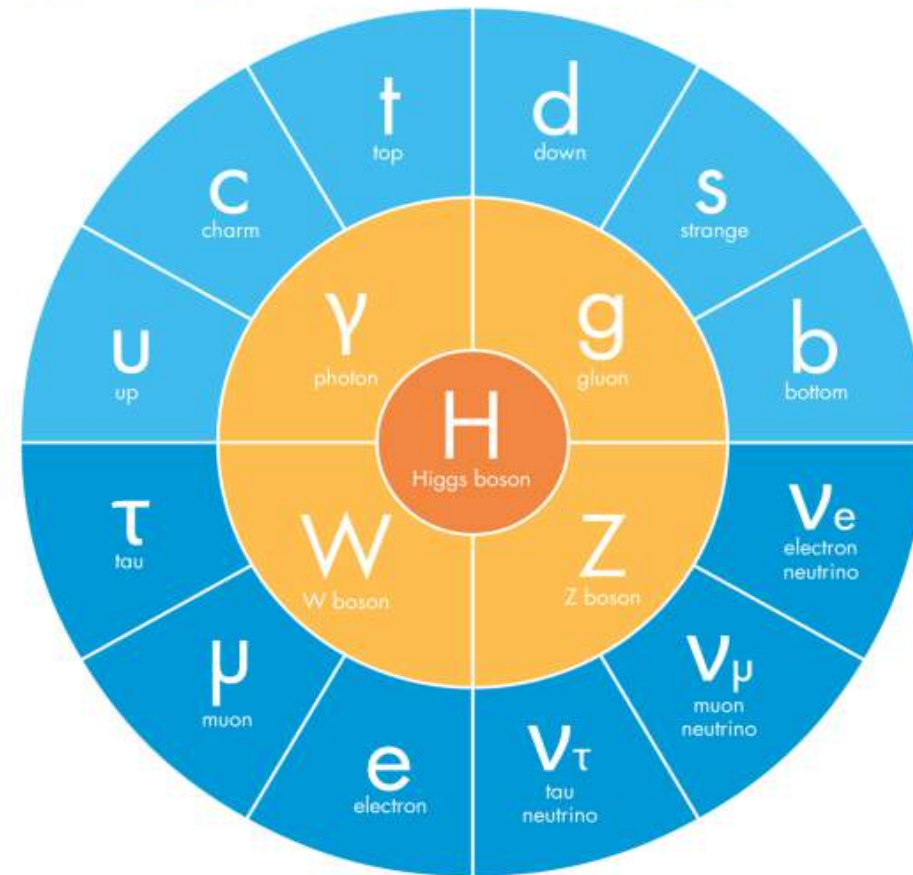
● Quarks

● Leptons

BOSONS (force carriers)

● Gauge bosons

● Higgs boson



However exciting ...



Higgs not the last brick !

Standard model

- Anomalies,
- unification,
- consistency of SSB,
- Gravity
- Flavor

THE STANDARD MODEL

FERMIONS (matter)

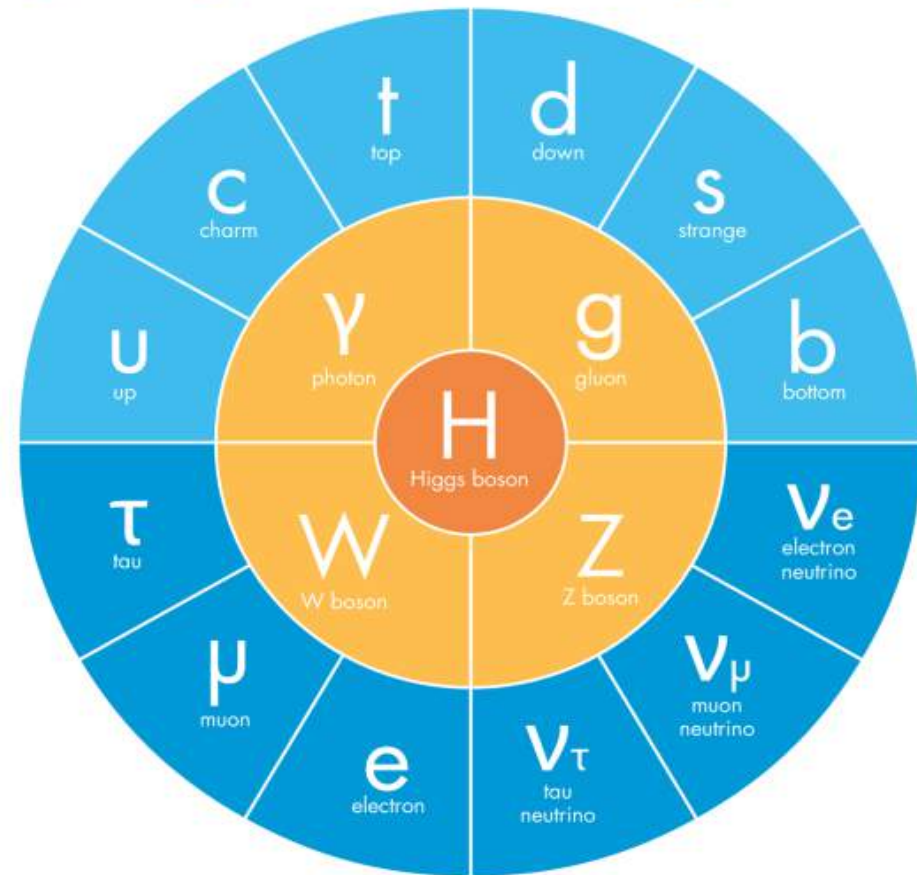
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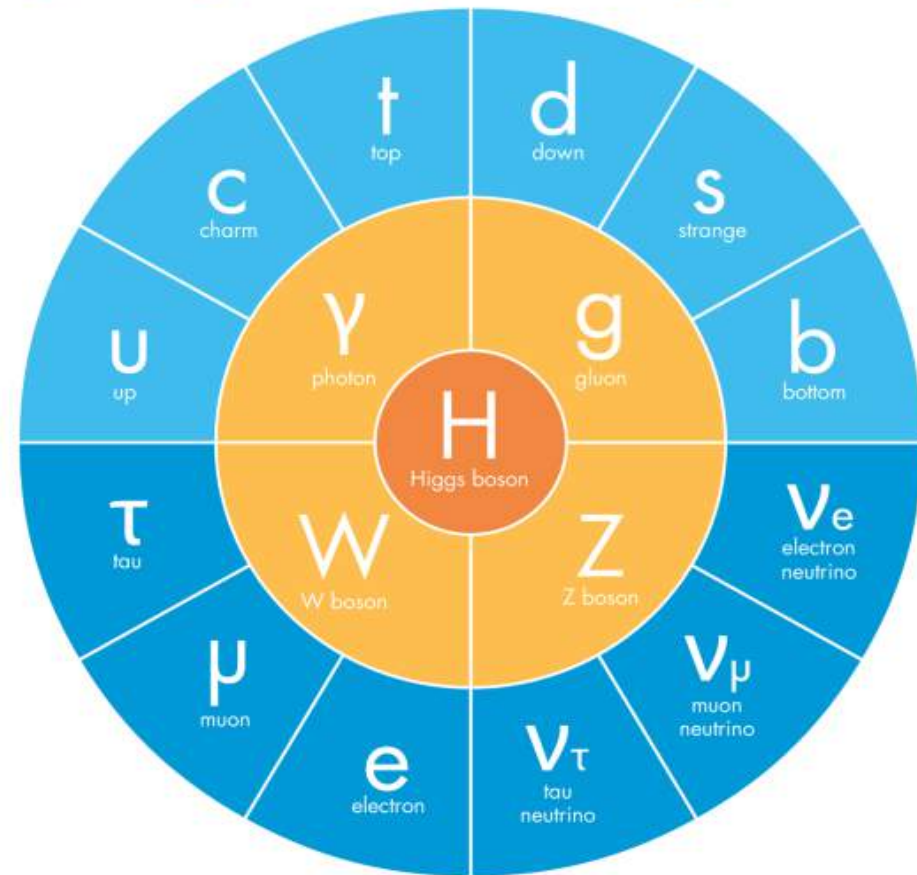
Quarks

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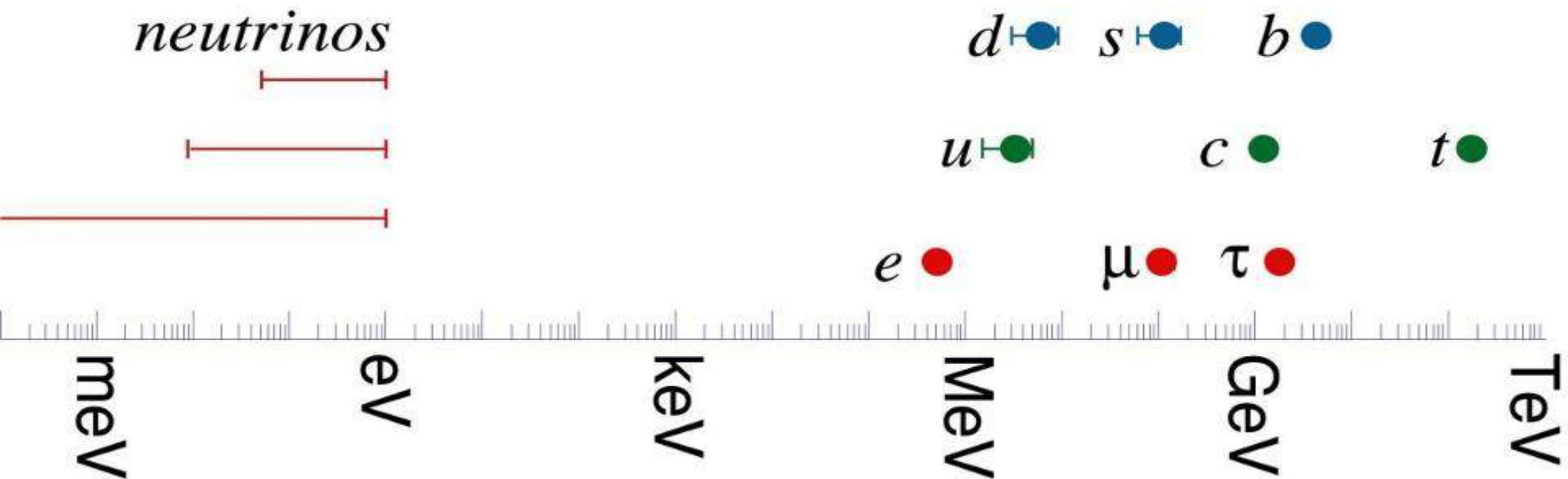
Gauge bosons

Higgs boson



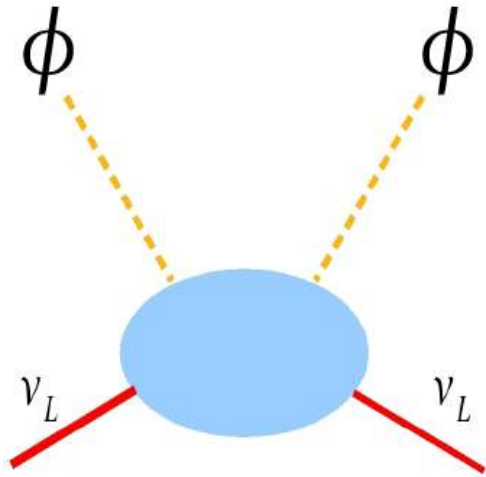
Despite its great success
SM does not explain neutrinos

... *nor cosmo: dark matter, inflation,
LG, dark energy*



Why tiny
neutrino masses

The origin of neutrino mass

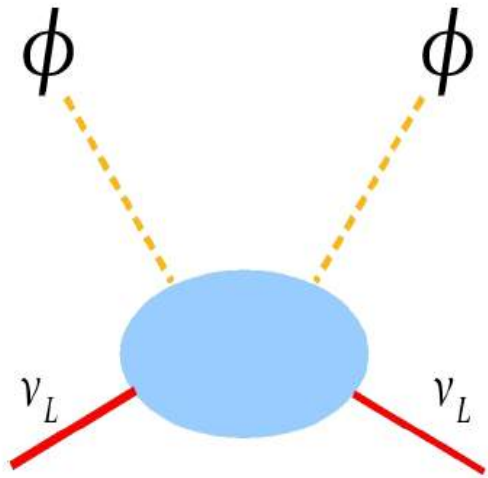


MECHANISM

SCALE

FLAVOR STRUCTURE

The origin of neutrino mass



MECHANISM

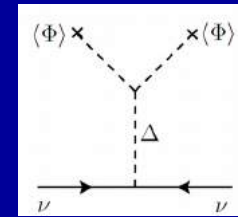
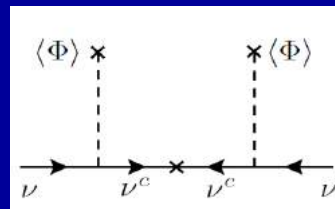
SCALE

FLAVOR STRUCTURE

Seesaw

$$v_3 v_1 \sim v_2^2$$

The origin of neutrino mass

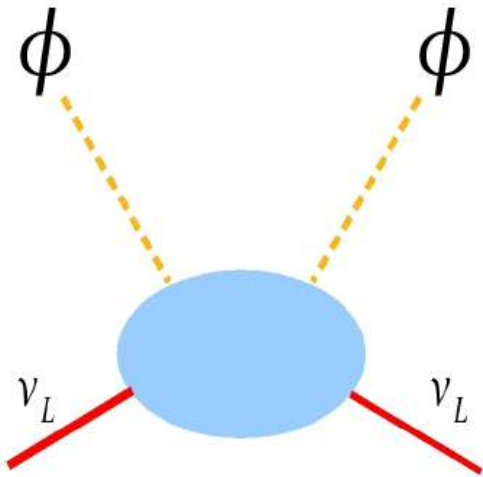


TYPE I

Minkowski 77
Gellman Ramond Slansky 80
Glashow, Yanagida 79
Mohapatra Senjanovic 80
Lazarides Shafi Wetterich 81
Schechter-Valle, 80 & 82

TYPE II

Schechter-Valle 80/82



MECHANISM

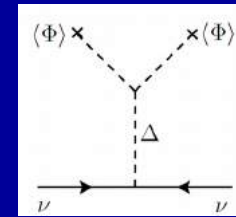
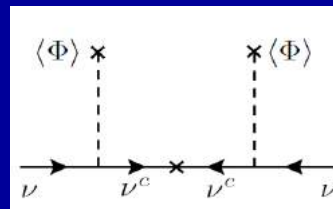
SCALE

FLAVOR STRUCTURE

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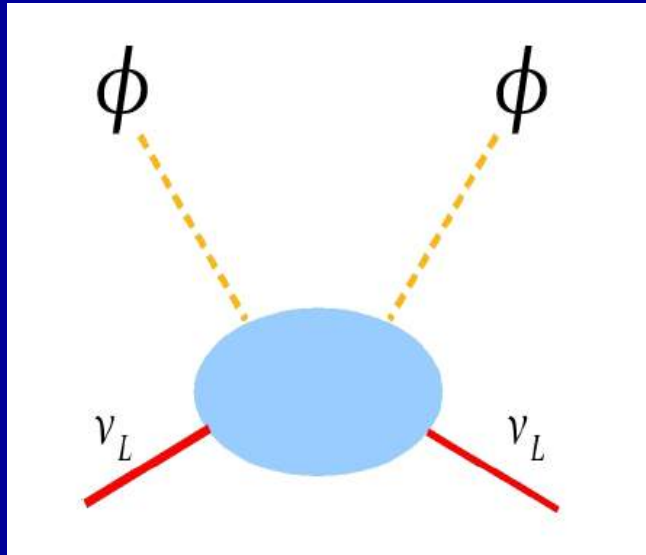


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MECHANISM

SCALE

FLAVOR STRUCTURE

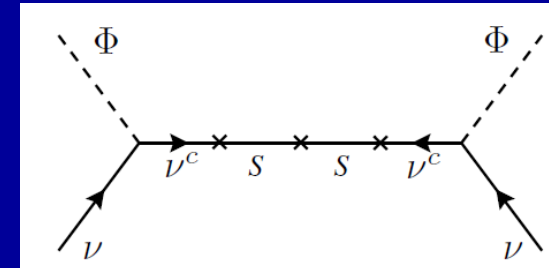
Seesaw

$$v_3 v_1 \sim v_2^2$$

Number & properties of messengers

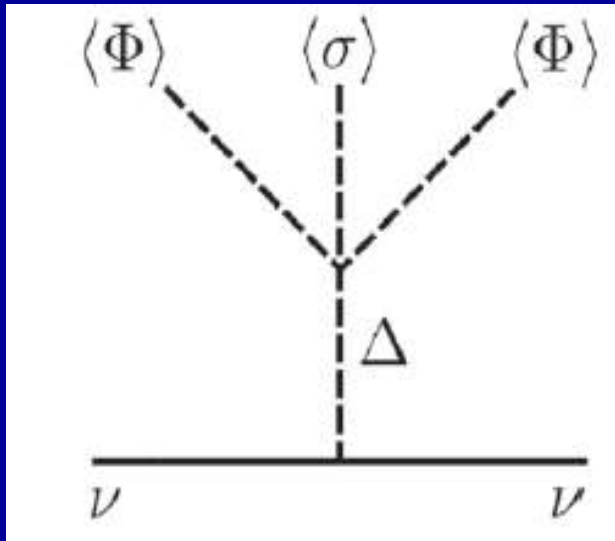
LOW-SCALE SEESAW

Mohapatra-Valle 86
Akhmedov et al PRD53 (1996) 2752
Malinsky et al PRL95(2005)161801
Bazzocchi et al, PRD81 (2010) 051701

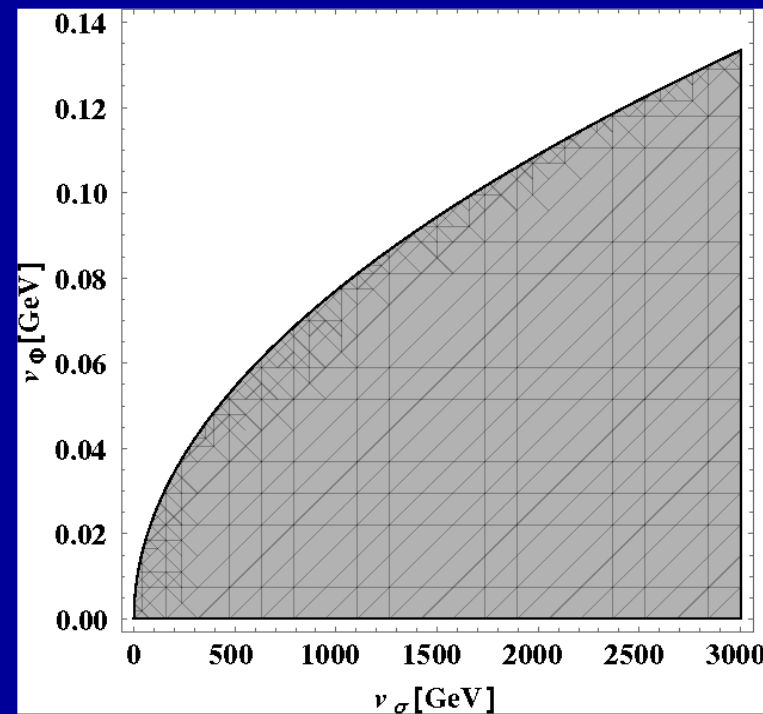


type-II seesaw with spont U(1) violation

Phys.Rev. D25 (1982) 774



Astrophysical limit



type-II seesaw with spont $U(1)$ violation

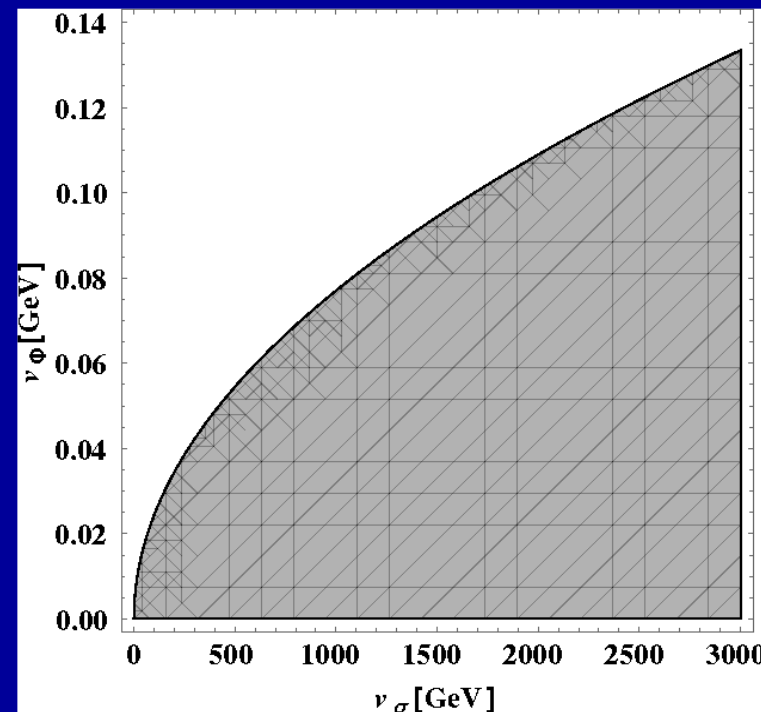
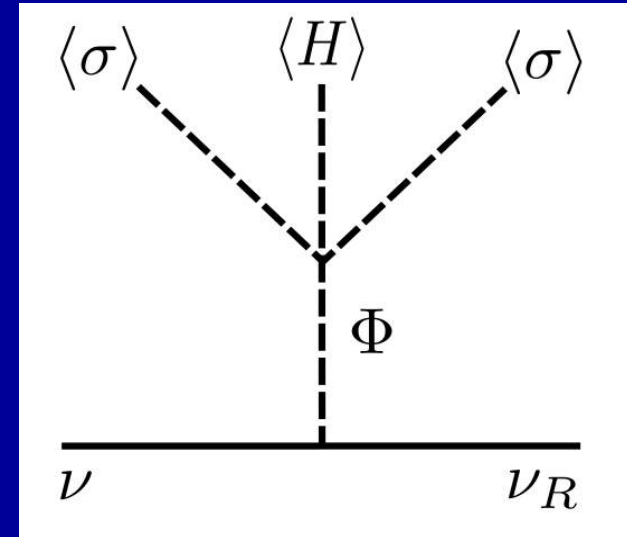
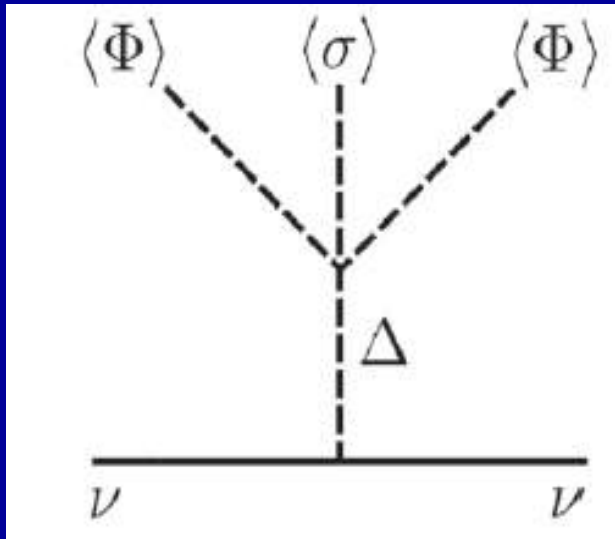
Phys.Rev. D25 (1982) 774

