

Neutrino models

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- Neutrino:

Since neutrinos are electrically neutral, there exist two types of masses; Dirac and Majorana.

This Majorana nature can be determined by measuring the signal of neutrinoless beta decay, although it is not detected yet...

I will focus on Majorana type of neutrino.

Majorana neutrinos

One of the renowned ideas is to induce the neutrino masses through Weinberg operator:

$$(LHLH)/\Lambda \Rightarrow v^2/\Lambda \quad (v \ll \Lambda).$$

S. Weinberg, Phys. Rev. Lett. 43, 1566 (1979).

We have no way how *theoretically* to determine the scale and or structure of cut-off scale Λ .

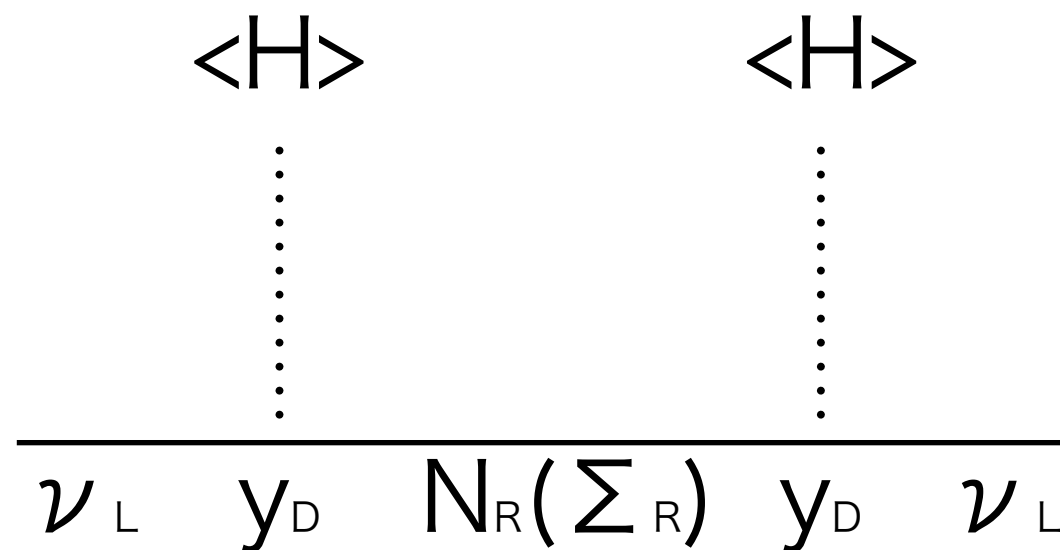
This question might lead to renormalizable theories.

Seesaw models

There are three representative types of models:

Mohapatra, Valle, Foot, Lew, He, Joshi, etc.

Type-1 (3):



$$m_\nu \equiv -m_D^T M_N^{-1} m_D; \quad D_\nu = U_{MNS}^T m_\nu U_{MNS}.$$

The neutrino mass matrix can be found by replacing

Λ into heavy mass of $N_R(\Sigma_R)$.

Casas-Ibarra(CI) parametrization:

$$y_D = i \frac{\sqrt{2}}{v} \sqrt{M_N} \mathcal{O} \sqrt{D_\nu} U_{MNS}^\dagger,$$

3 $N_R(\Sigma_R)$ case

$$\mathcal{O} = \begin{bmatrix} c_{\theta_{13}} c_{\theta_{12}} & c_{\theta_{13}} s_{\theta_{12}} & s_{\theta_{13}} \\ -c_{\theta_{23}} s_{\theta_{12}} - s_{\theta_{23}} s_{\theta_{13}} c_{\theta_{12}} & c_{\theta_{23}} c_{\theta_{12}} - s_{\theta_{23}} s_{\theta_{13}} s_{\theta_{12}} & s_{\theta_{23}} c_{\theta_{13}} \\ s_{\theta_{23}} s_{\theta_{12}} - c_{\theta_{23}} s_{\theta_{13}} c_{\theta_{12}} & -s_{\theta_{23}} c_{\theta_{12}} - c_{\theta_{23}} s_{\theta_{13}} s_{\theta_{12}} & c_{\theta_{23}} c_{\theta_{13}} \end{bmatrix},$$

2 $N_R(\Sigma_R)$ case

$$\mathcal{O} = \begin{bmatrix} 0 & 0 \\ \cos z & -\sin z \\ \pm \sin z & \pm \cos z \end{bmatrix}^T, \quad O = \begin{bmatrix} \cos z & -\sin z \\ \pm \sin z & \pm \cos z \\ 0 & 0 \end{bmatrix}^T,$$

