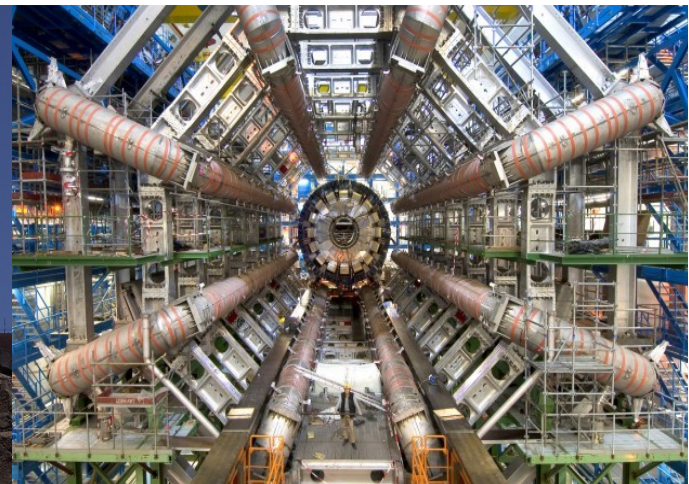
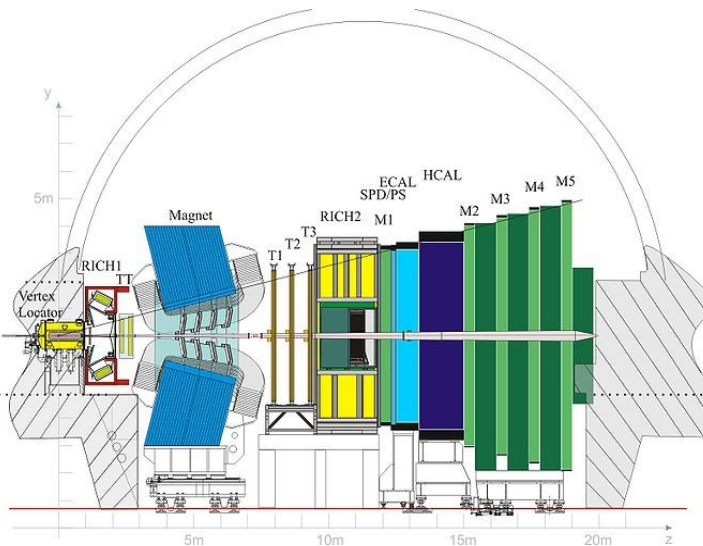


Hidden Valley Search Meets Cosmological Constraint

Lingfeng Li (Hong Kong U. of Sci. and Tech.)
and Yuhsin Tsai (U. of Maryland)

arXiv: 1812.XXXXX or 1901.XXXXX?

**Dec. 2018, NCTS Annual Theory Meeting:
Particles, Cosmology and Strings**



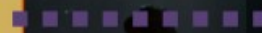
What's in the dark?