naturally small induced vev

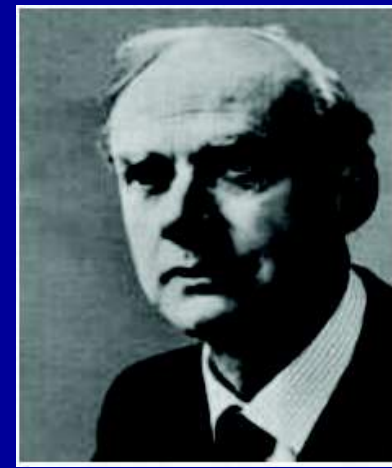
accidental global $U(1)$

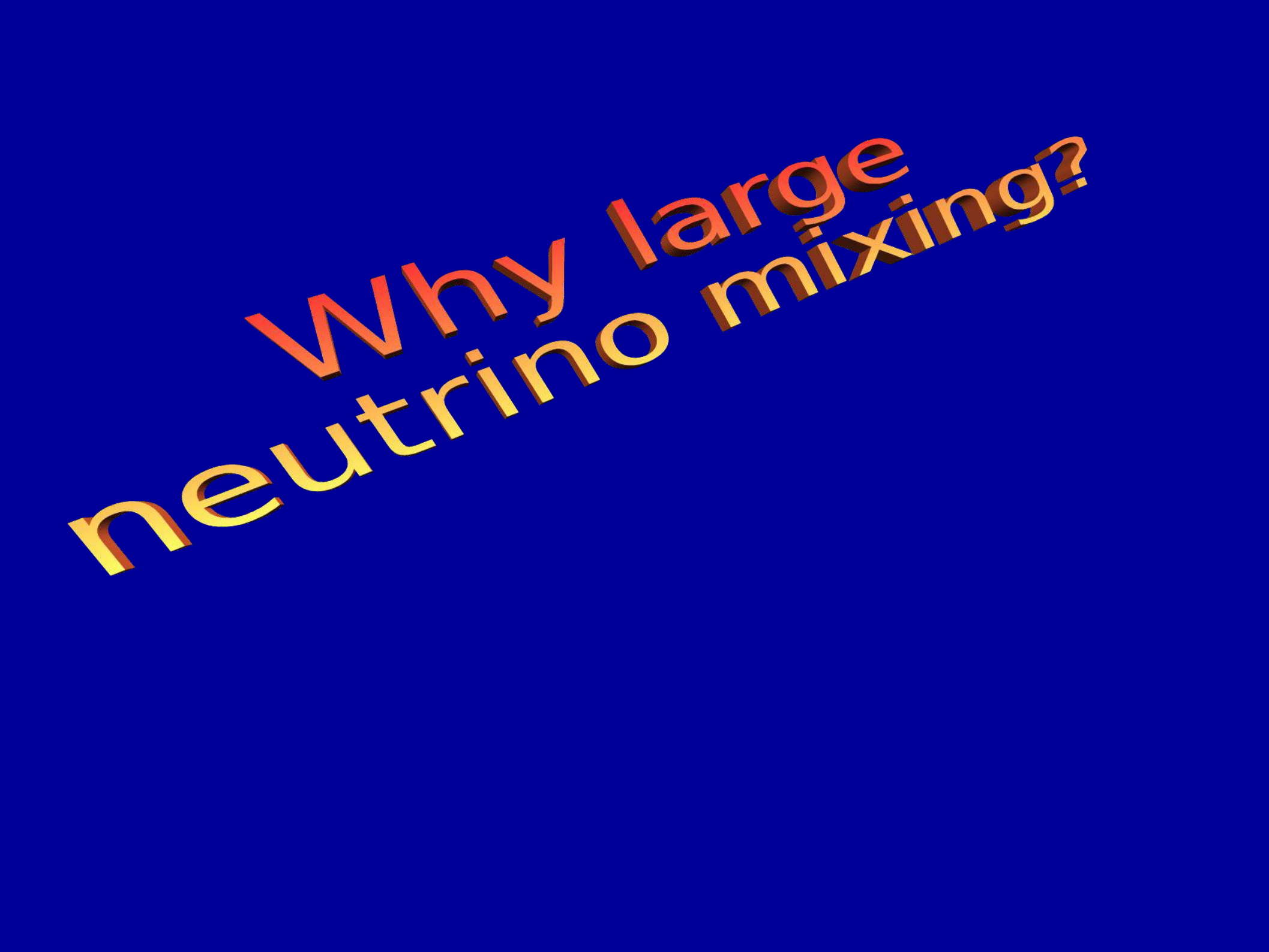
Majoron
vs
Diracon

Astrophysical limit



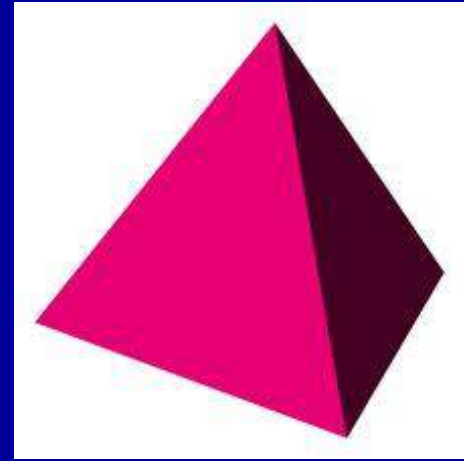
ArXiv:1605.08362 **PLB**





Why large
neutrino mixing?

Flavor Symmetry



A4

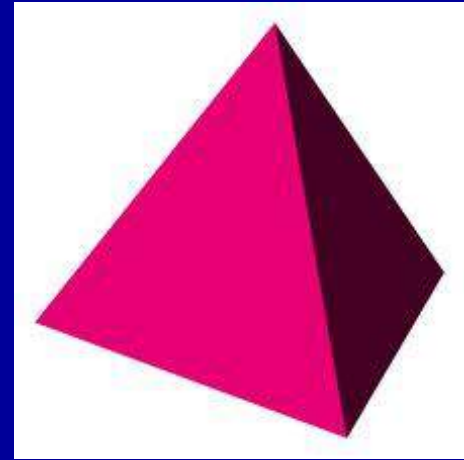
$$\begin{array}{ccc}
 \begin{pmatrix} \nu_e \\ e \\ e_R \end{pmatrix}_L & \begin{pmatrix} \nu_\mu \\ \mu \\ \mu_R \end{pmatrix}_L & \begin{pmatrix} \nu_\tau \\ \tau \\ \tau_R \end{pmatrix}_L \\
 \begin{pmatrix} u \\ d \\ u_R \end{pmatrix}_L & \begin{pmatrix} c \\ s \\ c_R \end{pmatrix}_L & \begin{pmatrix} t \\ b \\ t_R \end{pmatrix}_L \\
 & & \begin{pmatrix} d_R \\ s_R \\ b_R \end{pmatrix}
 \end{array}$$

Babu-Ma-Valle PLB552 (2003) 207
 Hirsch et al PRD69 (2004) 093006

$$\sin^2 \theta_{23} = 0.5$$

$$\sin^2 \theta_{13} = 0$$

Flavor Symmetry



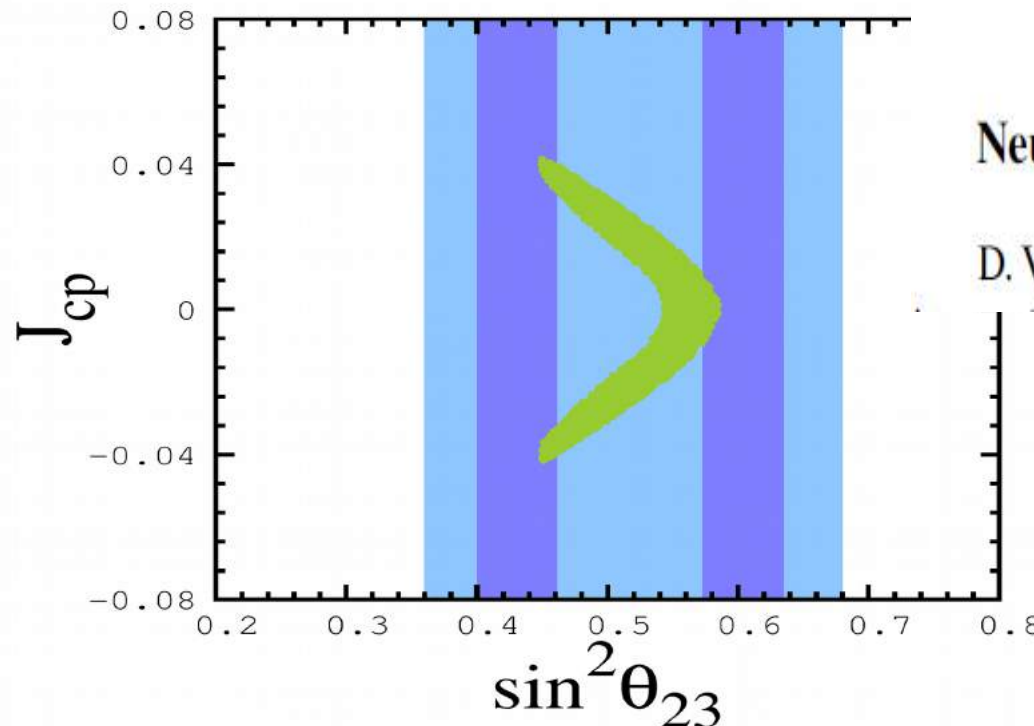
A4

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 \begin{pmatrix} u \\ d \\ u_R \end{pmatrix}_L & \begin{pmatrix} c \\ s \\ c_R \end{pmatrix}_L & \begin{pmatrix} t \\ b \\ t_R \end{pmatrix}_L \\
 \begin{pmatrix} u \\ d \\ u_R \end{pmatrix}_R & \begin{pmatrix} c \\ s \\ c_R \end{pmatrix}_R & \begin{pmatrix} t \\ b \\ t_R \end{pmatrix}_R
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Babu-Ma-Valle PLB552 (2003) 207
Hirsch et al PRD69 (2004) 093006

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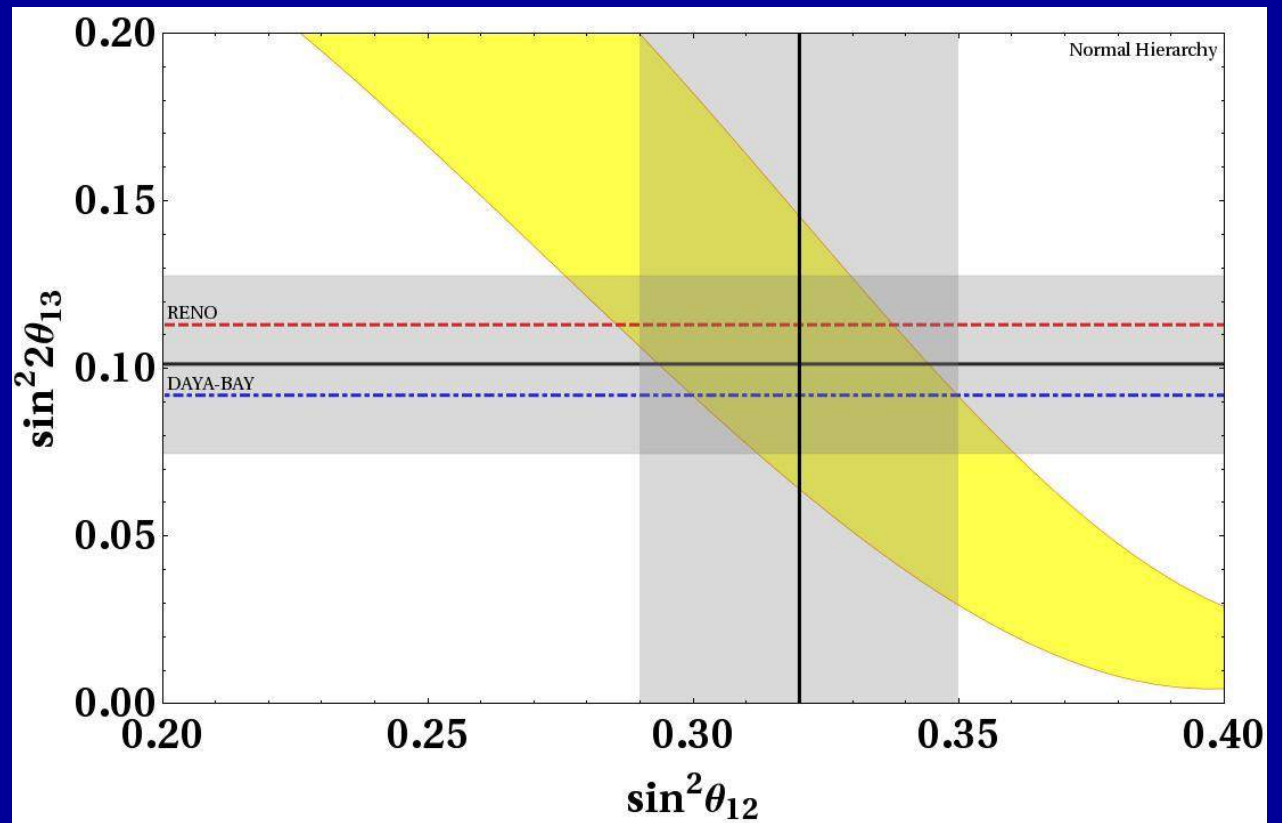
PHYSICAL REVIEW D 88, 016003 (2013)

Neutrino mixing with revamped A_4 flavor symmetry

D. V. Forero,^{1,2,*} S. Morisi,^{3,†} J. C. Romão,^{1,‡} and J. W. F. Valle^{2,§}

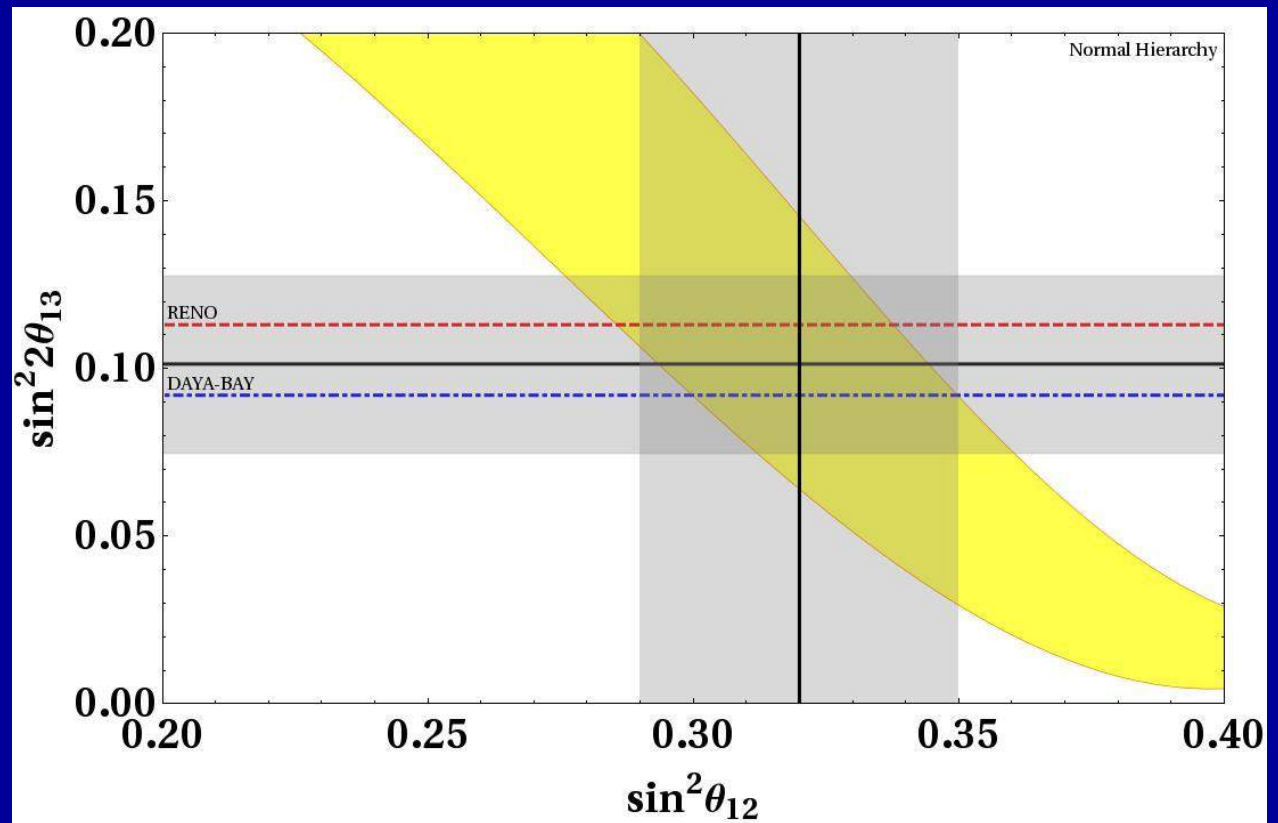
Flavor correlations

Boucenna et al
PhysRevD.86.073008



Flavor correlations

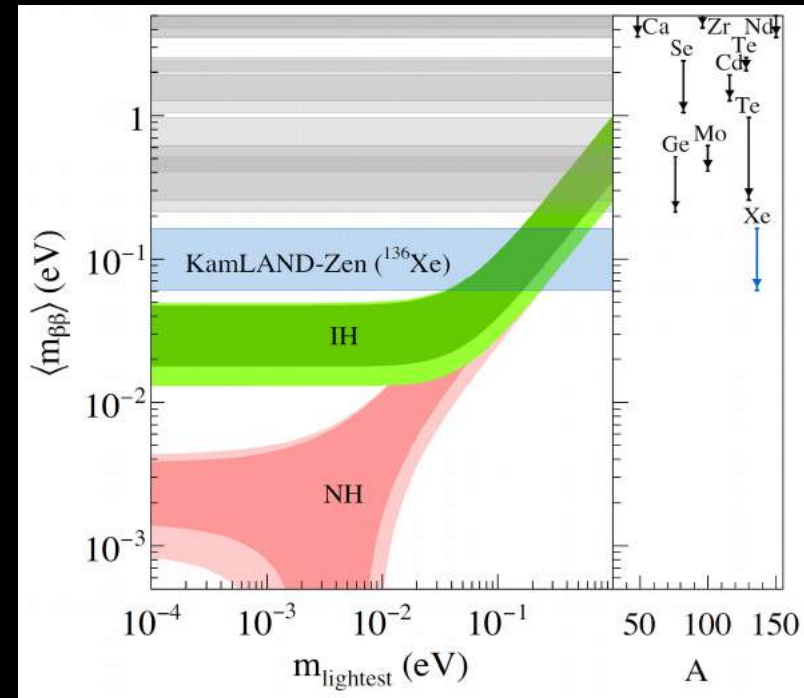
Boucenna et al
PhysRevD.86.073008



Model-independent flavor approach

P Chen et al
Phys.Lett. B753 (2016) 644-652
Phys.Rev. D94 (2016) no.3, 033002

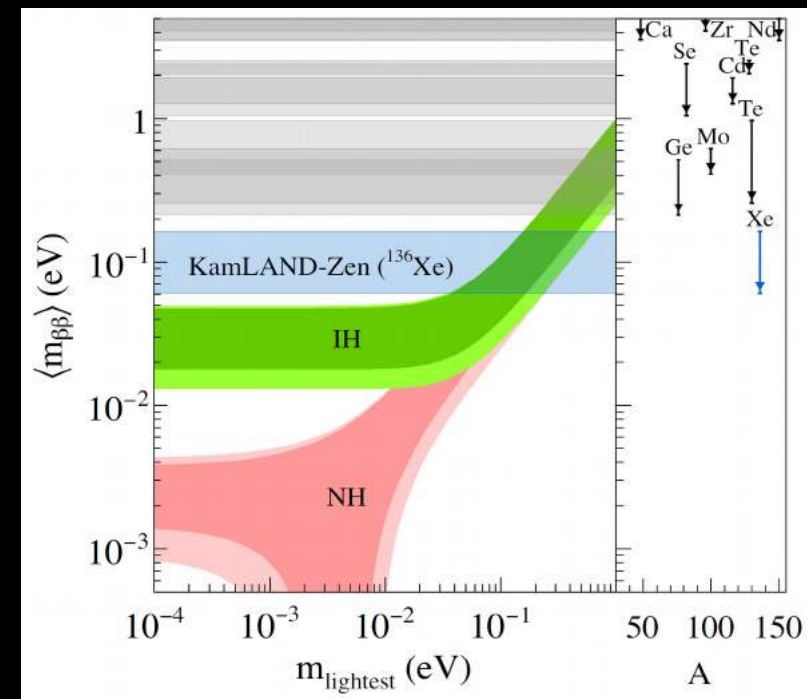
Flavor Sensitivity in neutrinoless double beta decay