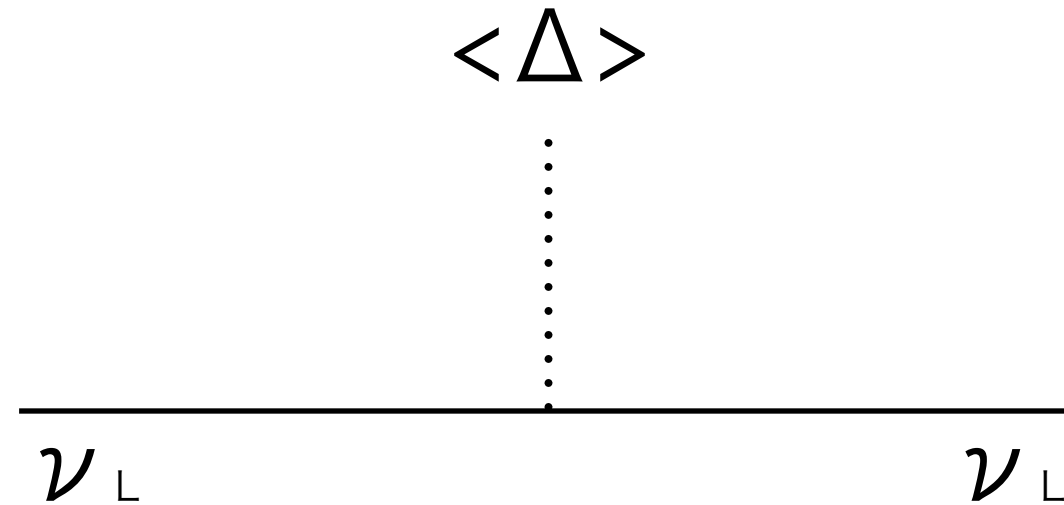
NH

IH

This is a useful parametrization not only to fit the current neutrino oscillation data but also to connect a theory and experiment.

Note: The mass of $N_R(\Sigma_R)$ has to be diagonal !

Type-2:



The neutrino mass matrix can be found by replacing v^2/Λ into VEV of Δ .

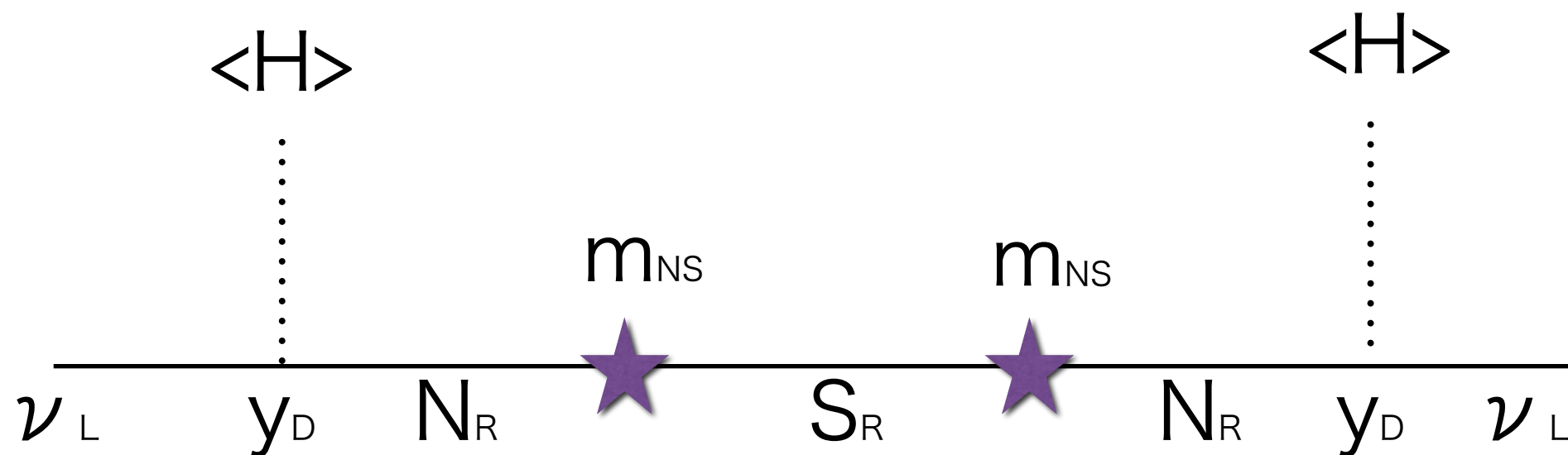
$$m_{\nu_{ab}} = \frac{y_{\Delta_{ab}}}{\sqrt{2}} v_{\Delta}.$$

CI parametrization: $y_{\Delta} = \frac{\sqrt{2}}{v_{\Delta}} U_{MNS}^* D_{\nu} U_{MNS}^{\dagger}.$

This is completely determined by the experimental values except VEV of Δ !