Standard Model

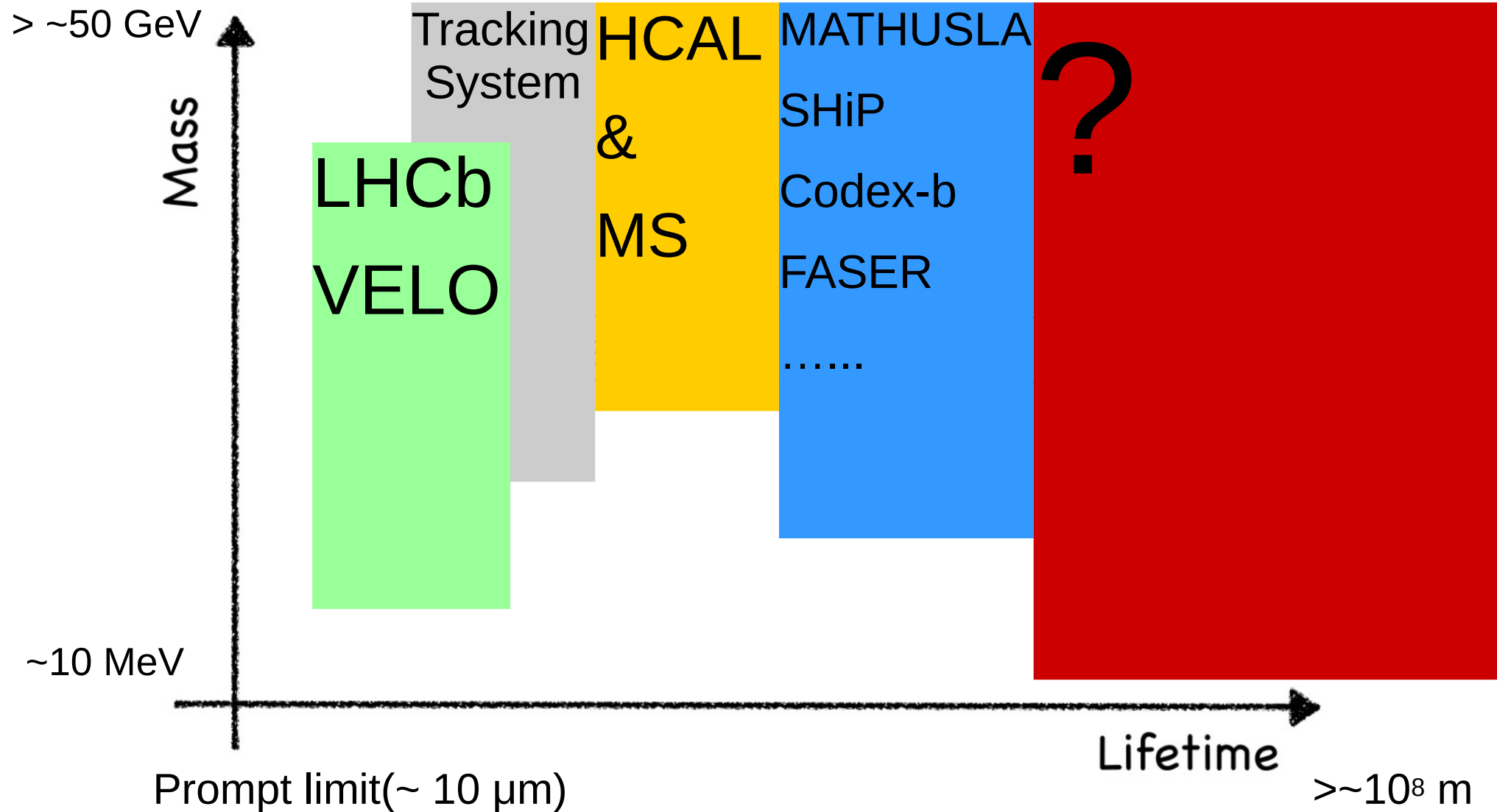
$U(1)$ $SU(2)_L$
 $SU(3)_c$ leptons
quarks
Higgs

Hidden Sector

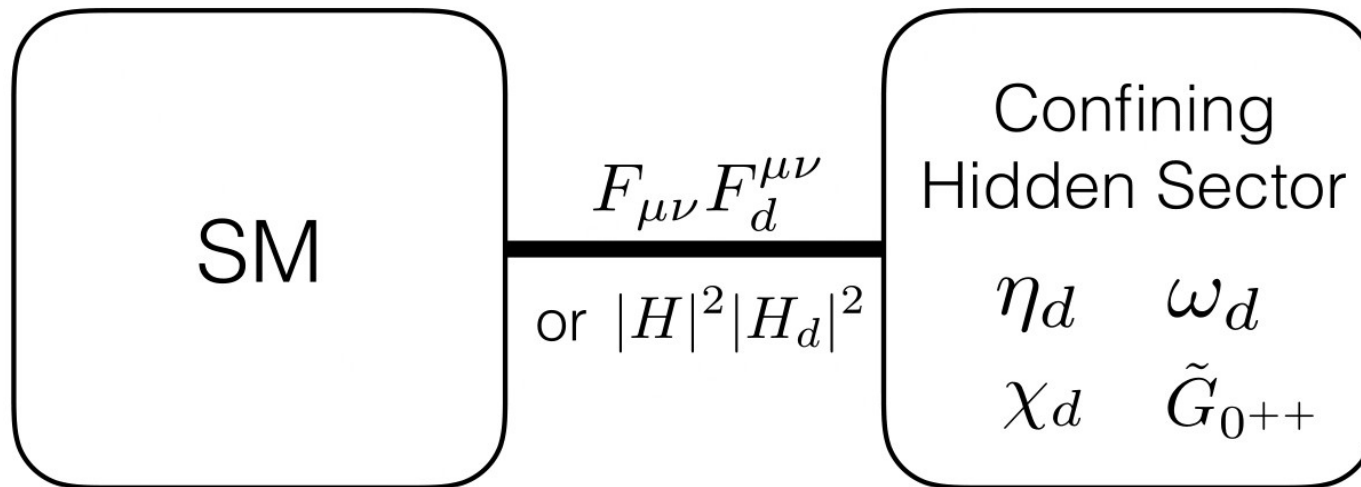
Hidden Force
dark photon?
dark QCD?