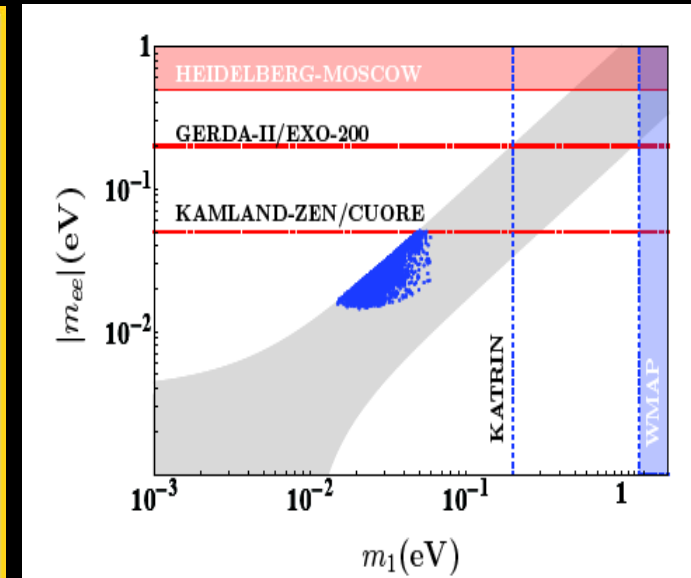
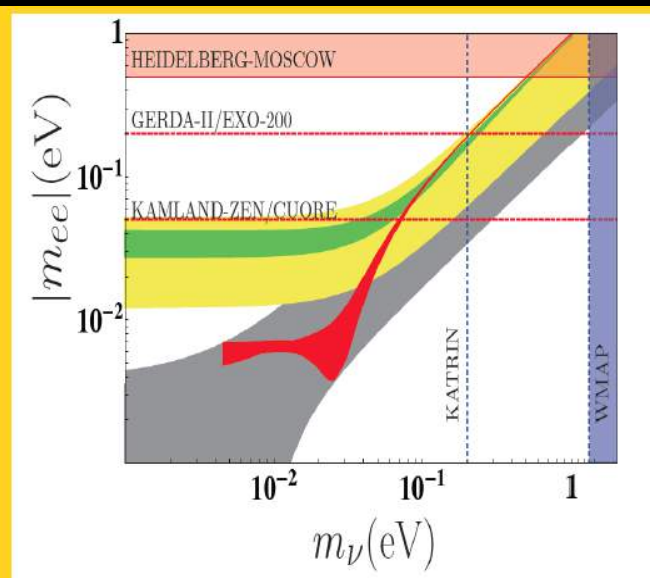
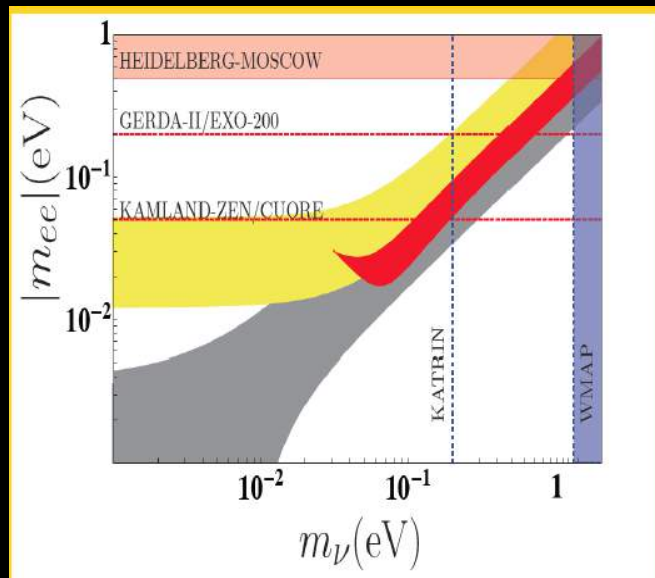
A.S. Barabash arXiv:1104.2714

Flavor Sensitivity in neutrinoless double beta decay

Lower bounds ...



A.S. Barabash arXiv:1104.2714

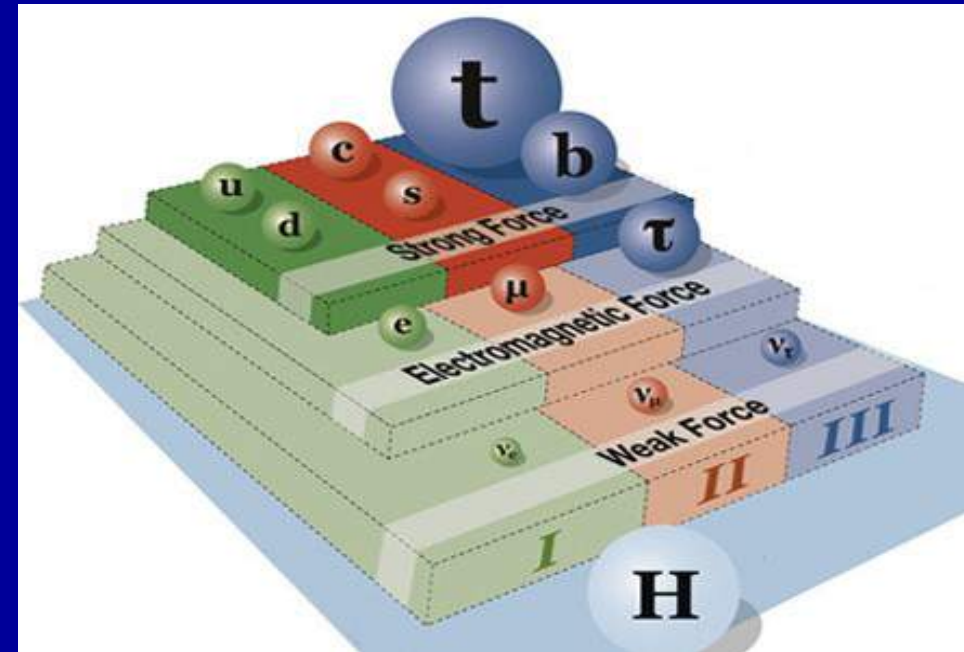


Dorame et al
NPB861 (2012) 259-270

PhysRevD.86.056001

King et al Phys. Lett. B 724 (2013) 68

Can neutrinos shed light on charged fermion masses?

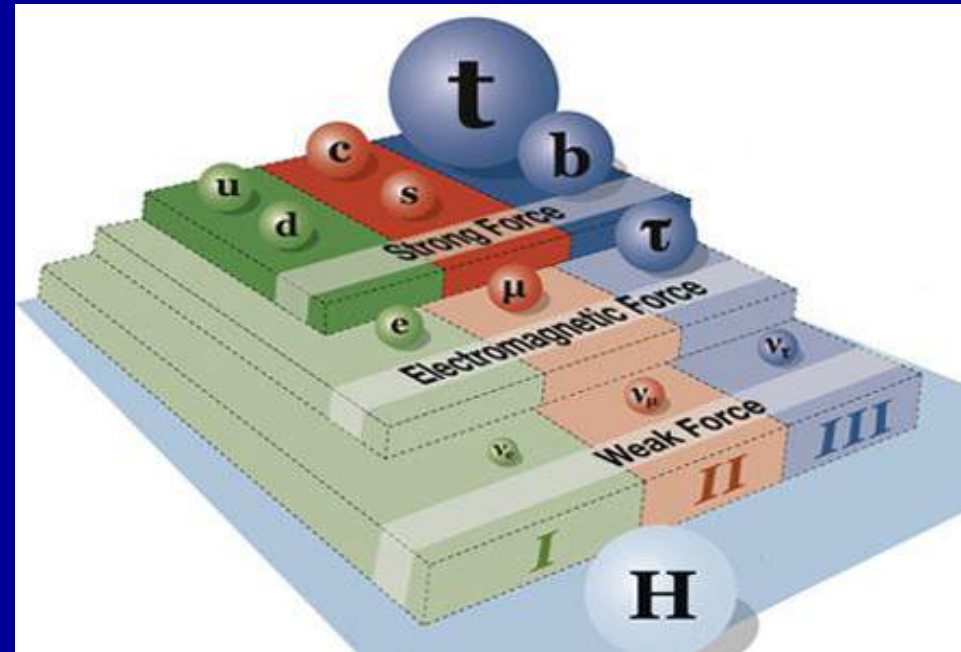


Neutrinos : Lepton number?

Can neutrinos shed light on charged fermion masses?

Flavor dependent b-tau unification

$$\frac{m_\tau}{\sqrt{m_e m_\mu}} \approx \frac{m_b}{\sqrt{m_d m_s}}$$



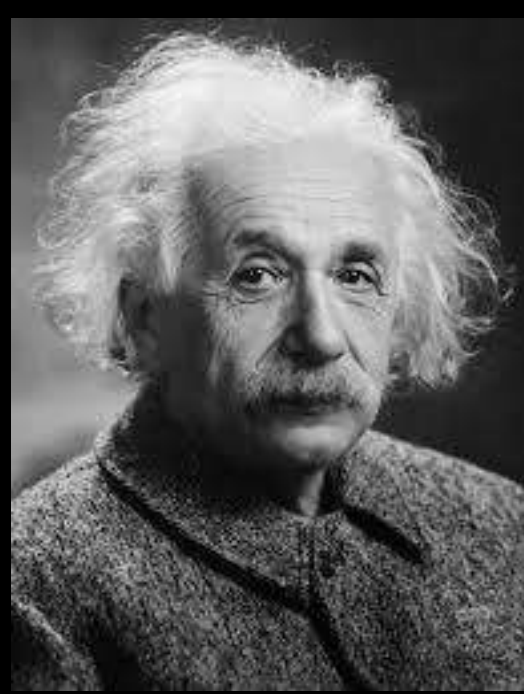
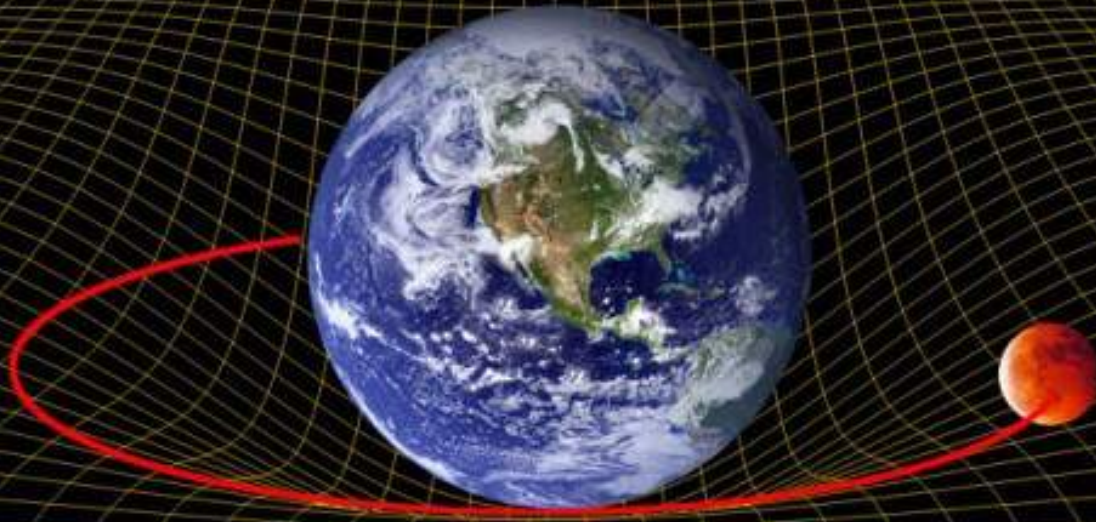
Morisi et al Phys.Rev. D84 (2011) 036003

King et al Phys. Lett. B 724 (2013) 68

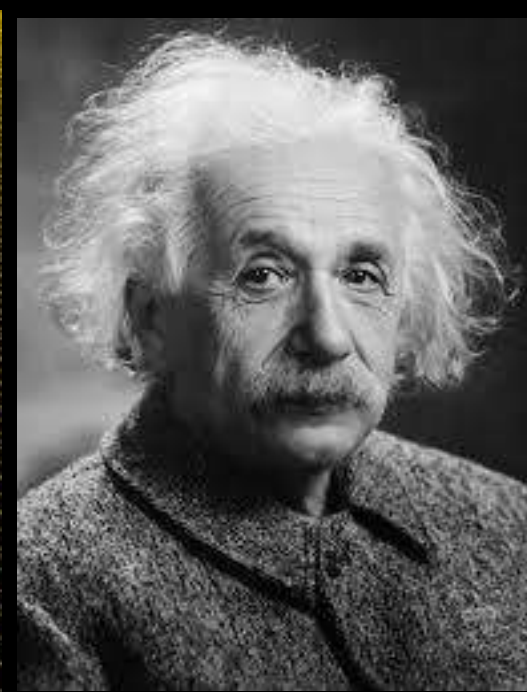
Morisi et al Phys.Rev. D88 (2013) 036001

Bonilla et al Phys.Lett. B742 (2015) 99

including Gravity



including Gravity



Neutrinos in the theory of everything

Chen et al arXiv:1509.06683
JHEP01(2016)007

Addazi et al

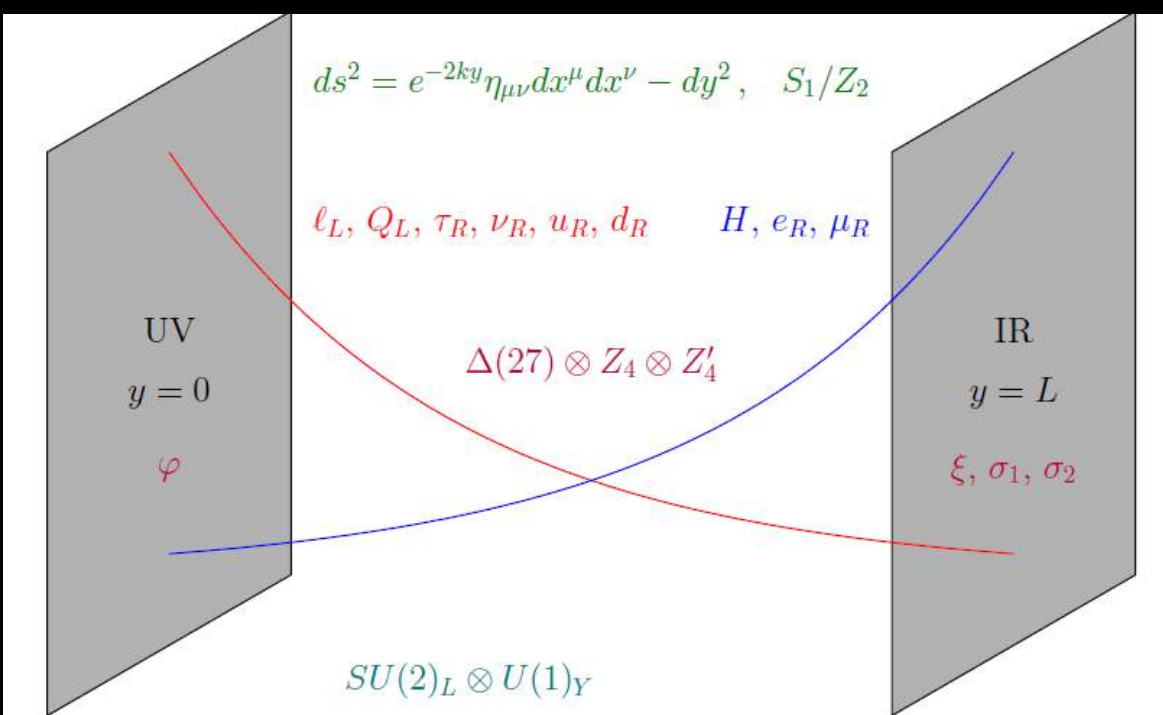
Phys.Lett. B759 (2016) 471-478



Warped standard model

Chen et al arXiv:1509.06683

JHEP01(2016)007

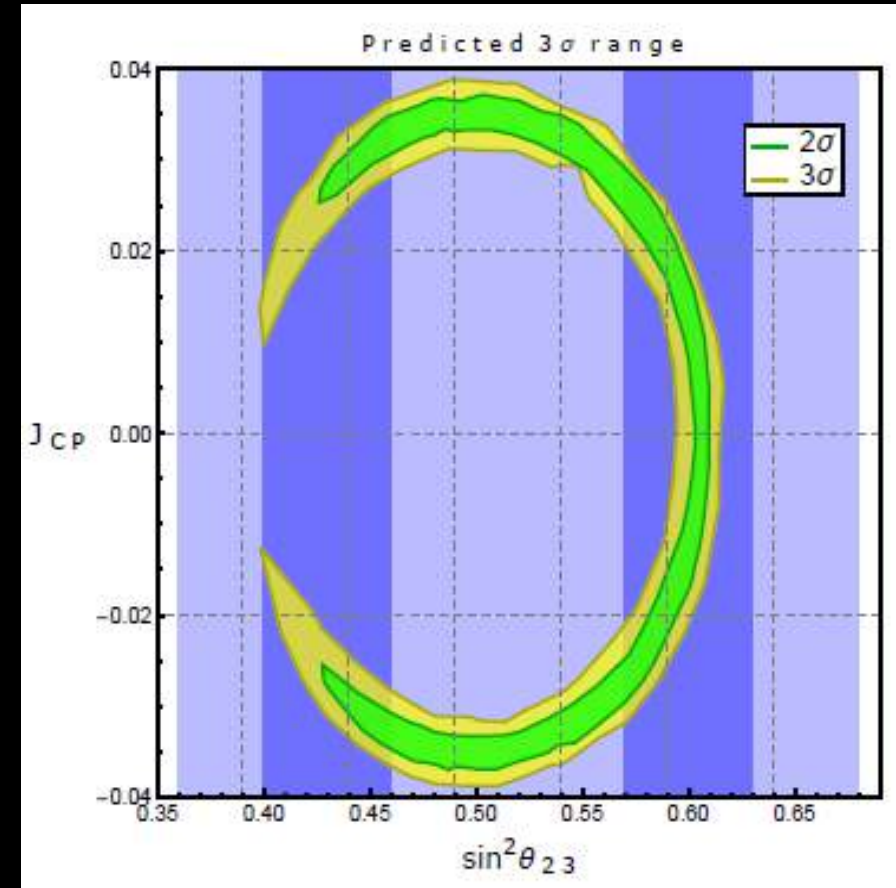
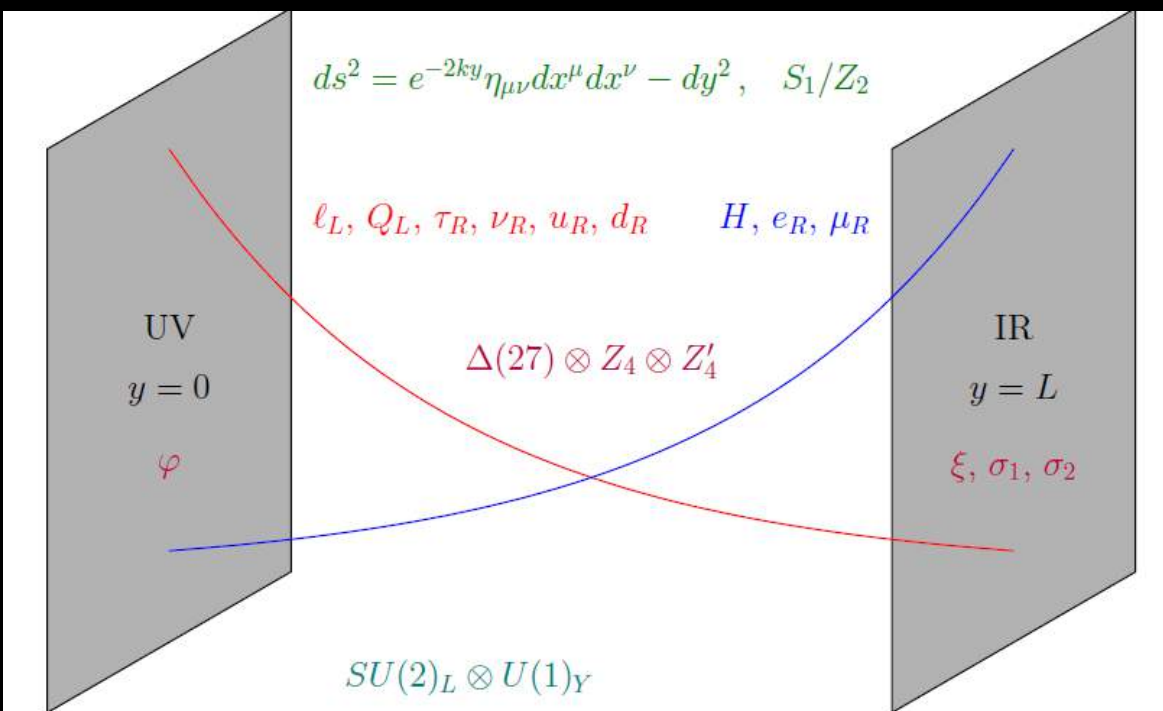