The other complicated but representative models introducing additional tiny Majorana fermions: S_R

1. Inverse seesaw:



Assuming $M_N=0$ and hierarchy: $\mu_S \ll y_D v < m_{NS}$,

$$m_\nu \approx m_D^T m_{NS}^{-1} \mu_S (m_{NS}^T)^{-1} m_D,$$

CI parametrization of Inverse seesaw

$$m_\nu \approx m_D^T \boxed{m_{NS}^{-1} \mu_S (m_{NS}^T)^{-1}} m_D,$$
$$= M_N^{-1}$$

Let us identify type-I like seesaw model by regarding the
red part as inverse of M_N .

However since inverse of M_N is not diagonal but symmetric,
one cannot apply the simplest CI !

One of the conventional solution is to apply Cholesky
Factorization(CF) into M_N

$$M_N^{-1} \equiv L_N^T D_N^{-1} L_N \quad \times \text{ } L_N \text{ is an triangular matrix}$$

Then neutrino mass matrix can be redefined by

$$m_\nu \approx m_D^T L_N^T D_N^{-1} L_N m_D \equiv M_D^T D_N^{-1} M_D$$

Thus one finds the CI parametrization just by the following replacements:

$$m_D \rightarrow M_D \text{ and } M_N \rightarrow D_N$$

Point: (DN , LN) can uniquely be fixed by components of MN, when DN is taken to be positive real.

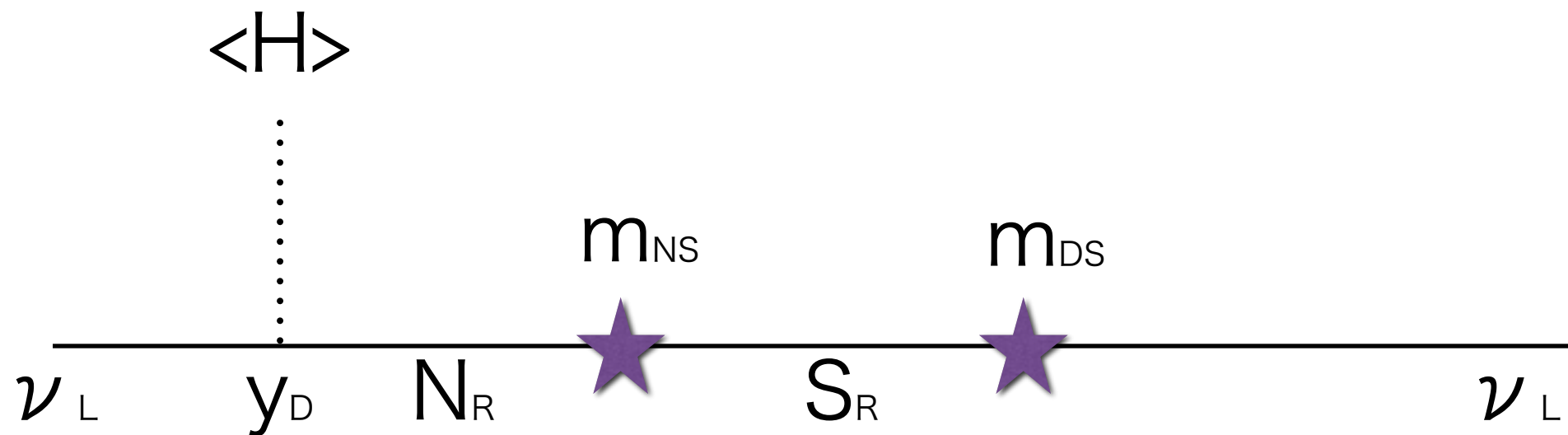
Generalized CI parameterization:

One can generalize 3xN right-handed matrix.

In this case, O is extended by 3 by N mixing matrix and its degrees of freedom is given by $3(N-2)$.

Note: O is not orthogonal matrix but just satisfies $OO^T = 1_{\{3 \times 3\}}$.

2. Linear seesaw:



Assuming $M_N=0$ and hierarchy: $m_{DS} \ll y_D v < m_{NS}$,

$$m_\nu \approx -m_{DS}^T (m_{NS}^T)^{-1} m_D - m_D^T (m_{NS})^{-1} m_{DS}$$

CI parametrization of linear seesaw is given by introducing anti-symmetric matrix A as:

$$m_D = -\frac{1}{2}m_{NS}^T(m_{DS}^T)^{-1}(U_{MNS}^*D_\nu U_{MNS}^\dagger + A).$$

$$\text{where} \quad A + A^T = 0.$$

Careful points to applying CI!

If you have a specific texture of a Yukawa coupling related to the neutrino mass matrix (such as two-zero texture) or the undefined inverse matrix (due to the reduction of rank), applying “CI” is not recommended.