Hidden Valley models are good, but..



Basic Idea



- Cosmology **can** tell us where and how to look for dark hadrons at a collider.
- We consider the **heavy** mediator case (mediator mass $>\sim$ dark hadron masses)

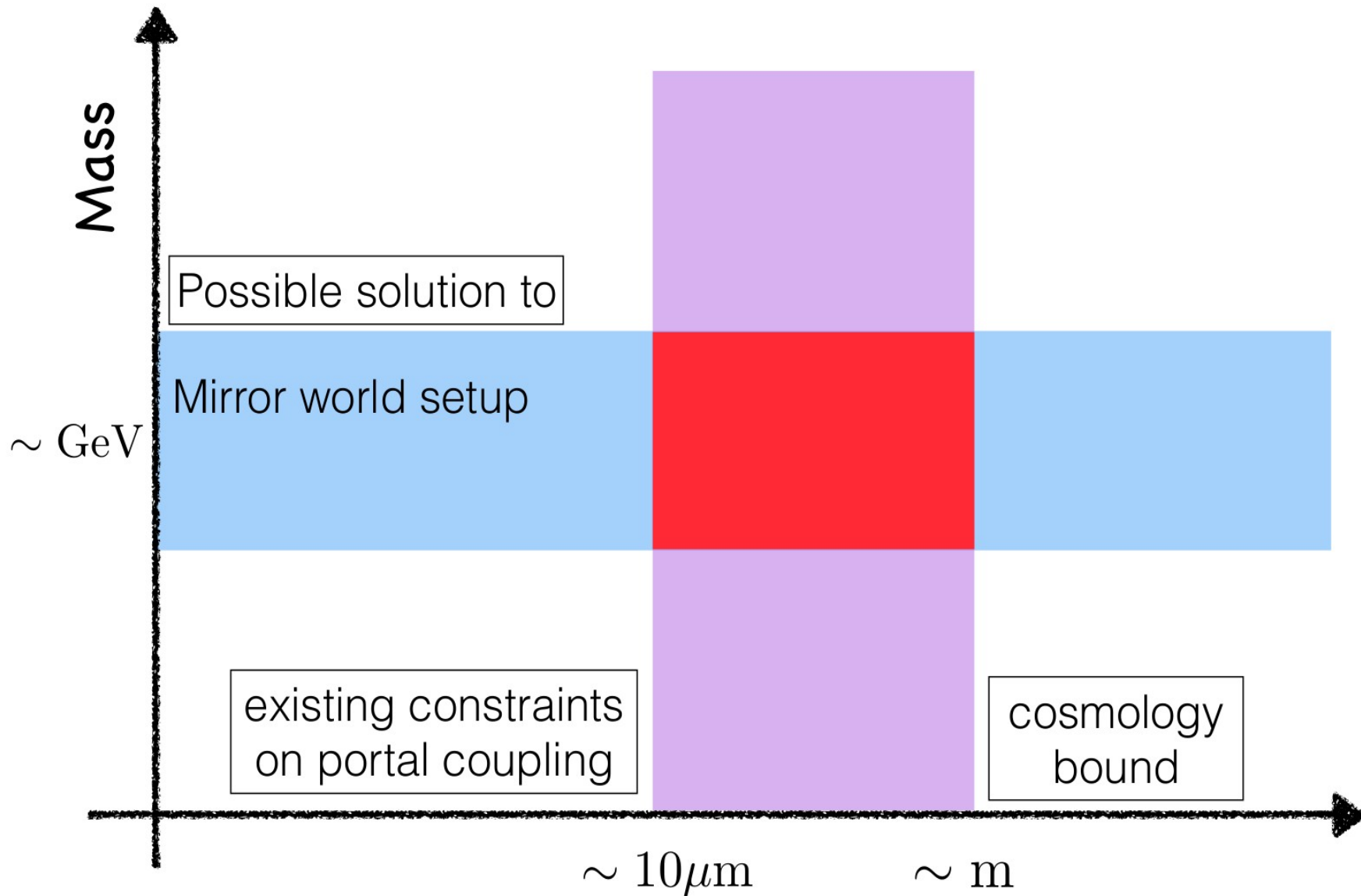
η_d : The pseudoscalar meson
(Lightest dark hadron)