Masses explained by choices
of the bulk parameters

Warped standard model

Chen et al arXiv:1509.06683

JHEP01(2016)007



Masses explained by choices
of the bulk parameters

<http://arxiv.org/abs/arXiv:1610.05962>

Warped flavor predictions

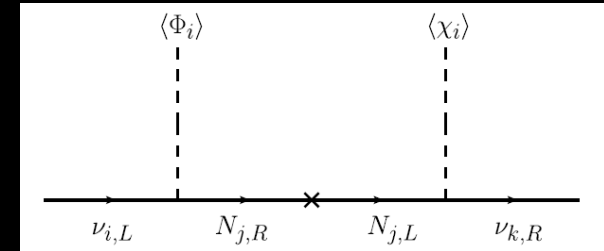
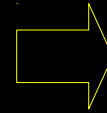
even if neutrinos form only
tiny DM fraction

they can hold the key
to Dark matter problem

Dark Matter Stability from Diracness

Chiulia et al arXiv:1606.04543
Phys.Lett. B761 (2016) 431

Fields	Z_4	Z_2	Fields	Z_4	Z_2
$\bar{L}_{i,L}$	$\mathbf{z}^3$	$\mathbf{1}$	$\nu_{i,R}$	$\mathbf{z}$	$-\mathbf{1}$
$l_{i,R}$	$\mathbf{z}$	$\mathbf{1}$	$\bar{N}_{i,L}$	$\mathbf{z}^3$	$\mathbf{1}$
$N_{i,R}$	$\mathbf{z}$	$\mathbf{1}$			
Φ	$\mathbf{1}$	$\mathbf{1}$	χ	$\mathbf{1}$	$-\mathbf{1}$
ζ	$\mathbf{z}$	$\mathbf{1}$	η	$\mathbf{z}^2$	$\mathbf{1}$

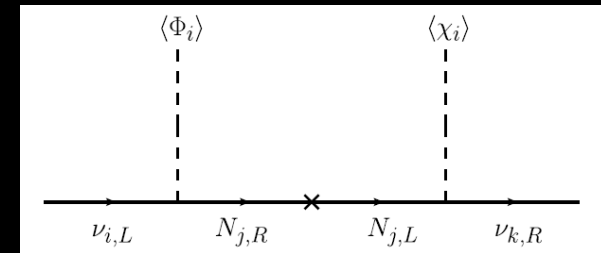
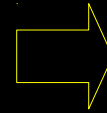


Lepton Quarticity vs Lepton number

Dark Matter Stability from Diracness

Chiulia et al arXiv:1606.04543
Phys.Lett. B761 (2016) 431

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Φ	$\mathbf{1}$	$\mathbf{1}$	χ	$\mathbf{1}$	$-\mathbf{1}$
ζ	$\mathbf{z}$	$\mathbf{1}$	η	$\mathbf{z}^2$	$\mathbf{1}$



Lepton Quarticity vs Lepton number

non SUSY scalar WIMP

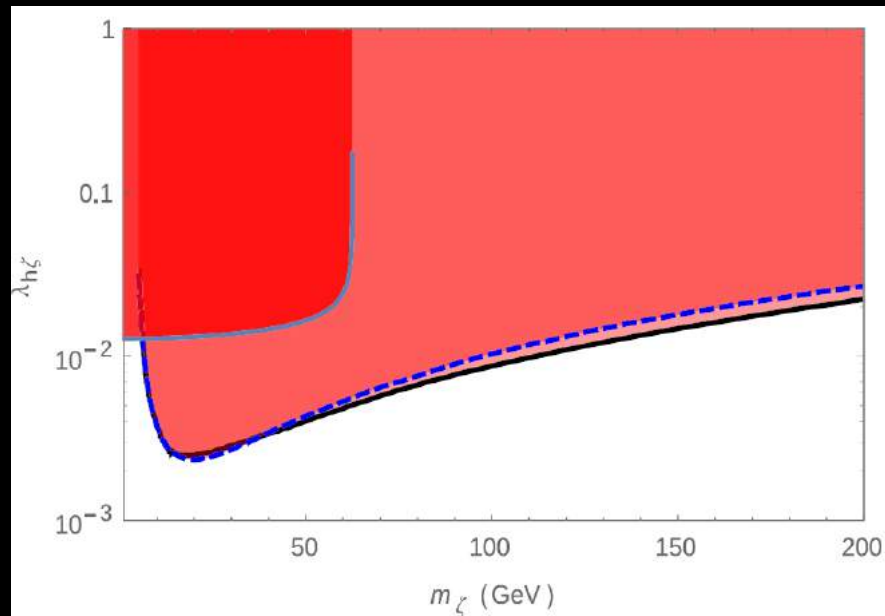
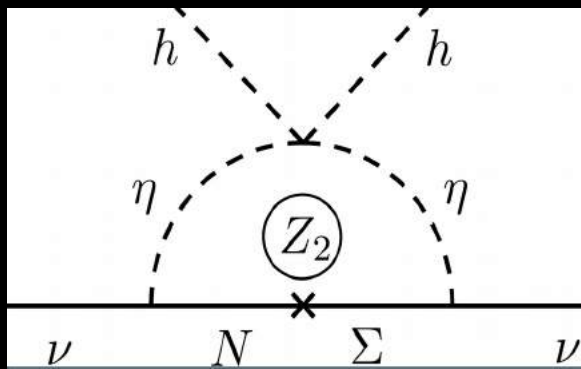


Fig. 5. The experimental sensitivity of our WIMP scalar dark matter candidate to invisible Higgs decay and direct detection. The light shaded region is ruled out by LUX (black continuous line) [34] and PandaX (blue dashed line) [35] data whereas the dark shaded region is ruled out by the bound on the Higgs invisible decay width from the LHC [36].

scotogenic dark matter

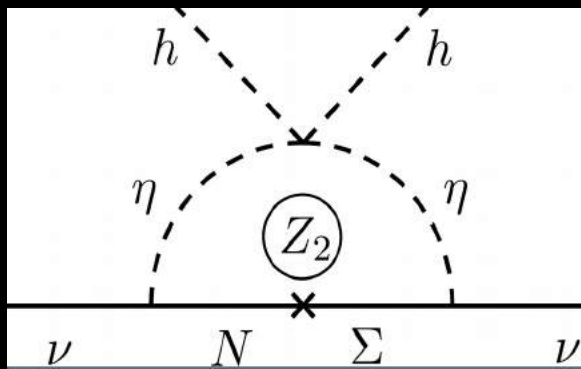


E Ma, Hirsch et al JHEP 1310 (2013) 149

	Standard Model			Fermions		Scalars	
	L	e	ϕ	Σ	N	η	Ω
Generations	3	3	1	1	1	1	1
$SU(2)_L$	2	1	2	3	1	2	3
$U(1)_Y$	-1/2	-1	1/2	0	0	1/2	0
Z_2	+	+	+	-	-	-	+

WIMP dark Matter as radiative
neutrino mass messenger

scotogenic dark matter



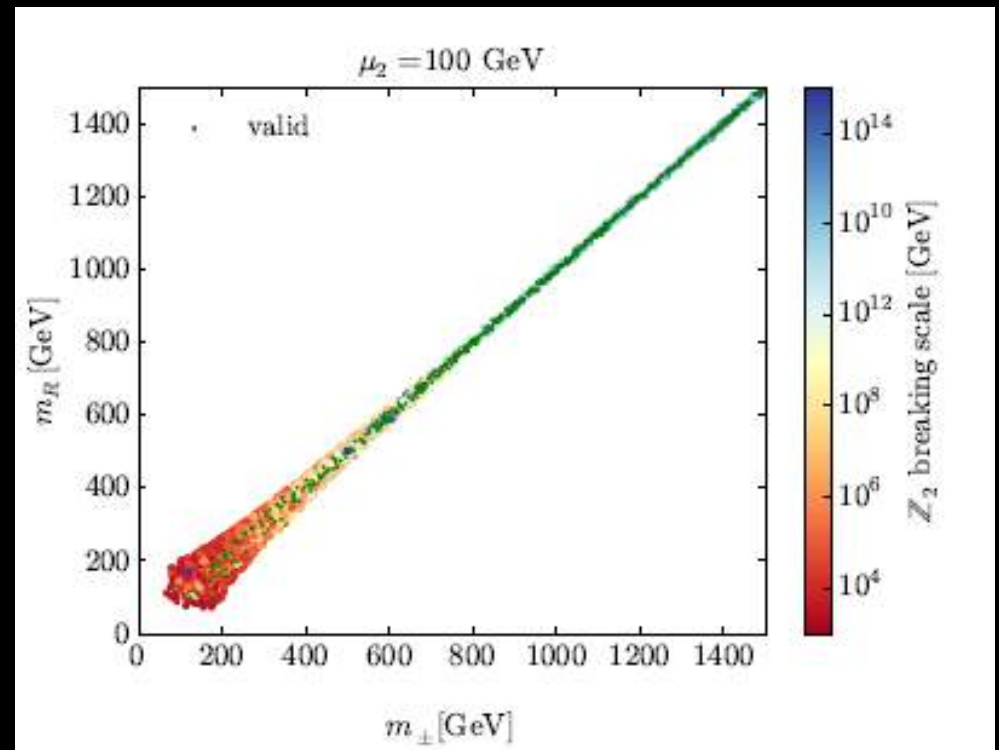
E Ma, Hirsch et al JHEP 1310 (2013) 149

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WIMP dark Matter as radiative neutrino mass messenger

Merle et al JHEP 1607 (2016) 013

Either scalar or fermion messenger
“susy” without susy



Scotogenic dark matter stability from Diracness

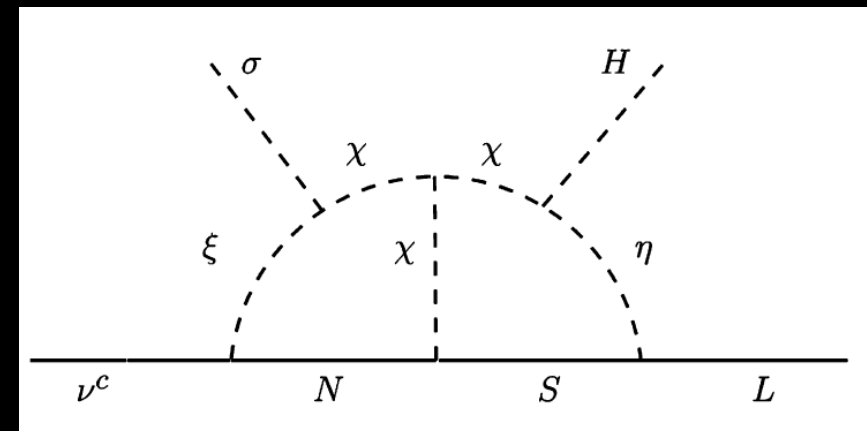
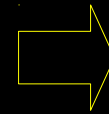
Lepton triality vs Lepton number

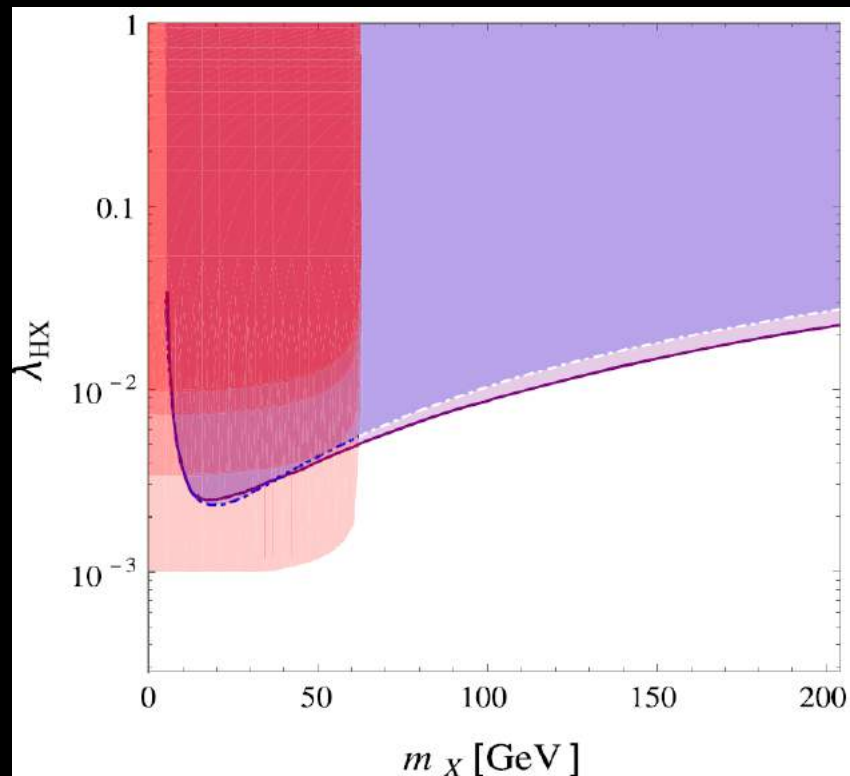
C. Bonilla et al. / Physics Letters B 762 (2016) 214–218

Table 1

Relevant particle content and quantum numbers of the model.

	$\bar{L}$	ν^c	H	η	N	S	σ	ξ	χ
$SU(2)_L$	2	1	2	2	1	1	1	1	1
$U(1)_D$	-1	3	0	0	-1	1	2	-2	0
Z_3^{DM}	1	1	1	α	α	α	1	α^2	α
Z_3	ω	ω^2	1	1	ω	ω^2	1	1	1





non SUSYscalar WIMP

Scotogenic dark matter stability from Diracness

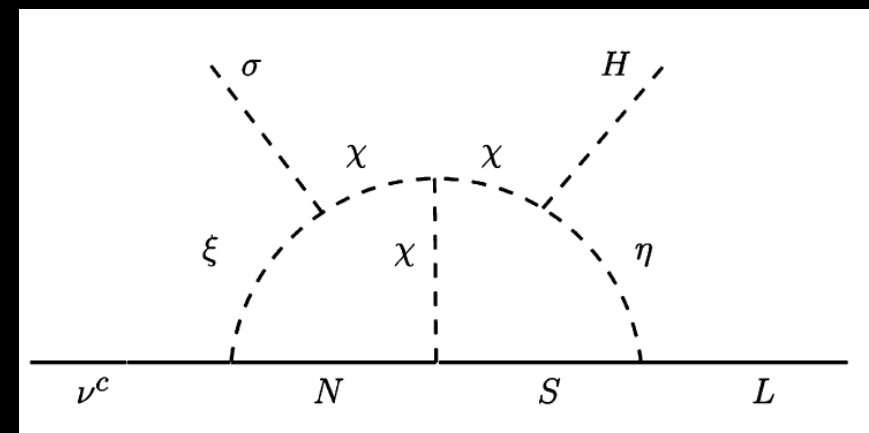
Lepton triality vs Lepton number

C. Bonilla et al. / Physics Letters B 762 (2016) 214–218

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Z_3	ω	ω^2	1	1	ω	ω^2	1	1	1



The keV majoron as a dark matter particle

V. Berezhinsky¹

*INFN, Laboratori Nazionali del Gran Sasso, I-67010, Assergi (AQ), Italy
and Institute for Nuclear Research, Moscow, Russia*

and

J.W.F. Valle²

*Instituto de Física Corpuscular - IFIC/CSIC, Dept. de Física Teòrica, Universitat de València,
46100 Burjassot, València, Spain*

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46100 Burjassot, València, Spain*

Decaying Warm Dark Matter and Neutrino Masses

M. Lattanzi^{1,*} and J. W. F. Valle^{2,†}

¹*Oxford Astrophysics, Denis Wilkinson Building, Keble Road, OX1 3RH, Oxford, United Kingdom*

²*Instituto de Física Corpuscular-C.S.I.C./Universitat de València Campus de Paterna, Apt 22085, E-46071 València, Spain*

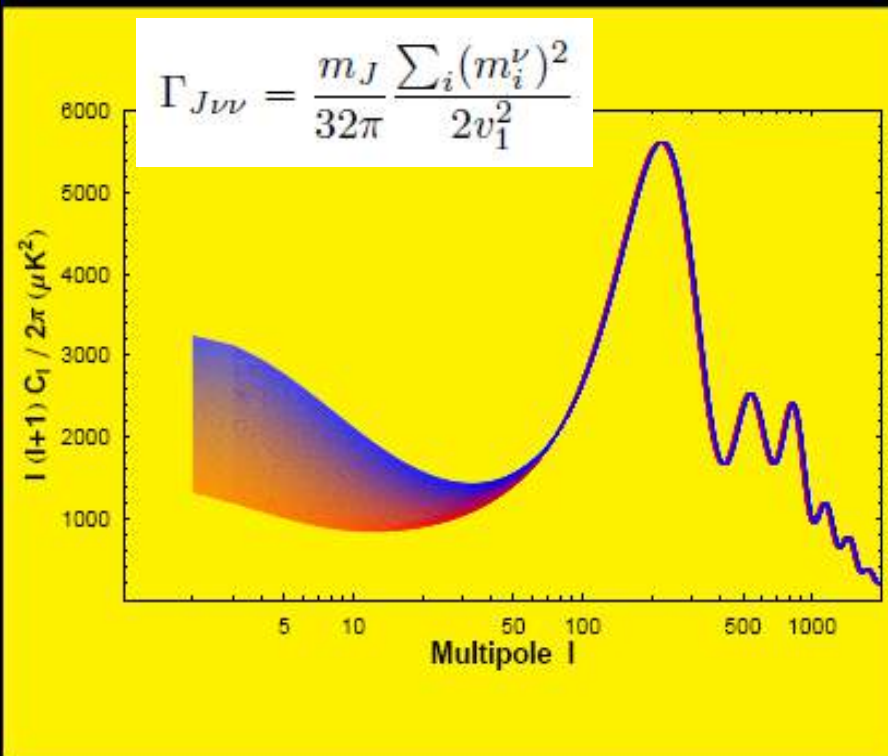
(Received 27 May 2007; published 20 September 2007)

Neutrino masses may arise from spontaneous breaking of ungauged lepton number. Because of quantum gravity effects the associated Goldstone boson—the majoron—will pick up a mass. We determine the lifetime and mass required by cosmic microwave background observations so that the massive majoron provides the observed dark matter of the Universe. The majoron decaying dark matter scenario fits nicely in models where neutrino masses arise via the seesaw mechanism, and may lead to other possible cosmological implications.