Textures

H. Fritzsch, Z. Xing, S. Zhou, etc,...

Reduced matrix of y_D

Juan Herrero-Garcia, Miguel Nebot, Nuria Rius, Arcadi Santamaria, etc,...

Radiatively induced neutrino mass models:

It provides a concrete structure of Λ at **loop levels**.
=> Theory can be within low energy scale(\sim TeV).

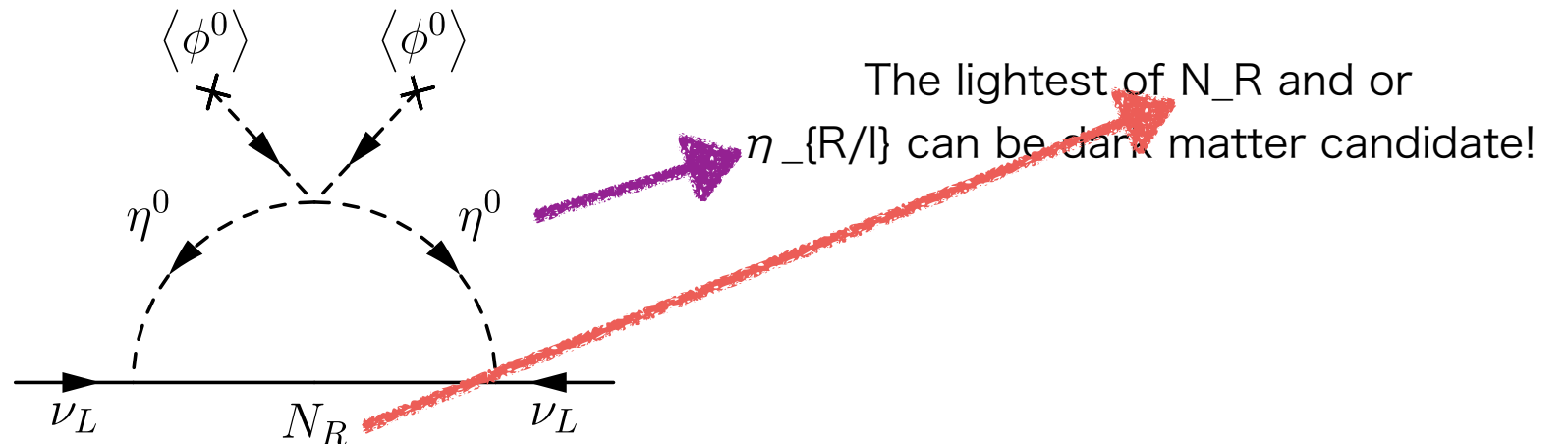
2. Dark matter candidate, muon anomalous magnetic dipole moment, origin of Baryon number asymmetry via leptogenesis, can be explained **linking to the active neutrinos!**

Note: One has to take care of Lepton flavor violations(LFVs), since they are always arisen from the related neutrino Yukawa couplings !!

Ma Model

SM+right-handed Neutrino +SU(2) doublet Scalar

	N^c	η
$SU(2)_L$	1	2
$U(1)_Y$	0	1/2
\mathbb{Z}_2	-1	-1



Interacting term

$$\mathcal{L} = (y_\nu)_{\alpha i} \eta^\dagger \bar{\ell}_\alpha N_i^c + \text{h.c.} - \mathcal{V}(\phi, \eta) + \frac{1}{2} M_i N_i N_i + H.c.$$

※ y_ν is restricted by LFVs at one-loop level.

Scalar potential term

$$\begin{aligned} \mathcal{V}(\phi, \eta) = & m_\phi^2 |\phi|^2 + m_\eta^2 |\eta|^2 + \frac{\lambda_1}{2} |\phi|^4 + \frac{\lambda_2}{2} |\eta|^4 + \lambda_3 |\phi|^2 |\eta|^2 \\ & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \frac{\lambda_5}{2} \left[(\phi^\dagger \eta)^2 + (\eta^\dagger \phi)^2 \right] \end{aligned}$$

Summary

Along this line of ideas, a vast of literature has arisen,
since nonzero θ_{13} has been measured (2012),

More than 110 ideas at one-loop level,

More than 40 ideas at two-loop level,

More than 30 ideas at three-loop level,

and a few ideas at four-loop level...

Another applications:

- Small mass SM fermions(e,u,d,muon) can be explained by introducing exotic charged fermions.
- Anomaly such as $b \rightarrow s, \mu \mu_{\bar{}}$ could be explained by introducing leptoquarks with/without flavor dependent gauged symmetries.
- Several anomalies of indirect detections from cosmic-rays as a DM candidate could be explained.
- Collider signatures could be expected in the near future.

Thanks.