χ_d : The scalar meson

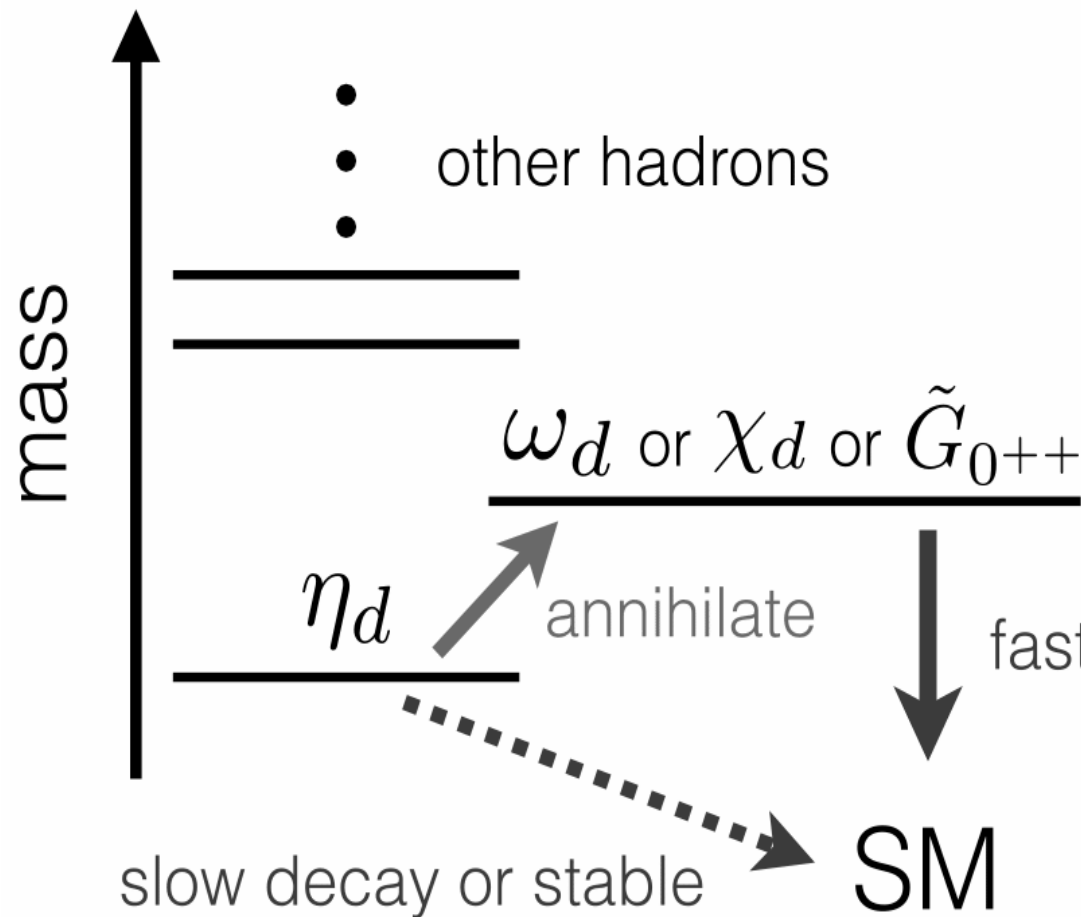
ω_d : The vector meson

G_{0++} : The dark glueball

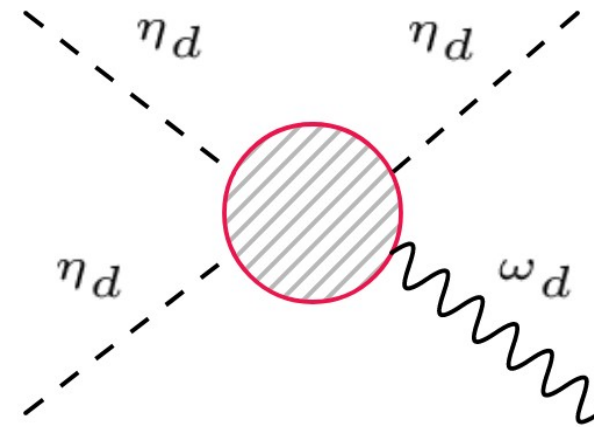
Thanks to Cosmological Constraints



Related discussions in [\[1512.02647 & 1809.10152\]](#)