Lattanzi & Valle, PRL99 (2007) 121301

Consistency with CMB

Lattanzi & Valle, PRL99 (2007) 121301

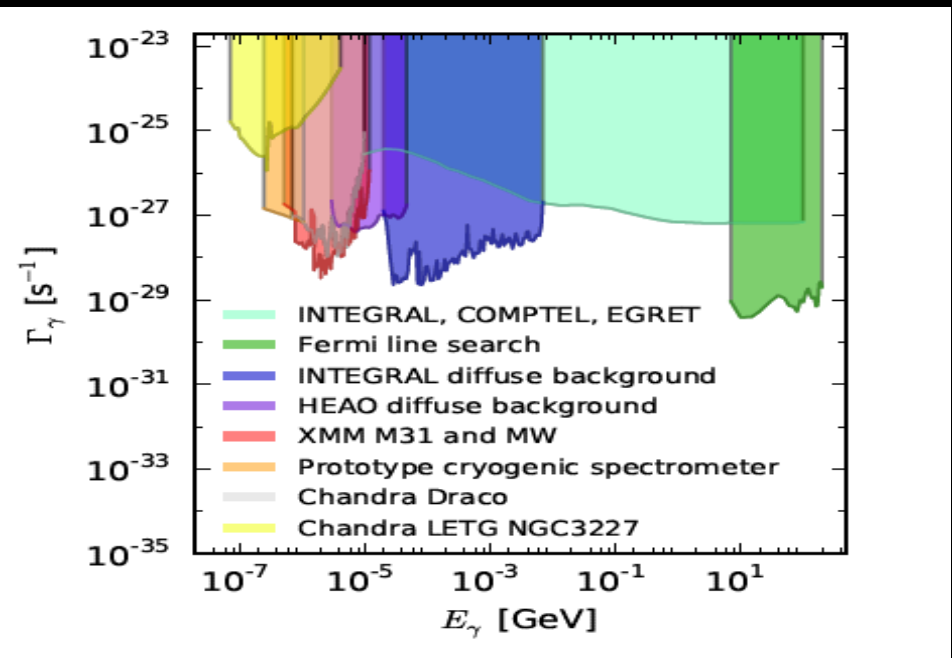


Bazzocchi & al JCAP 0808 (2008) 013

Esteves et al, PRD 82, 073008 (2010)

dark matter majorons

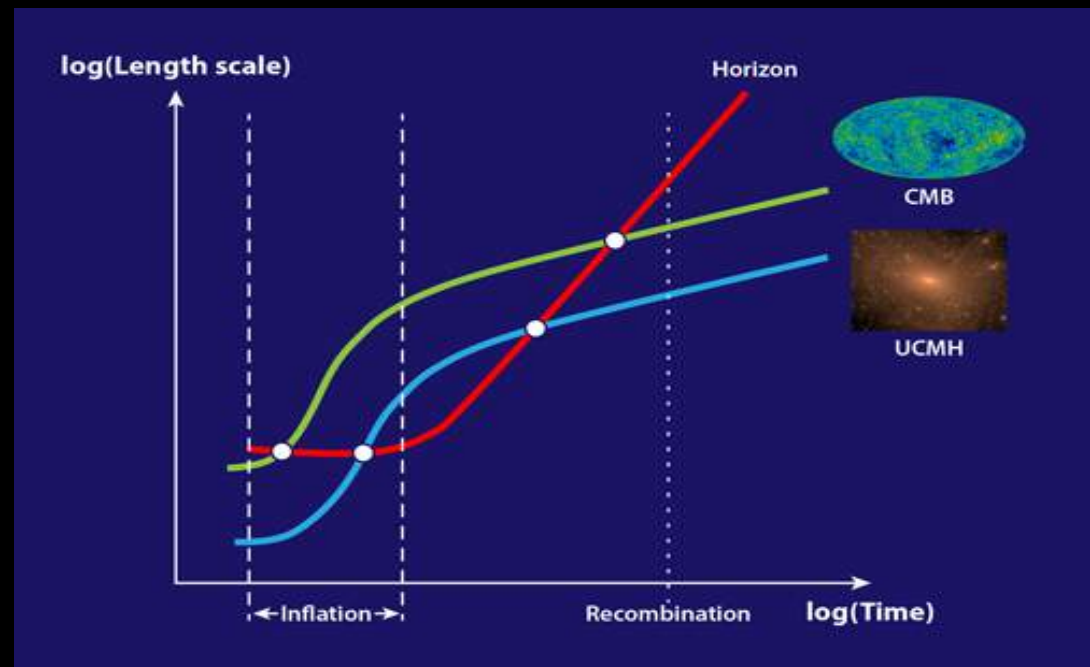
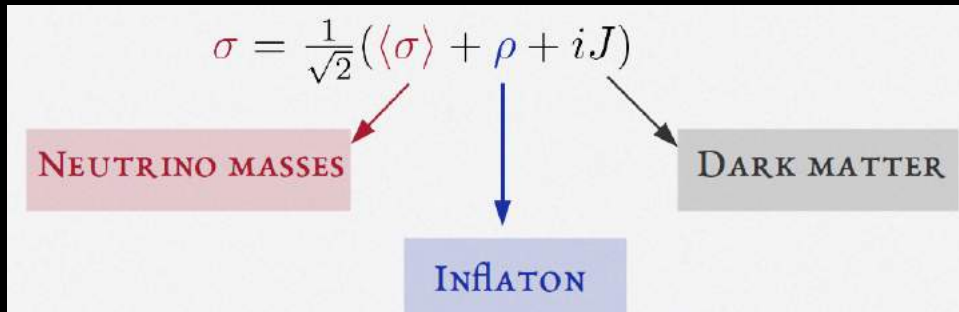
$J \rightarrow \gamma\gamma$



Lattanzi et al PRD88 (2013) 063528

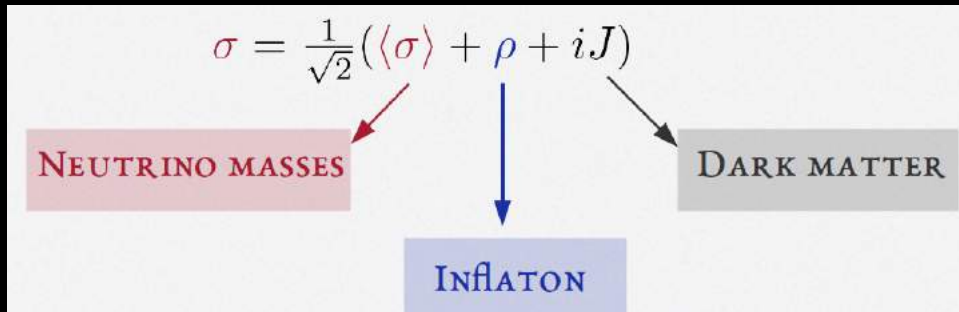
majoron dark matter & seesaw inflation

Boucenna, Morisi, Shafi, Valle
PRD90 (2014) 055023

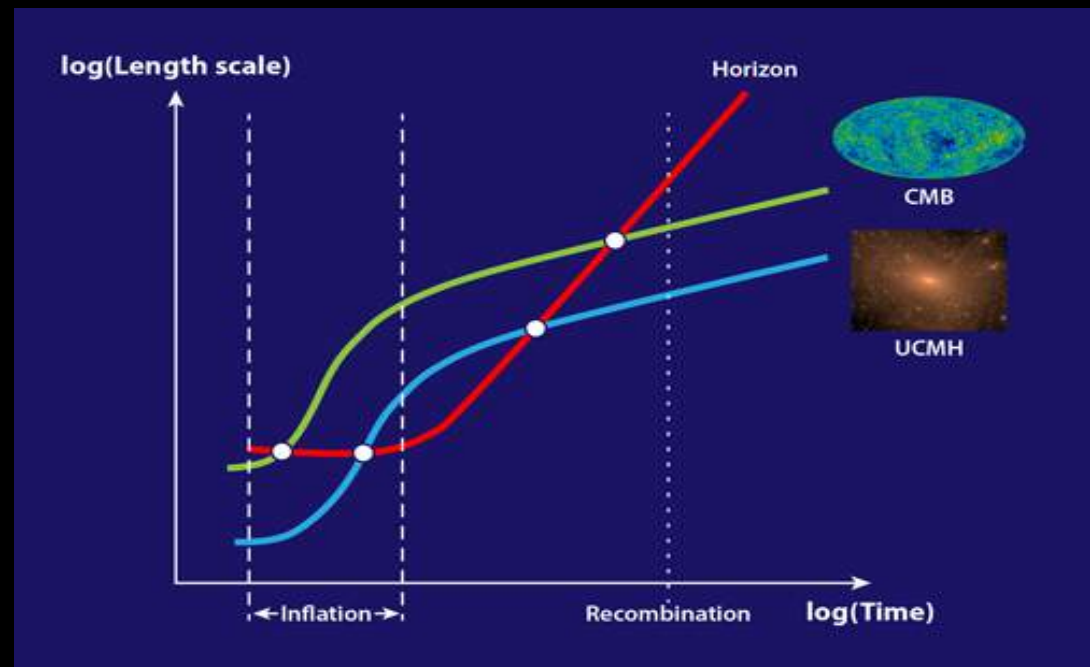
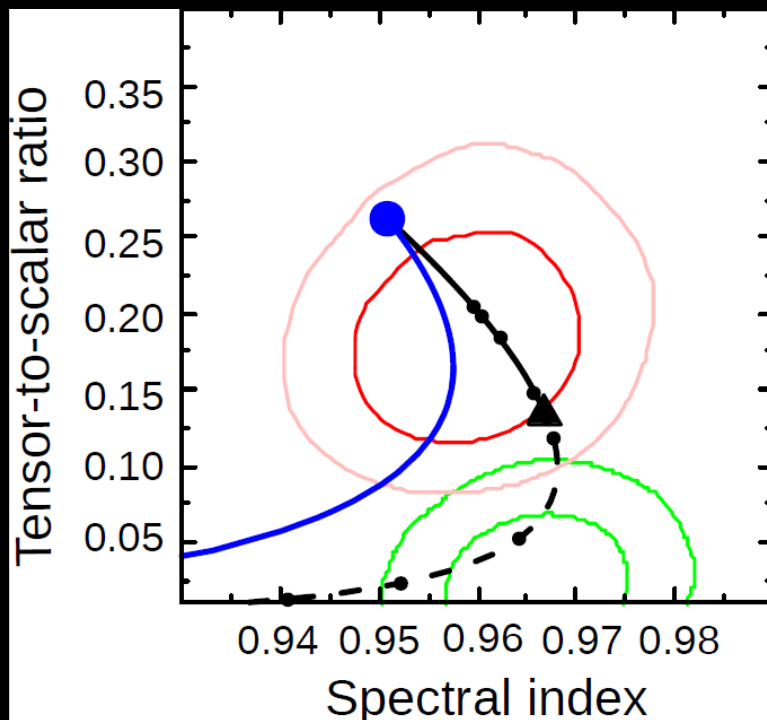


majoron dark matter & seesaw inflation

Boucenna, Morisi, Shafi, Valle
PRD90 (2014) 055023



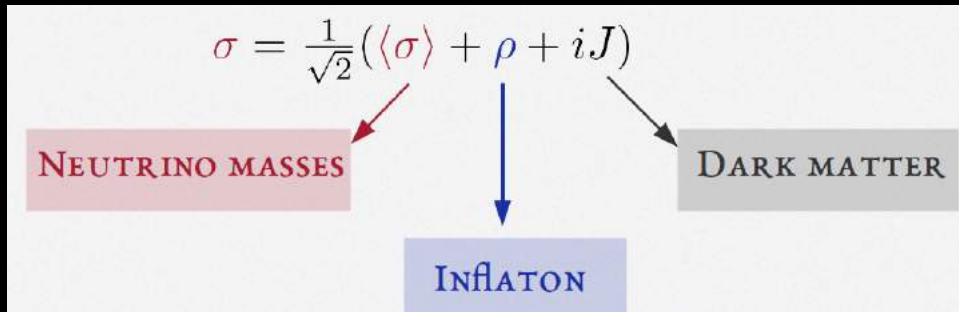
Quartic versus Higgs Inflation



<http://arxiv.org/pdf/1502.00612v1>

majoron dark matter & seesaw inflation

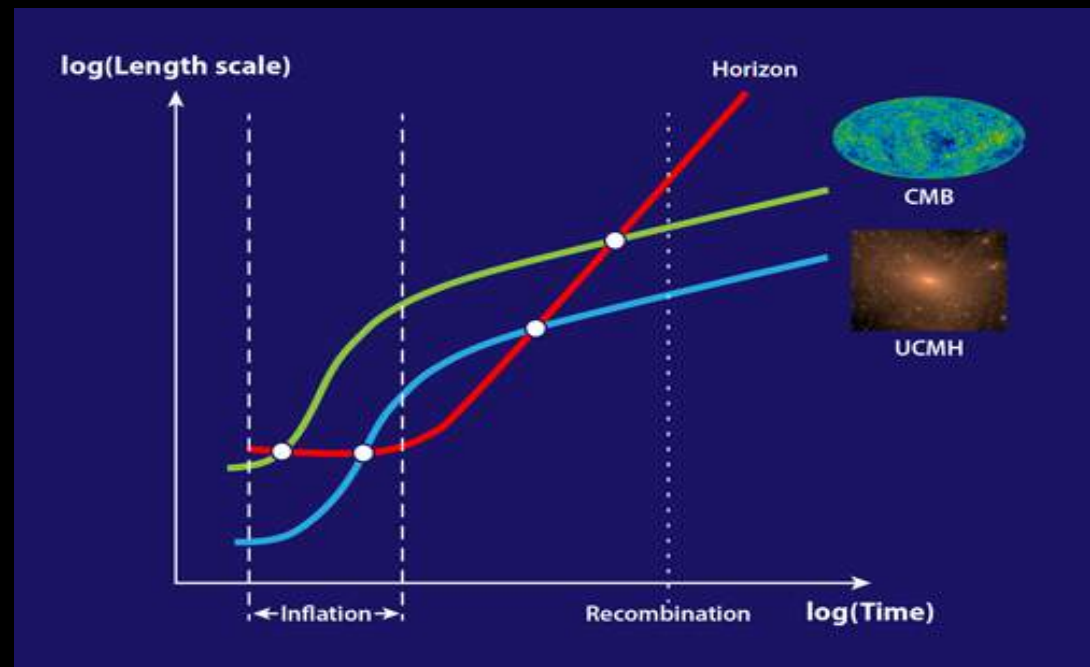
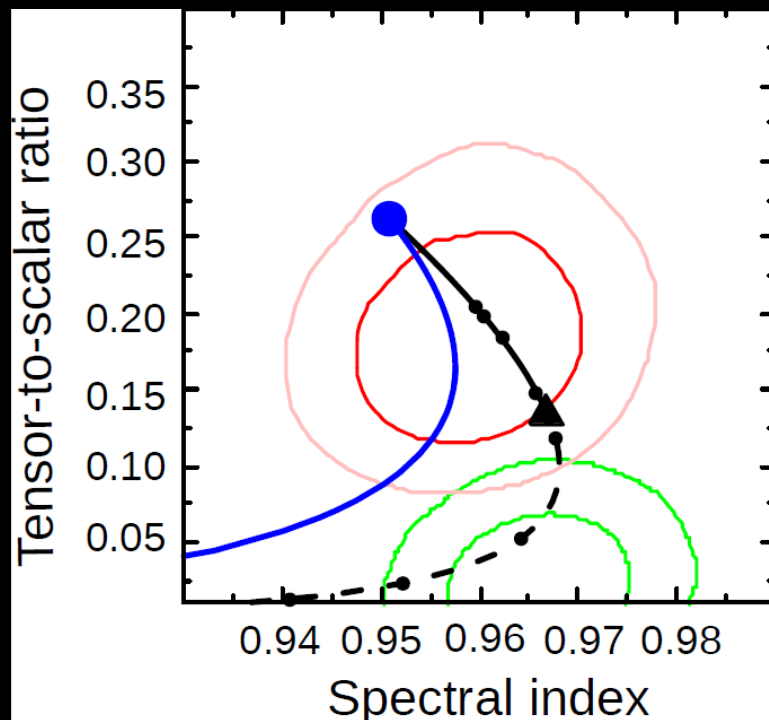
Boucenna, Morisi, Shafi, Valle
PRD90 (2014) 055023



type-I seesaw **Leptogenesis**

Aristizabal et al JCAP 1407 (2014) 052

Quartic *versus* Higgs Inflation



<http://arxiv.org/pdf/1502.00612v1>

DARK MATTER FROM FLAVOR SYMMETRY

- *Accidental ?*

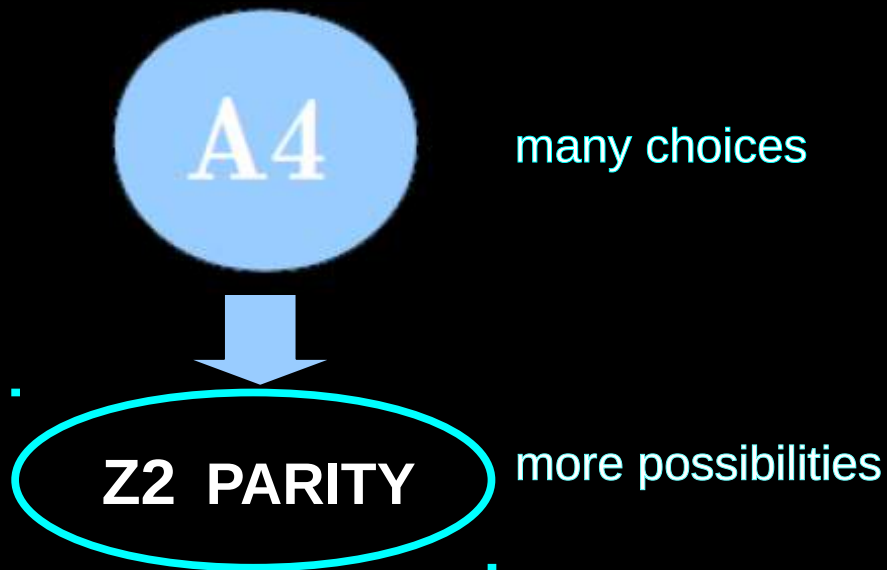
Lavoura, Morisi, JV JHEP 1302(2013) 118

DARK MATTER FROM FLAVOR SYMMETRY

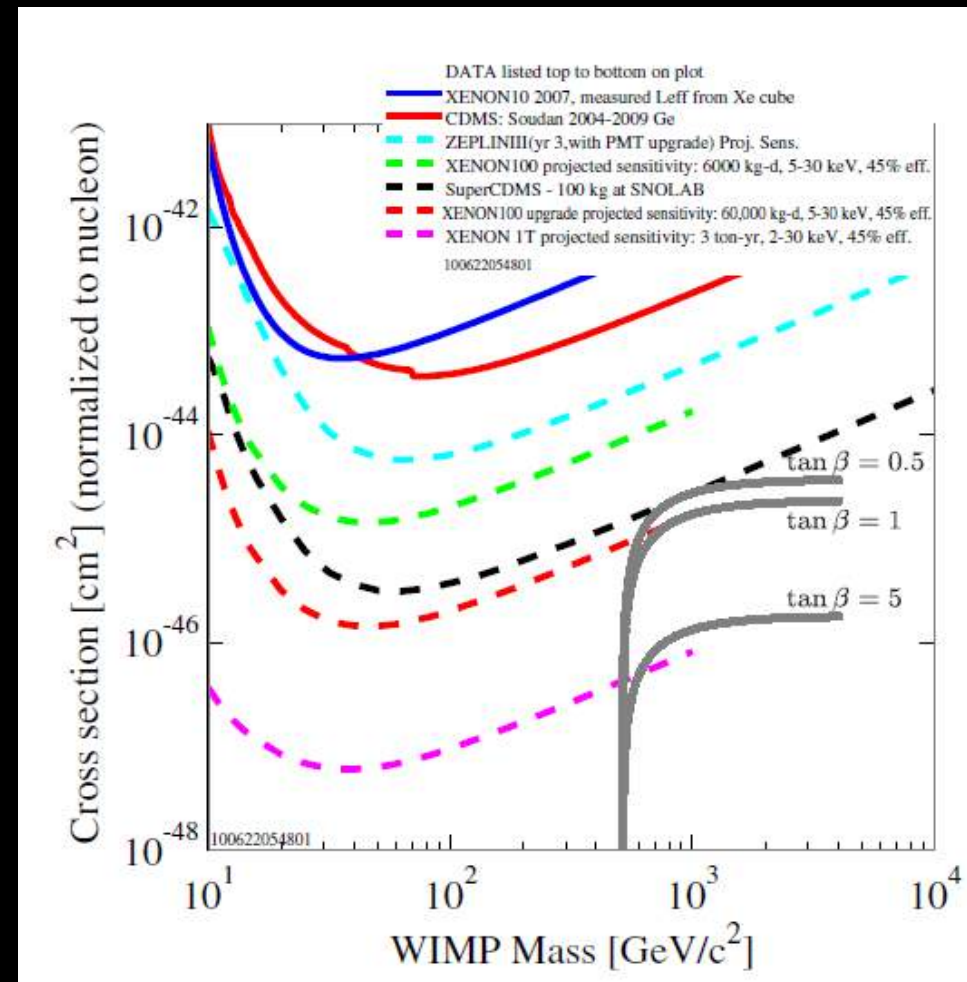
- *Accidental ?*
- *unbroken subgroup*

Lavoura, Morisi, JV JHEP 1302(2013) 118

Boucenna, et al JHEP 1105 (2011) 037
Hirsch, et al Phys.Rev. D82 (2010) 116003



**HIGGS PORTAL
DIRECT DETECTION**

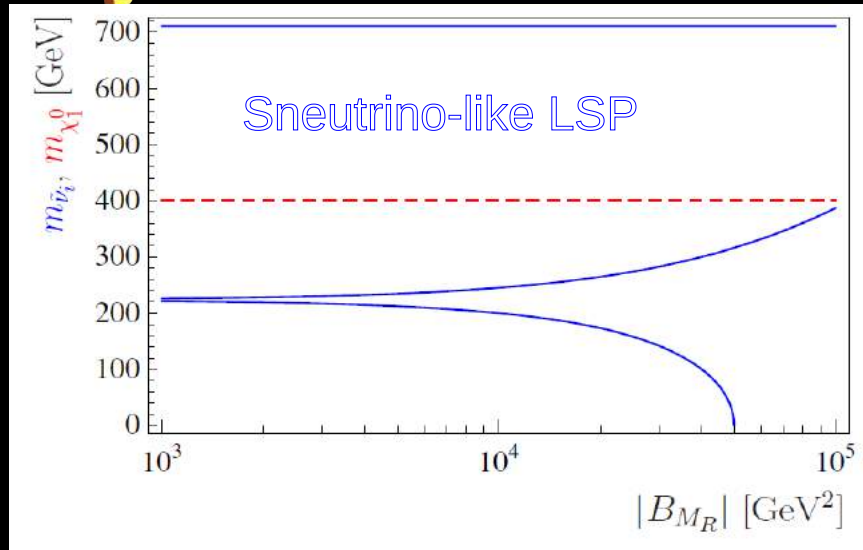


susy wimp dark matter

Arina et al PRL101 (2008) 161802

Bazzocchi, Cerdeno, Munoz, J.V., PRD81 (2010) 051701

De Romeri, Hirsch, JHEP 1212 (2012) 106



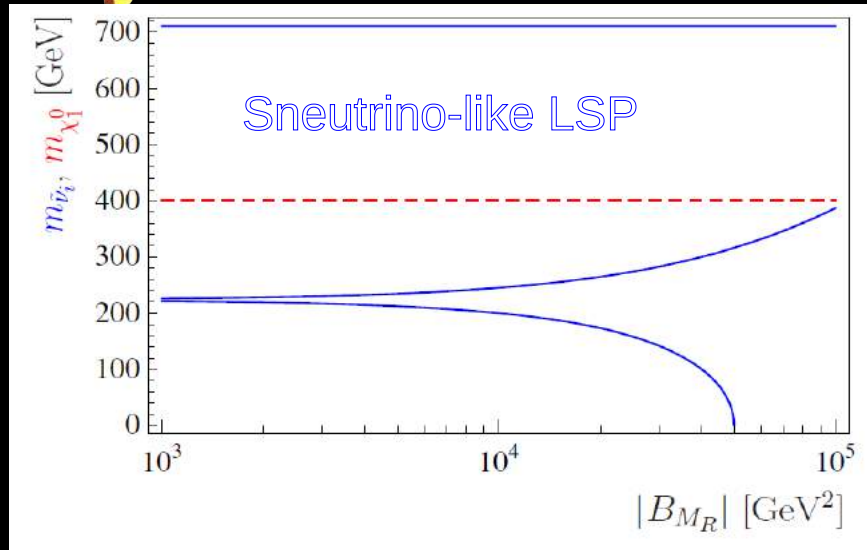
susy low-scale seesaw ...