Inelastic transition process:



Sensitive to mass gap, for related discussions see [\[1505.07107,1705.08450,1705.09292,1805.12139\]](#)

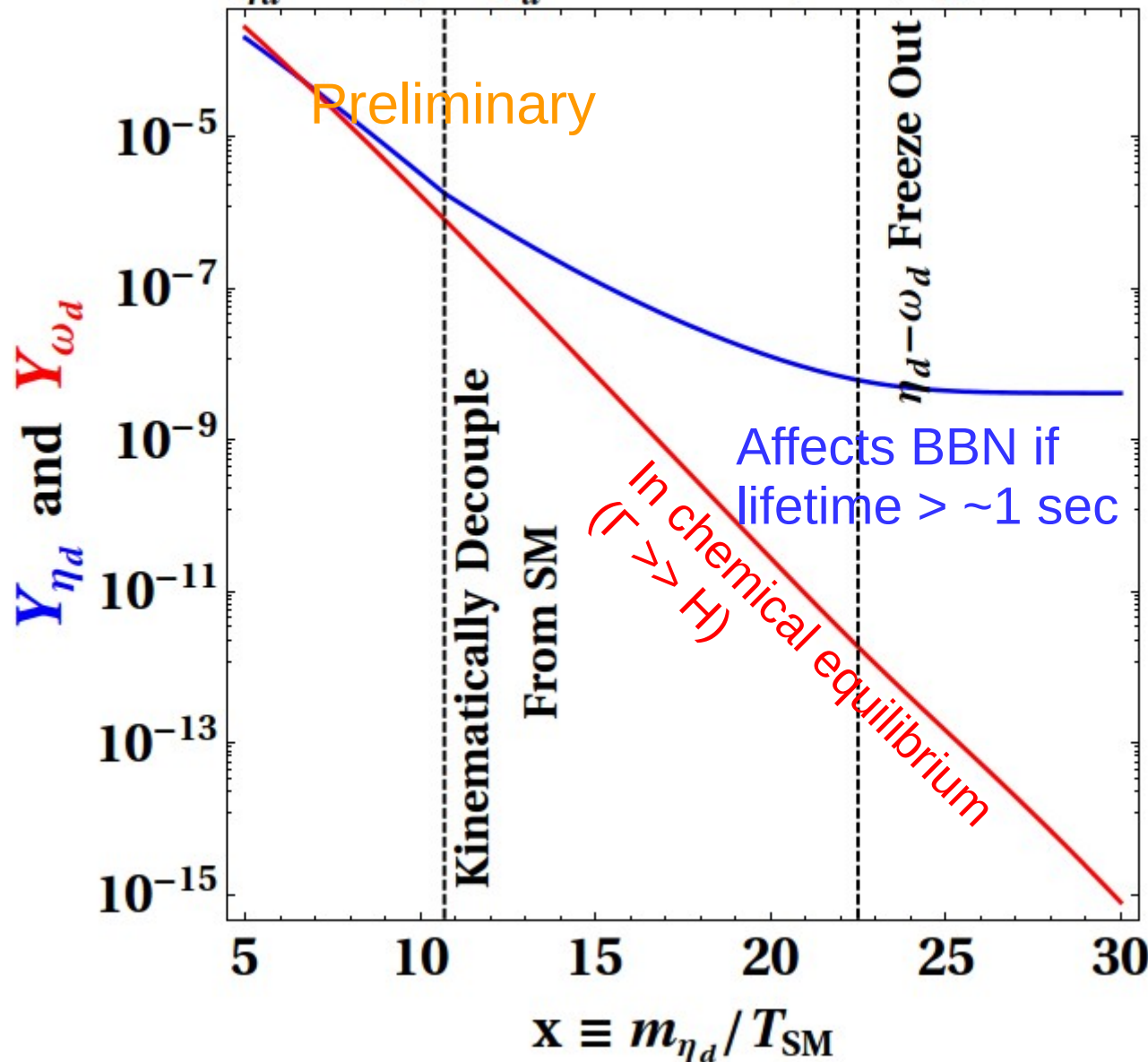
A fast decaying component (**Collider bounds apply here**) that depletes dark hadron number density

Slowly decaying/Stable ones become troublesome:

BBN/ CMB/ Ω_{CDM} bounds...

Thermal History

$$m_{\eta_d} = 10, m_{\omega_d} = 12 \text{ GeV}, c\tau(\omega_d) = 1 \text{ cm}$$

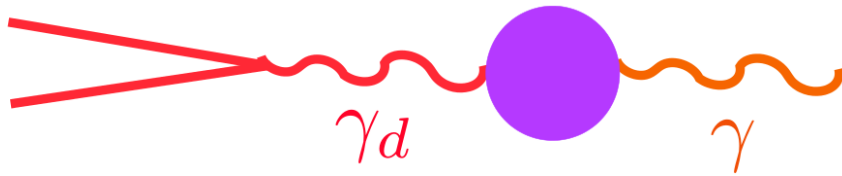


Sensitive to:

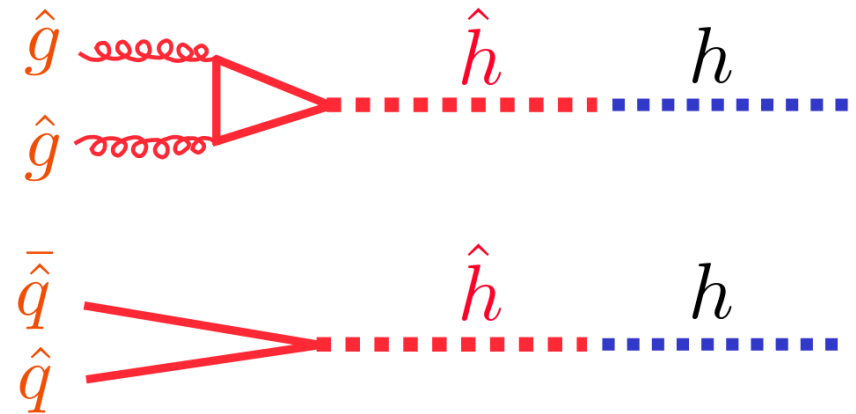
- Decay length Γ (kinematic decoupling time)
- ω_d - η_d mass difference Δm (when freeze out happens).