susy wimp dark matter

Arina et al PRL101 (2008) 161802

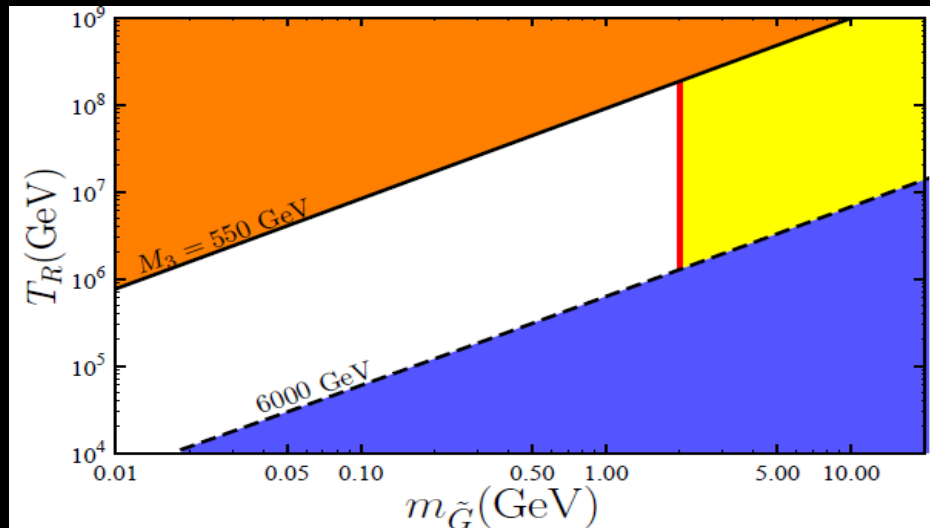
Bazzocchi, Cerdeno, Munoz, J.V., PRD81 (2010) 051701

De Romeri, Hirsch, JHEP 1212 (2012) 106



susy inverse seesaw ...

Restrepo et al PRD85 (2012) 023523



decaying Gravitino dark matter

doubly suppressed decays

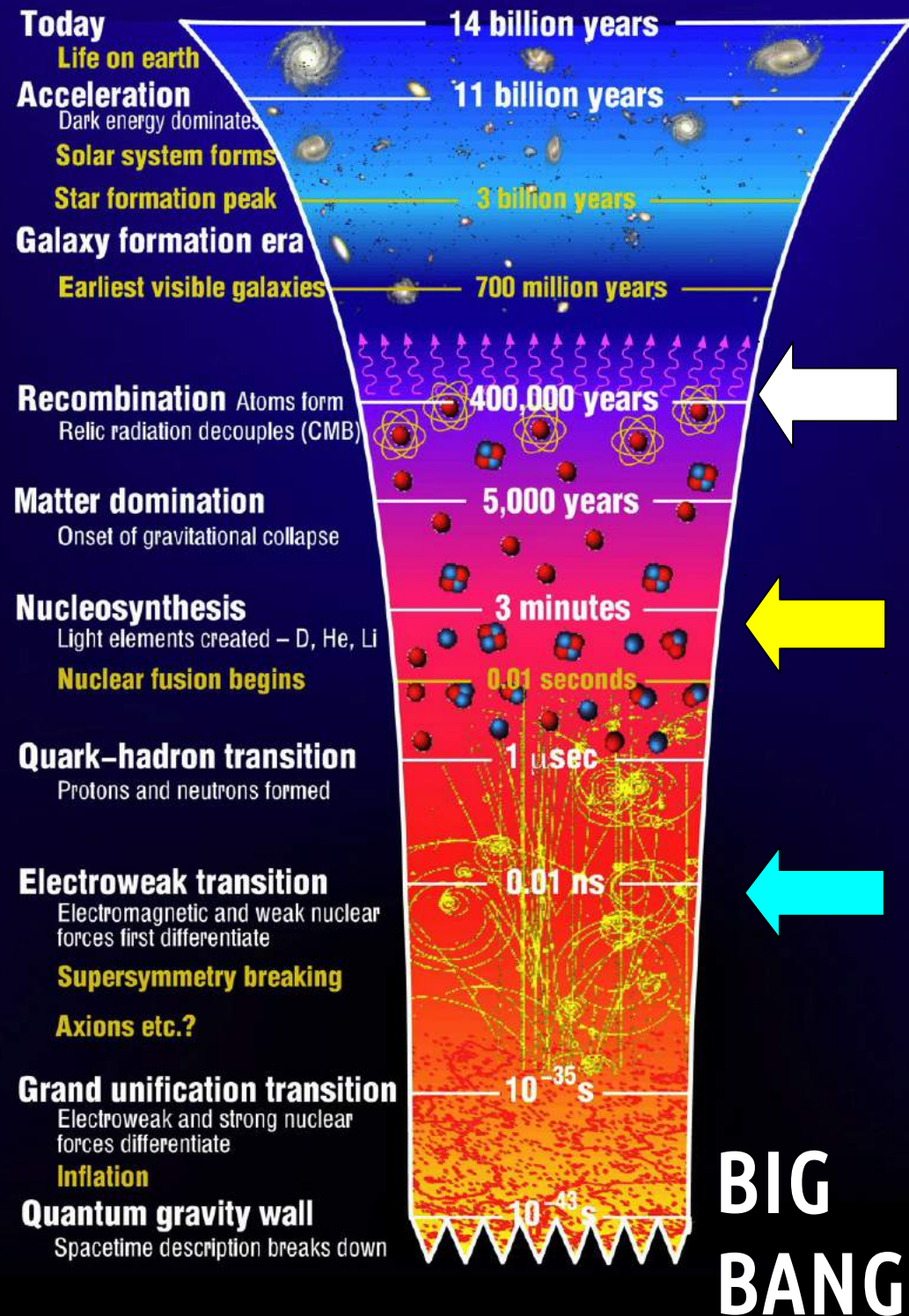
$$\Gamma = \Gamma(\tilde{G} \rightarrow \sum_i \nu_i \gamma) \simeq \frac{1}{32\pi} |U_{\tilde{\gamma}\nu}|^2 \frac{m_{\tilde{G}}^3}{M_P^2}$$

chosen to fit neutrino osc. data

conclusions

a most ubiquitous particle in the Universe

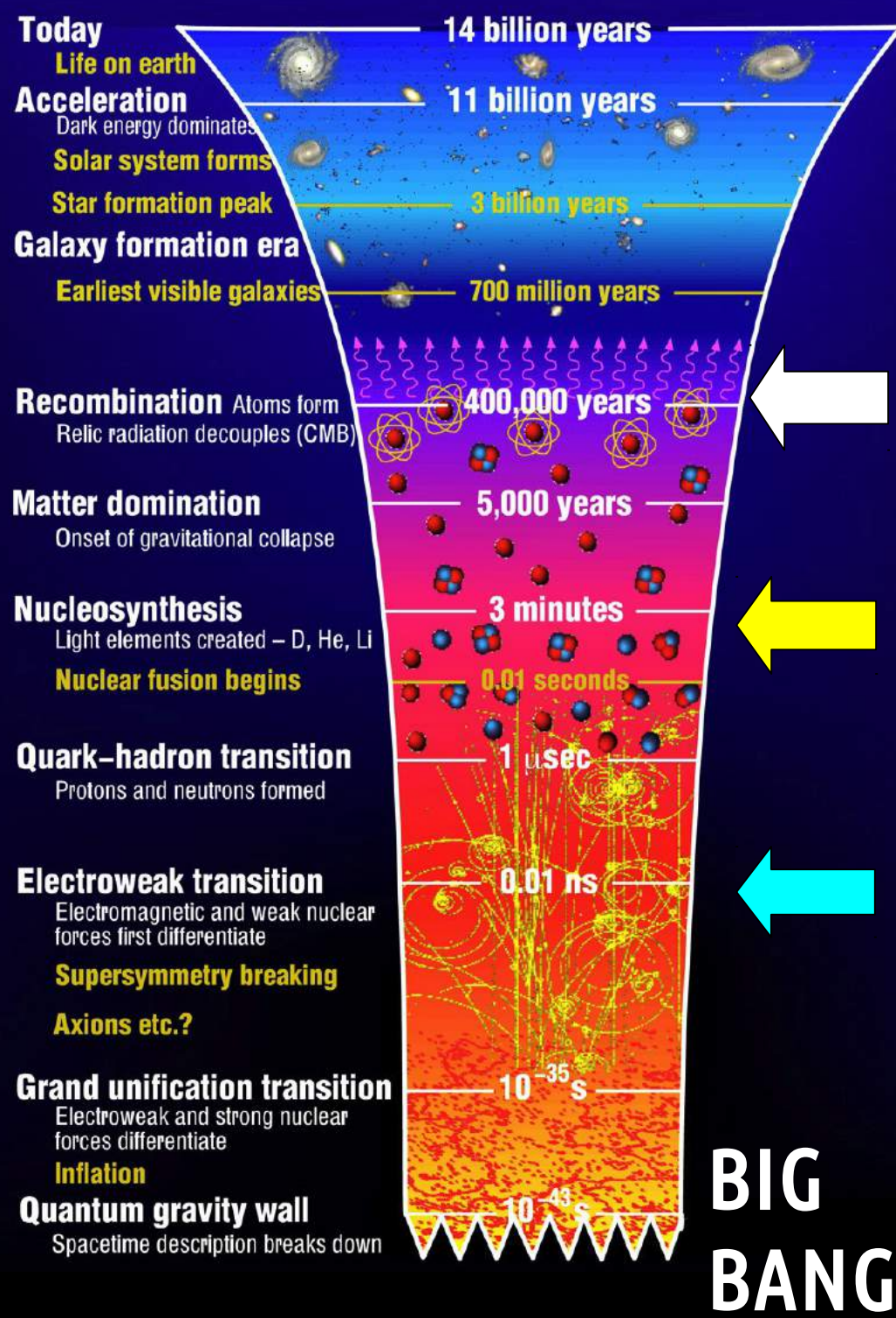
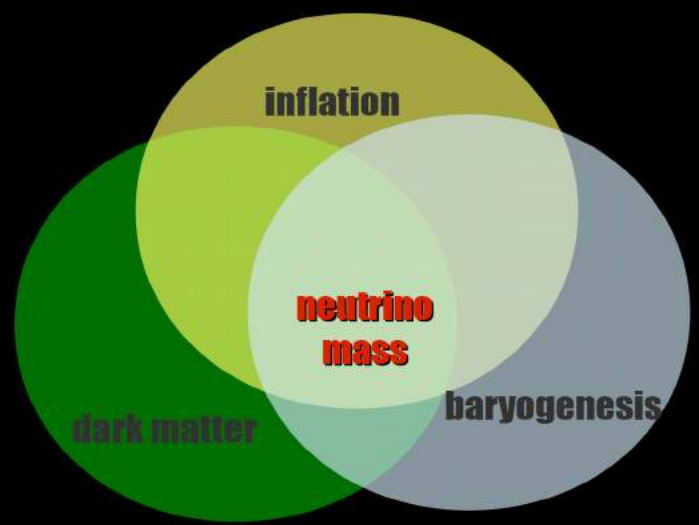
neutrinos may explain WDM or CDM
through an emergent theory ...



conclusions

a most ubiquitous particle in the Universe

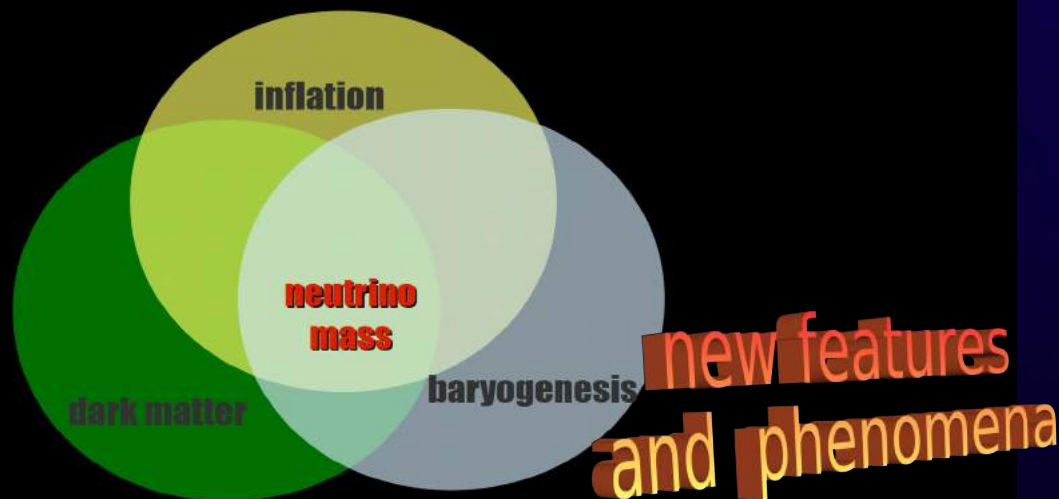
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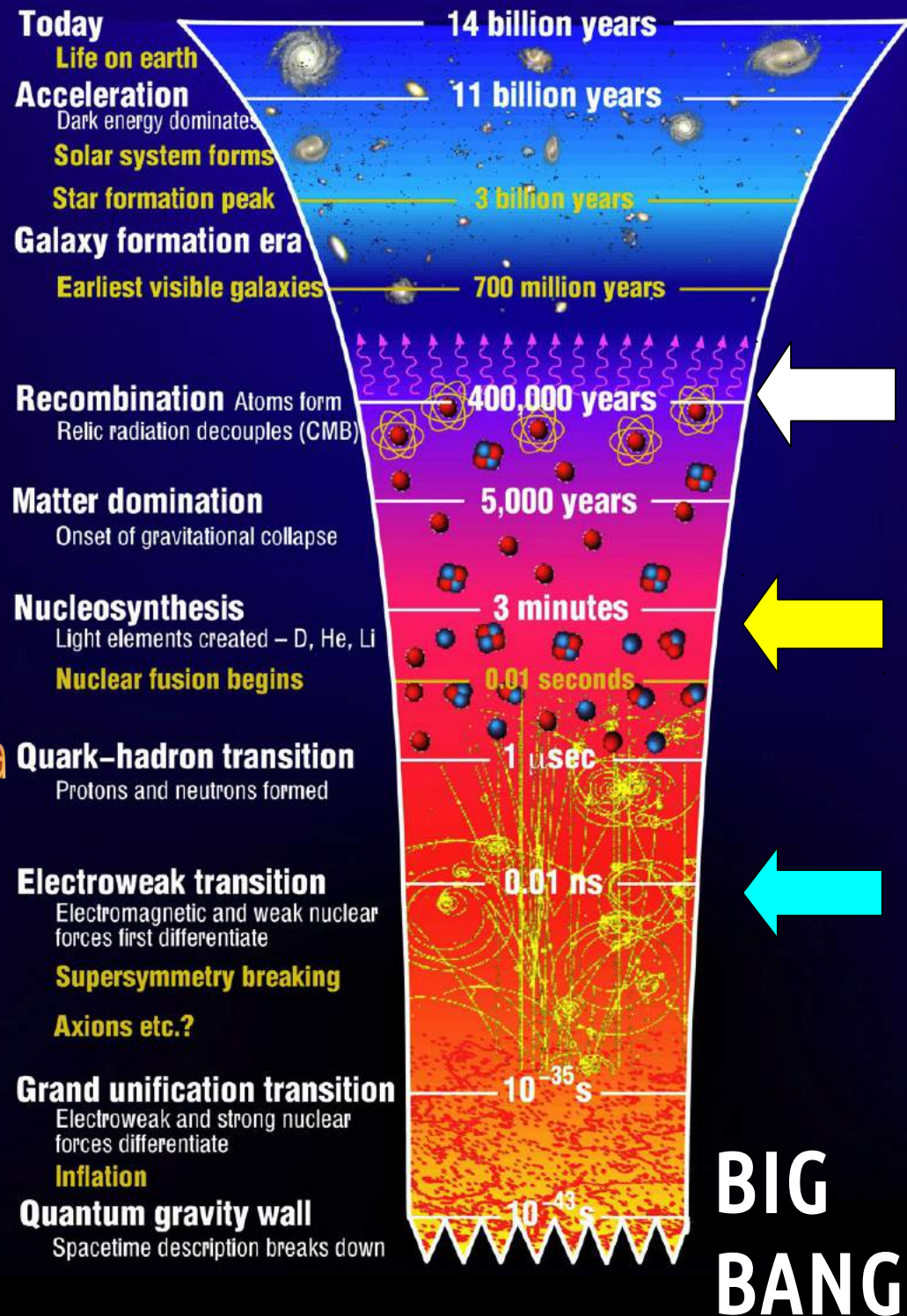
conclusions

a most ubiquitous particle in the Universe

neutrinos may explain WDM or CDM
through an emergent theory ...