The analytical form is **not always precise**

Different Models of Hadron Decays



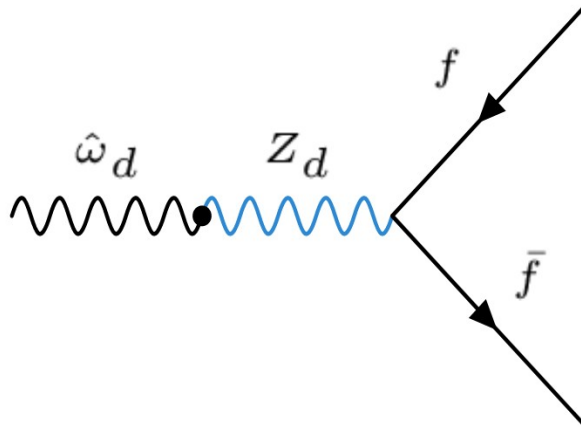
Dark Photon Portal
(A Decaying Vector)



SM Higgs Portal
(A Decaying Scalar)

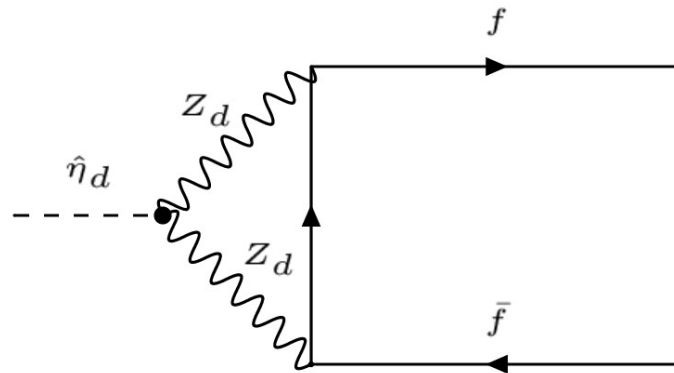
Dark Photon Portal Models

Vector and Pseudoscalar Decay



Vector Decay: **Fast**,
tree level process

- Still large range of possibility
- Can be out of collider reach if we do not consider the pseudoscalar decay.



$$\Gamma_{\eta_d \rightarrow f \bar{f}} = \frac{9\alpha_D^2 m_{\eta_d}^3}{8\pi^3 f_{\eta_d}^2} \left[\frac{\alpha}{2\pi} \epsilon^2 \left(\frac{m_f}{m_{\eta_d}} \right)^2 \left(\frac{m_{\eta_d}}{m_{Z_d}} \right)^2 \right]^2$$