- SSB consistency, new higgses, new decay modes
- unification without GUT embedding
- new gauge boson & fermion messengers @ LHC
- non-unitarity => new CPV in neutrino oscillations
- LFV mainly at HE colliders
- LFV/CPV unsuppressed by neutrino mass
- LNV either short or long-range



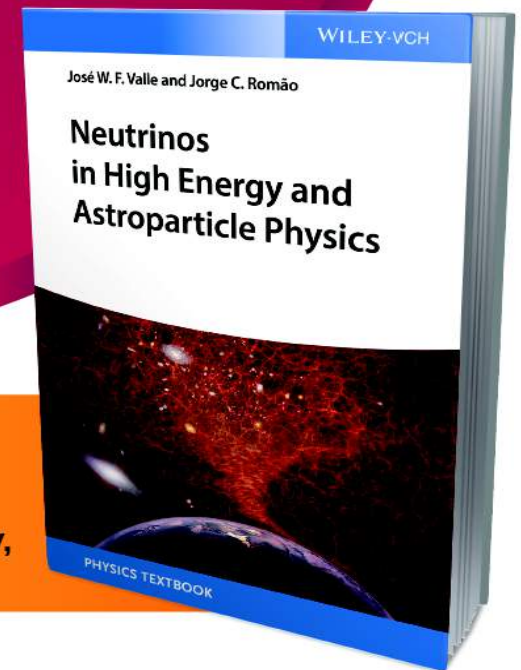
Thank you

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*Jose Wagner Furtado Valle,
Jorge Romao*

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Back up

Type I & non-unitarity of lepton mixing

$$\begin{pmatrix} \alpha_{11} & 0 & 0 \\ \alpha_{21} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

$$\alpha_{11}^2 \geq 0.989, \quad \alpha_{22}^2 \geq 0.999, \quad |\alpha_{21}|^2 \leq 6.6 \times 10^{-4}$$

Schechter & JV PRD22 (1980) 2227 & PDG
Rodejohann, JV Phys.Rev. D84 (2011) 073011

PLB199, 432 (1987)

PhysRevD.92.053009

non-unitary propagation hints associated
(relatively low-mass) type-I seesaw
messenger responsible for inducing neutrino mass

Type I & non-unitarity of lepton mixing

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PLB199, 432 (1987)

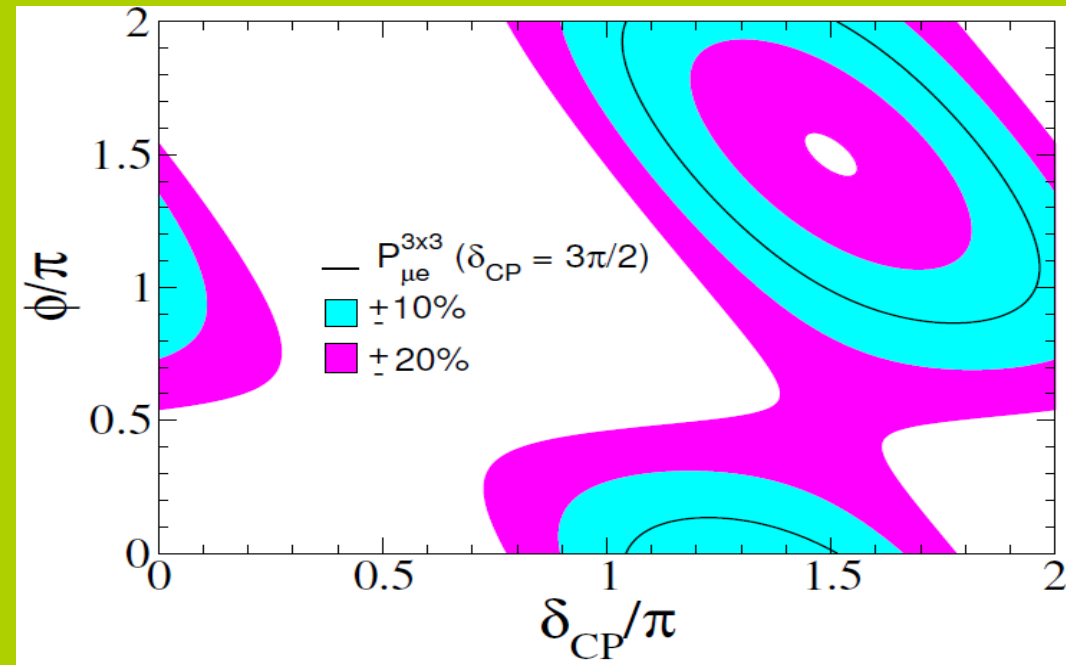
PhysRevD.92.053009

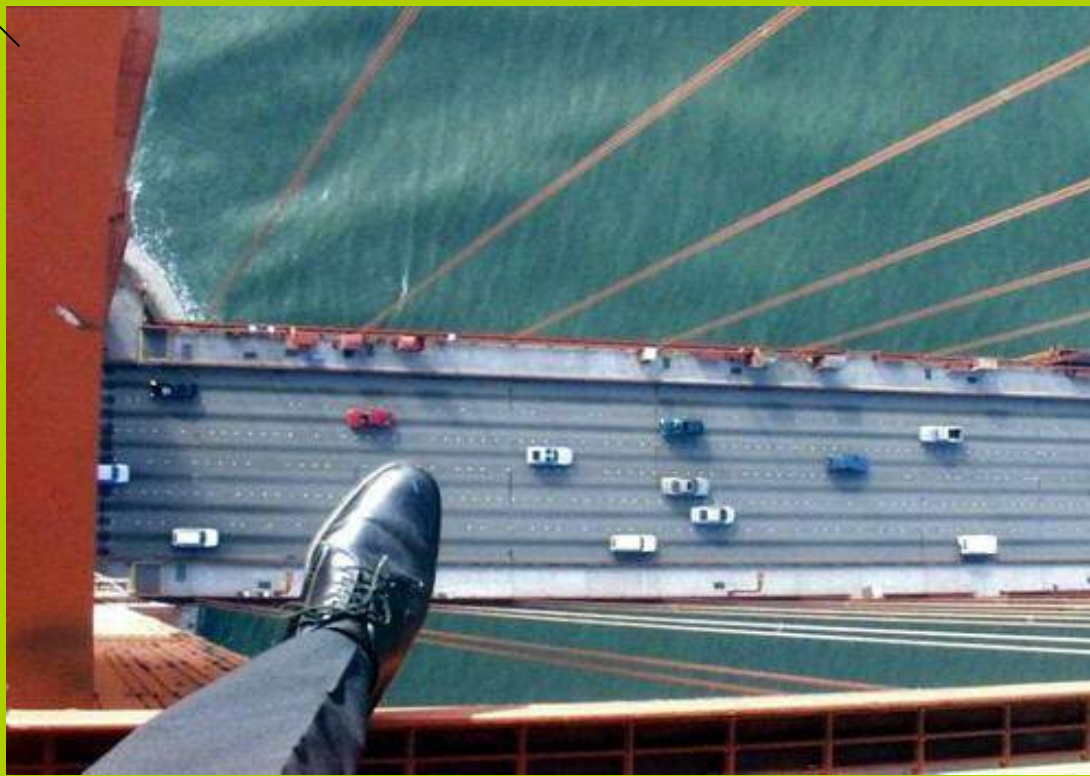
non-unitary propagation hints associated
(relatively low-mass) type-I seesaw
messenger responsible for inducing neutrino mass

$$P_{\mu e} = \alpha_{11}^2 \alpha_{22}^2 P_{\mu e}^{3 \times 3} + \alpha_{11}^2 \alpha_{22} |\alpha_{21}| P_{\mu e}^I + \alpha_{11}^2 |\alpha_{21}|^2$$

PhysRevLett.117.061804

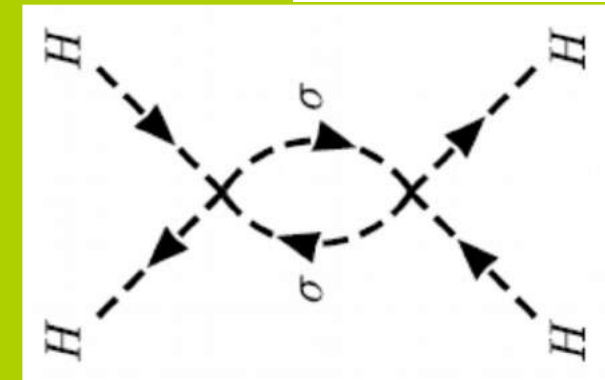
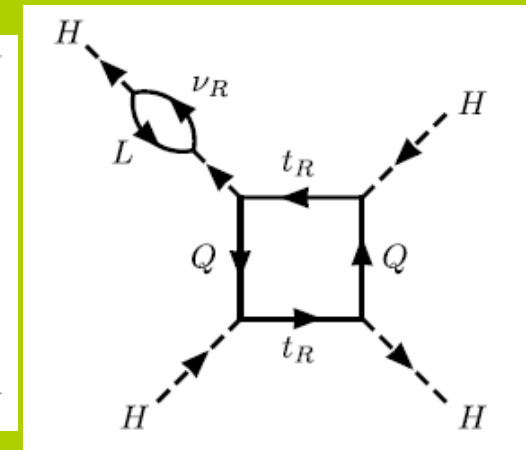
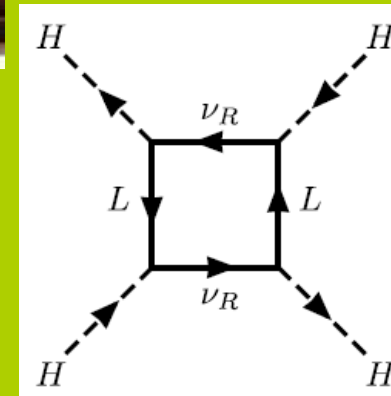
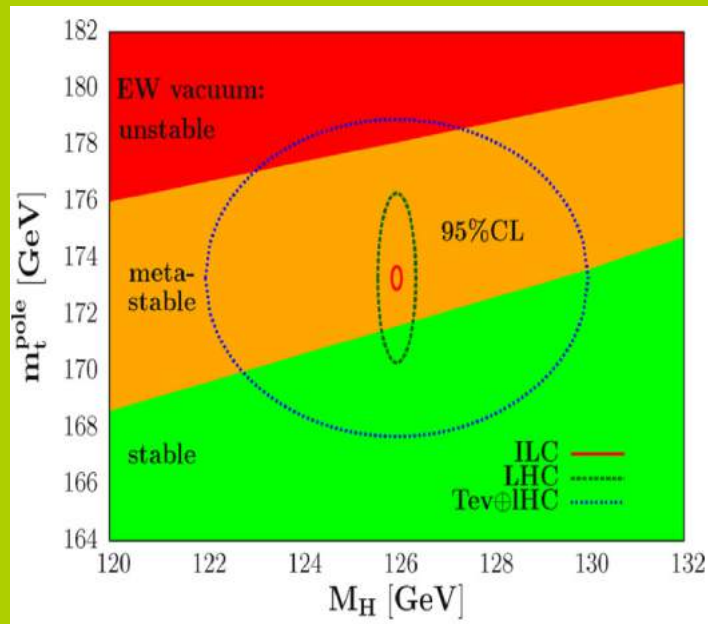
CP confusion





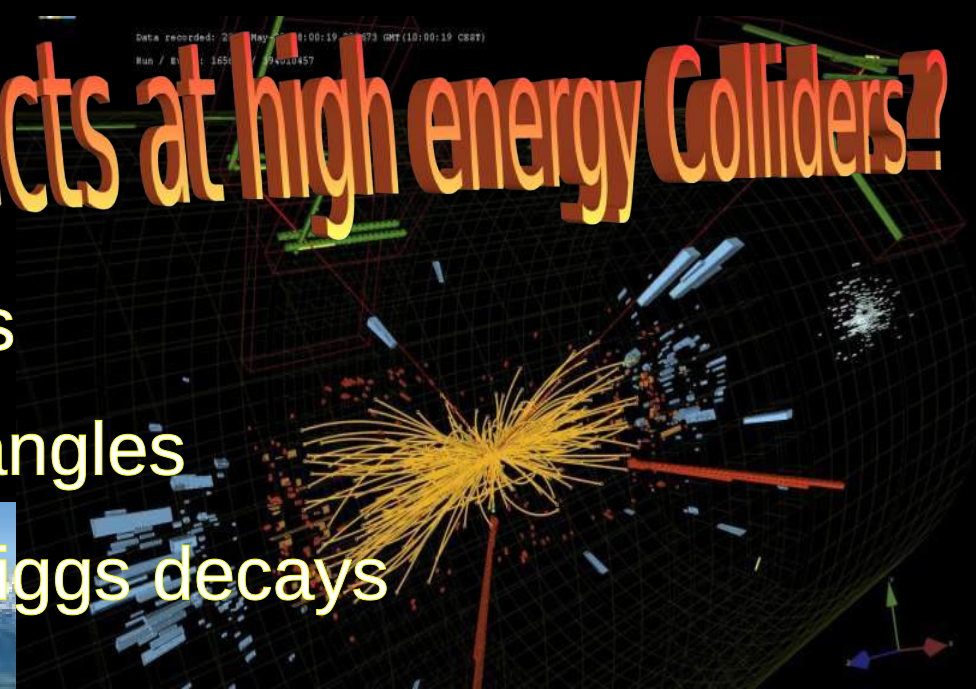
SM vacuum and neutrinos

Physics Letters B 756 (2016) 345–349



Neutrino effects at high energy Colliders?

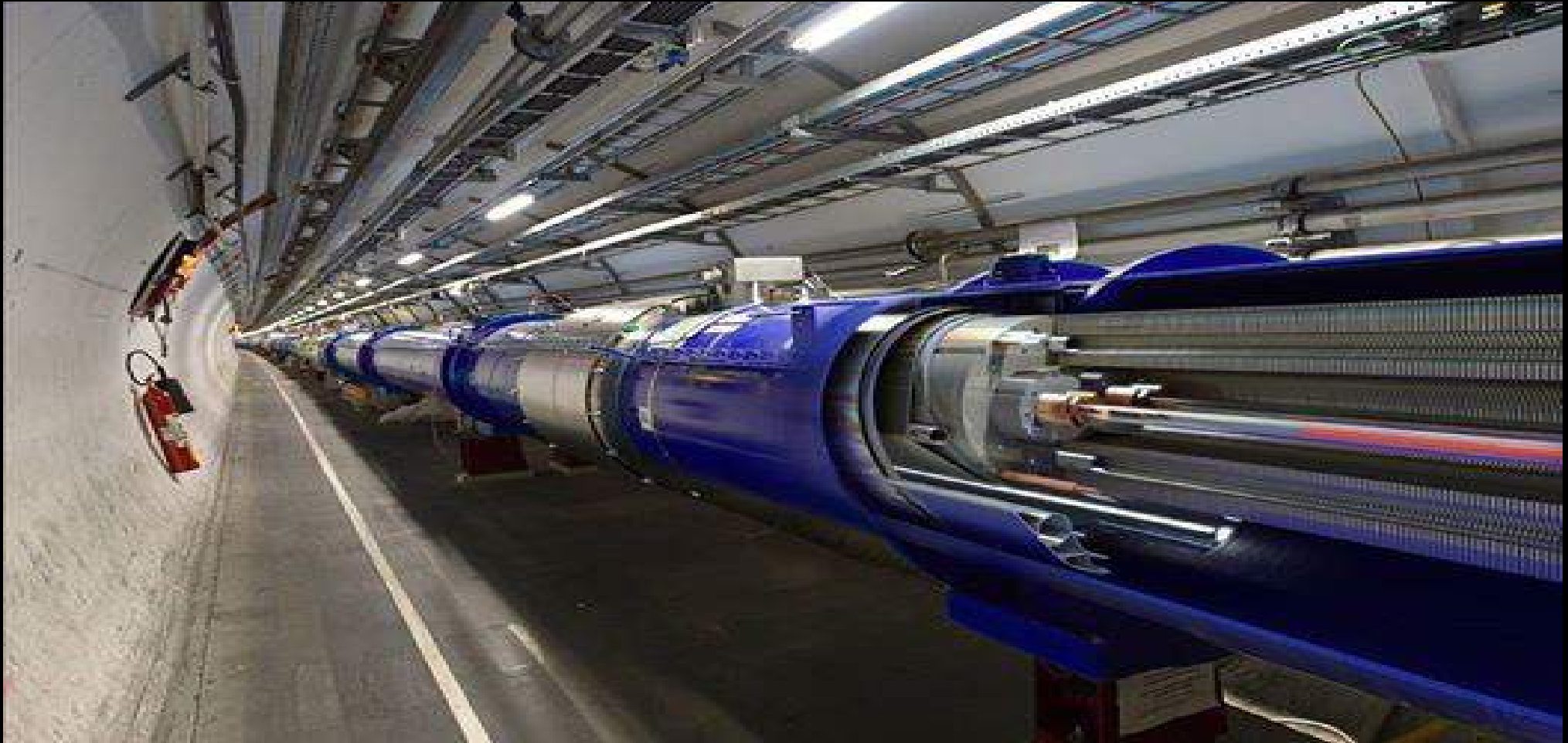
- Discover neutrino messengers
- Re-measure neutrino mixing angles
- Probe e-weak breaking e.g. Higgs decays



Invisible Higgs decays

Joshipura & J.V.

Nucl.Phys. B397 (1993) 105-122



Higgs searches 2016

Bonilla Fonseca & J.V.

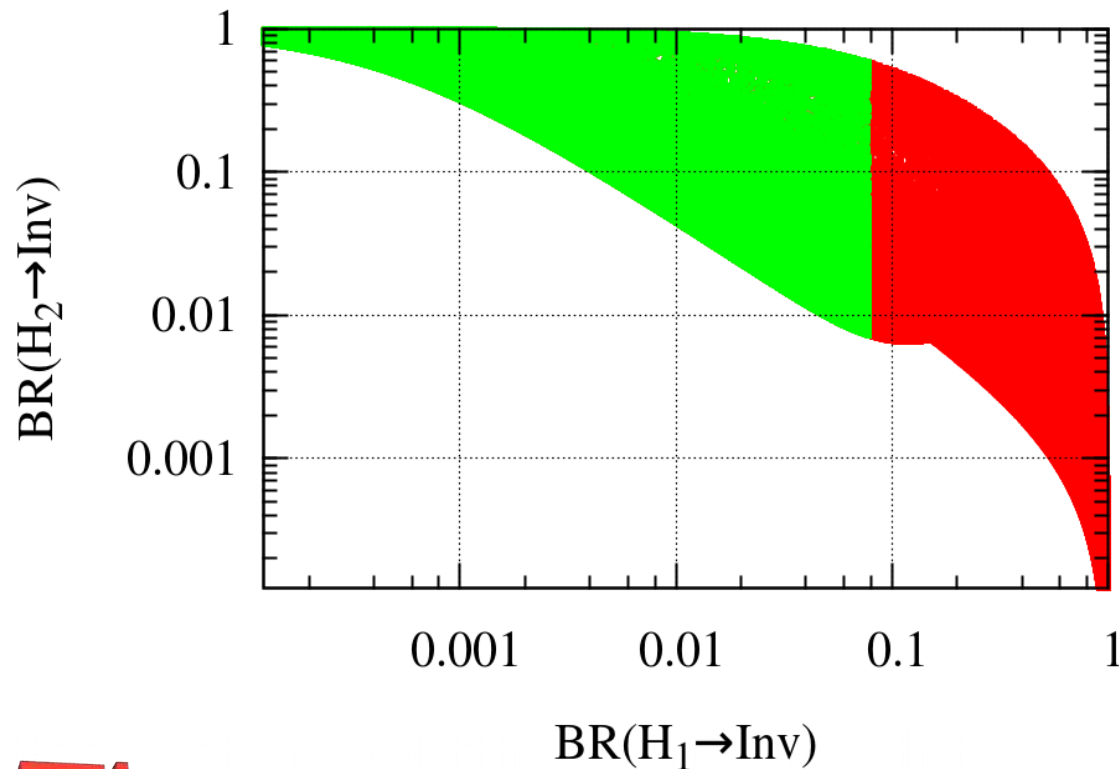
Phys.Lett. B756 (2016) 345-349 ...

Neutrino mass and invisible Higgs decays at the LHC

Cesar Bonilla,^{1,*} Jorge C. Romão,^{2,†} and José W.F. Valle^{1,‡}

$v_\sigma = 3 \text{ TeV}$

channel	ATLAS	CMS
$\mu_{\gamma\gamma}$	1.17 ± 0.27	$1.14^{+0.26}_{-0.23}$
μ_{WW}	$1.00^{+0.32}_{-0.29}$	0.83 ± 0.21
μ_{ZZ}	$1.44^{+0.40}_{-0.35}$	1.00 ± 0.29
$\mu_{\tau^+\tau^-}$	$1.4^{+0.5}_{-0.4}$	0.91 ± 0.27
$\mu_{b\bar{b}}$	$0.2^{+0.7}_{-0.6}$	0.93 ± 0.49

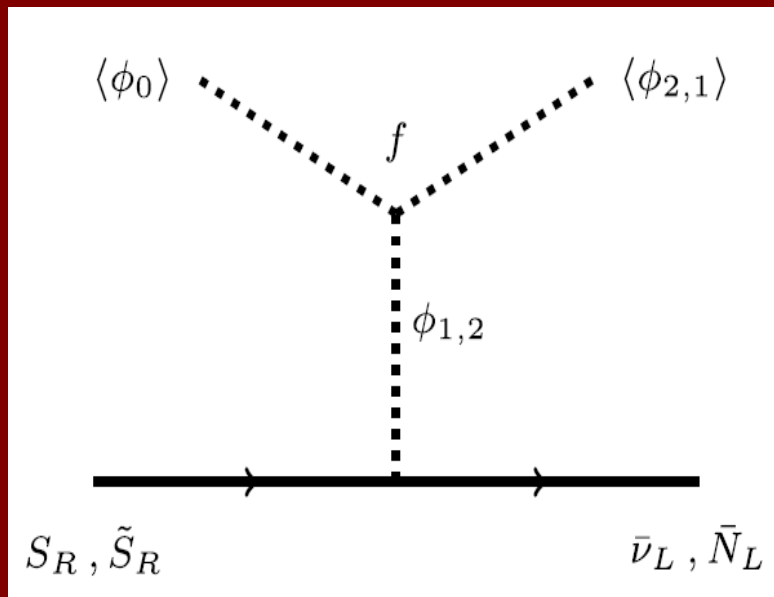


$$H_i \rightarrow JJ \quad \text{and} \quad H_2 \rightarrow 2H_1 \rightarrow 4J$$

$$\left(\text{when } m_{H_1} < \frac{m_{H_2}}{2} \right).$$

Theories of neutrino
as attractive higgs
search benchmarks

Dirac seesaw



Physics Letters B 755 (2016) 363–366

No conventional GUT embedding :

<http://arxiv.org/abs/arXiv:1608.05334>

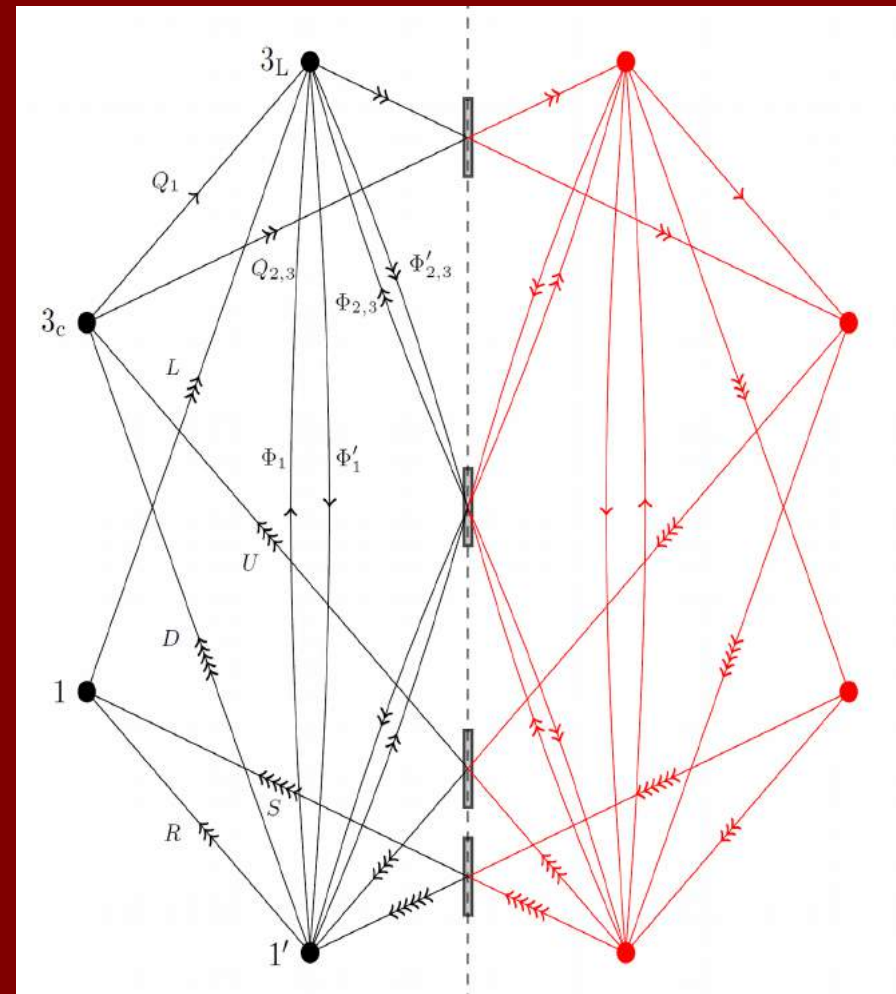
string completion Quiver setup

L and B conserved : no proton decay, no RPV ...

Addazi et al arXiv:1604.02117

331 from strings

10.1016/j.physletb.2016.06.015

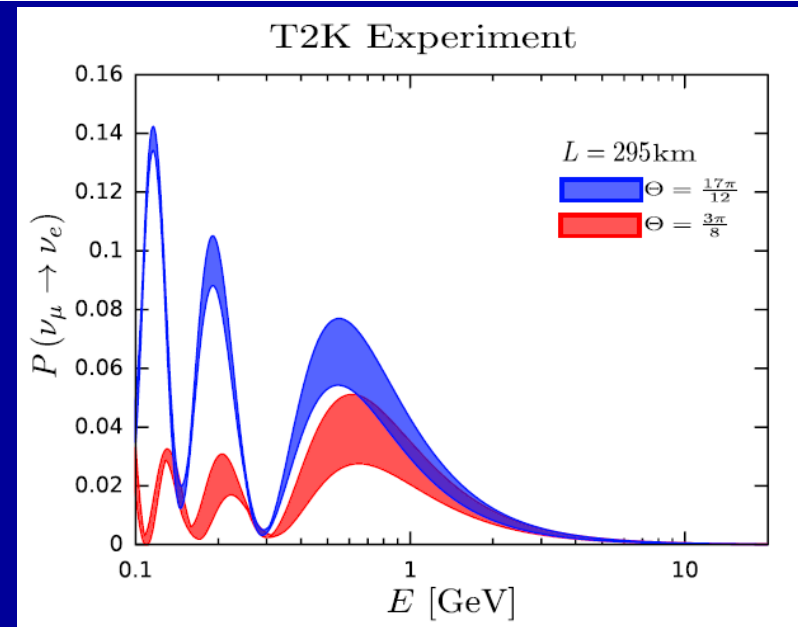
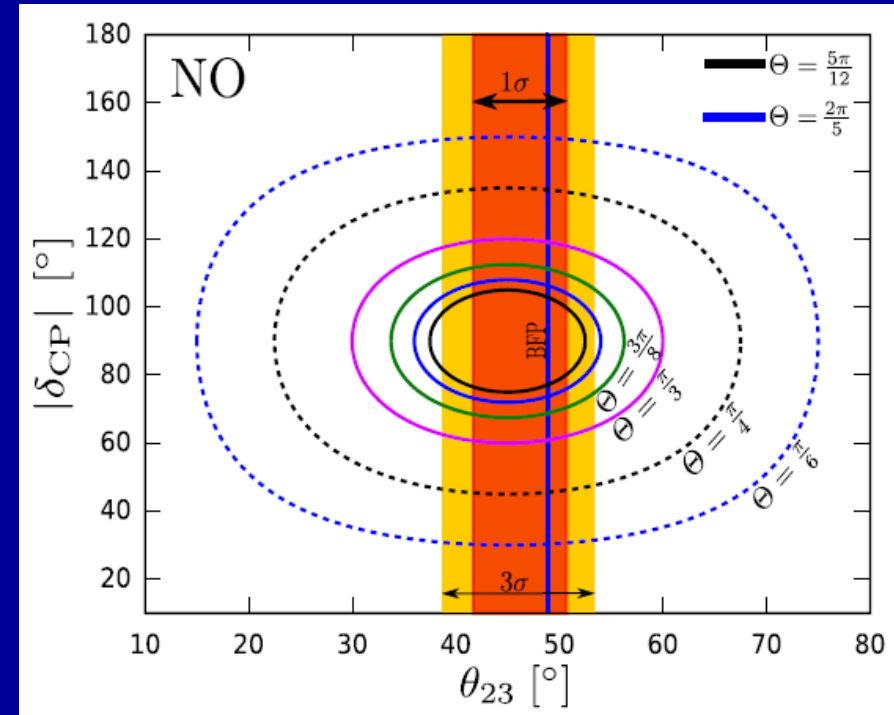
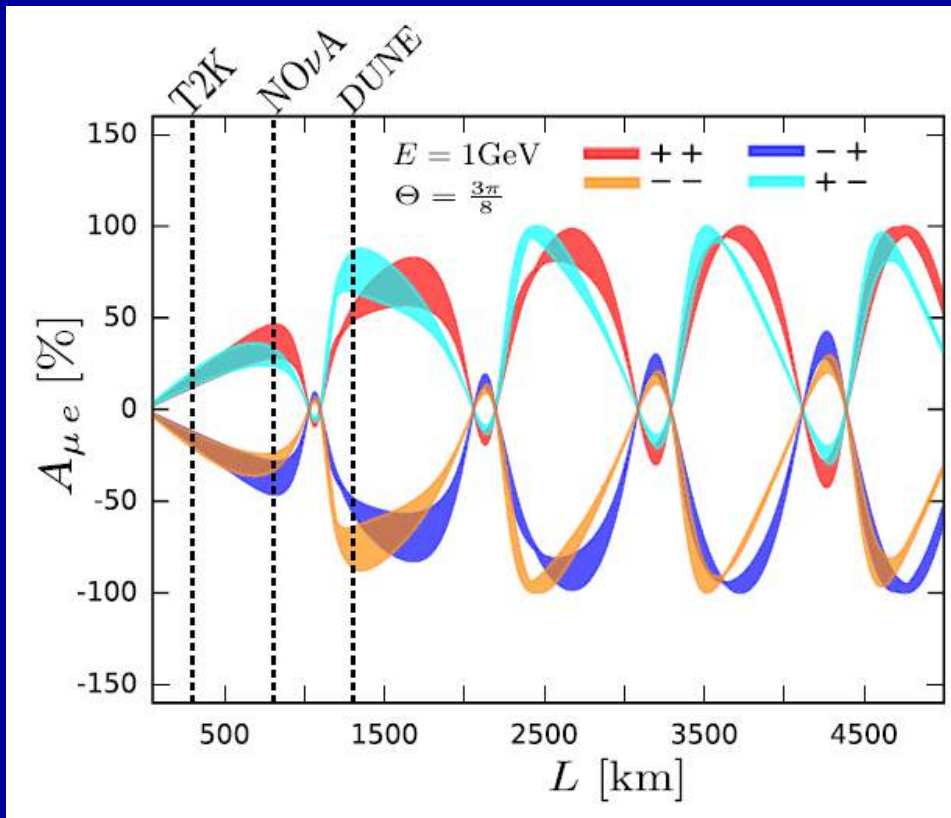


neutron-antineutron oscillations from exotic instantons

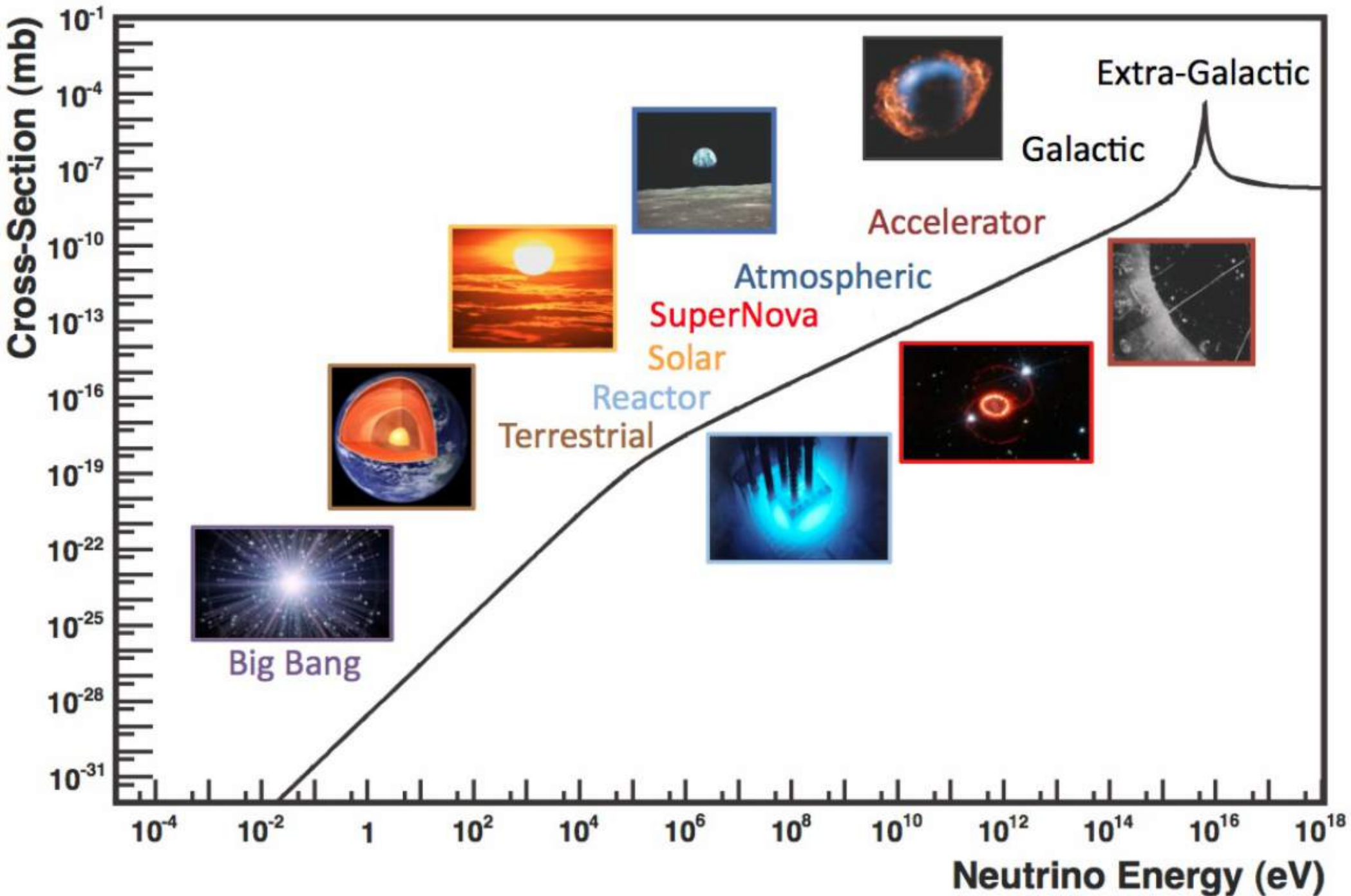
Model-independent flavor approach

$$\mathbf{X}^T \mathbf{m}_\nu \mathbf{X} = \mathbf{m}_\nu^*$$

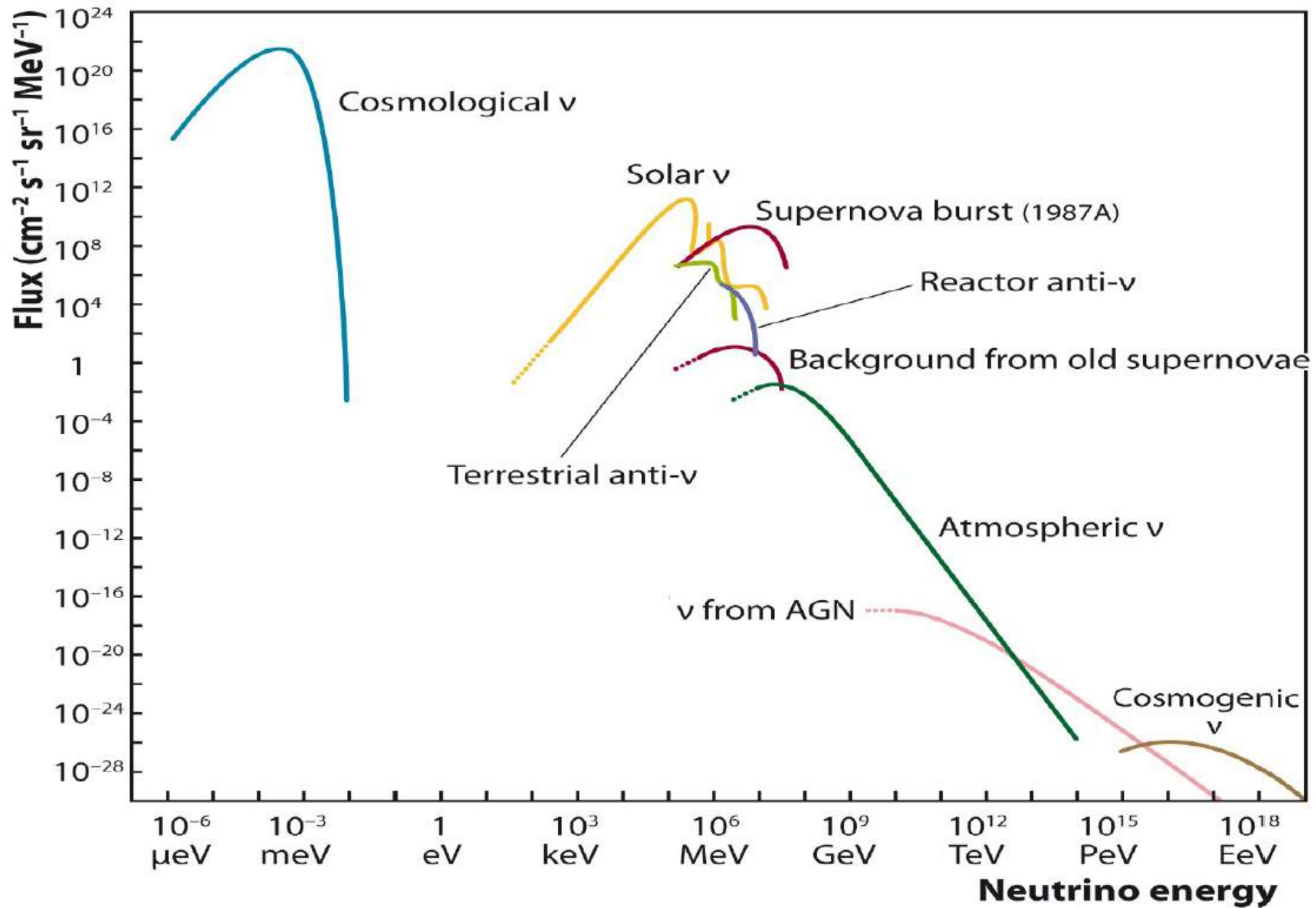
*Predicting neutrino mixing
from residual CP symmetries*



Neutrino sources & cross sections

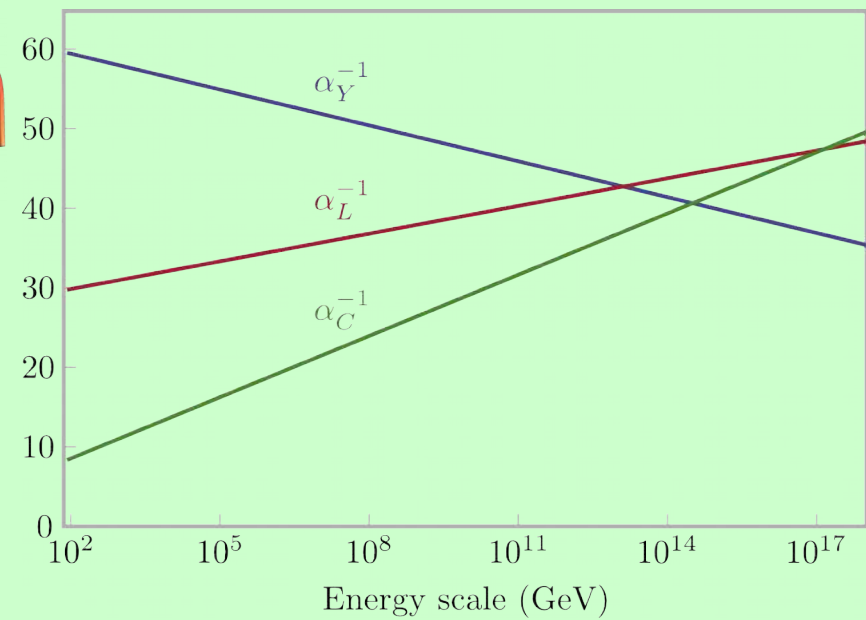


neutrino sources & fluxes



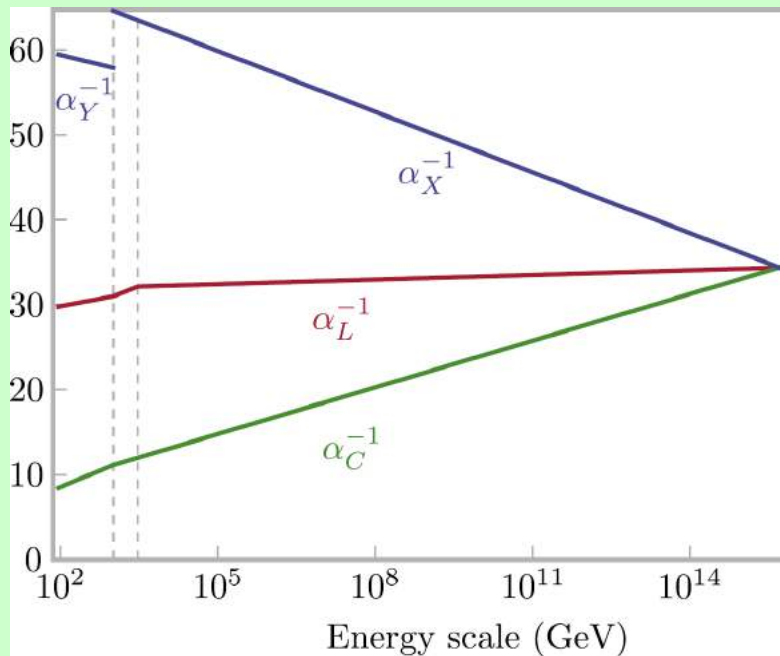
neutrinos may help unification

a near miss ...



What makes the gauge couplings unify? SUSY-GUT

But ... p decay, super-particles ...



The physics responsible for gauge coupling unification may also induce neutrino masses

Boucenna et al [Phys. Rev. D 91, 031702 \(2015\)](#)

Deppisch et al [Phys.Lett. B762 \(2016\) 432](#)