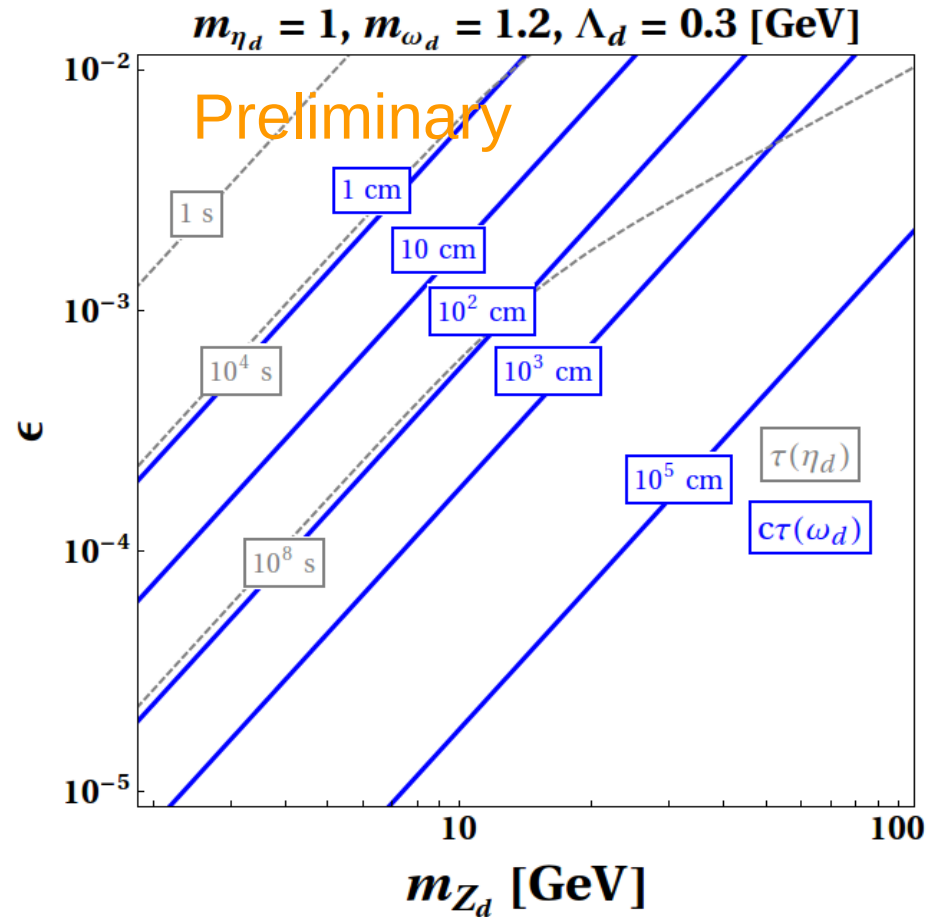
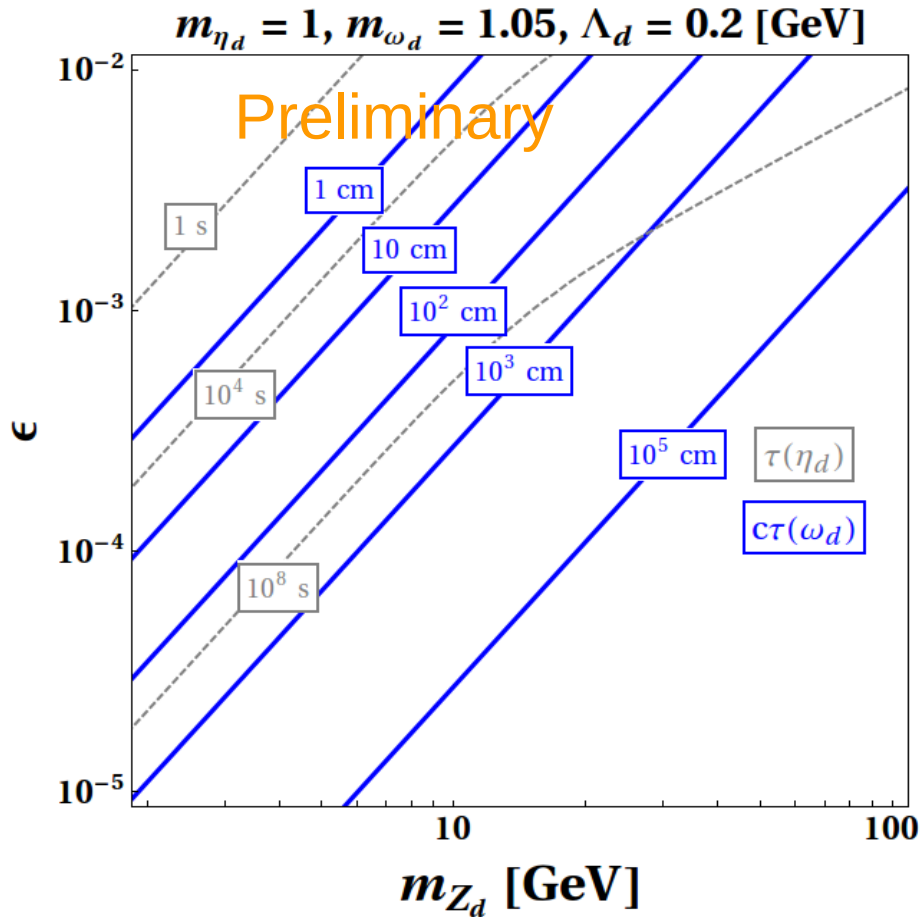
$$\Gamma_{\eta_d \rightarrow 4f} = \frac{9\alpha_D^2 m_{\eta_d}^3}{8\pi^3 f_{\eta_d}^2} \left[\frac{\alpha}{2\pi} \epsilon^2 \left(\frac{m_{\eta_d}}{2m_{Z_d}} \right)^4 \right]^2$$

Pseudoscalar Decay: **Slow**,
either 2 body (when Z_D is heavy)
or 4 body (lighter Z_D)
dominated

- Lifetime **longer than ~1 sec**,
- Introduces severe **constraints** during the BBN era

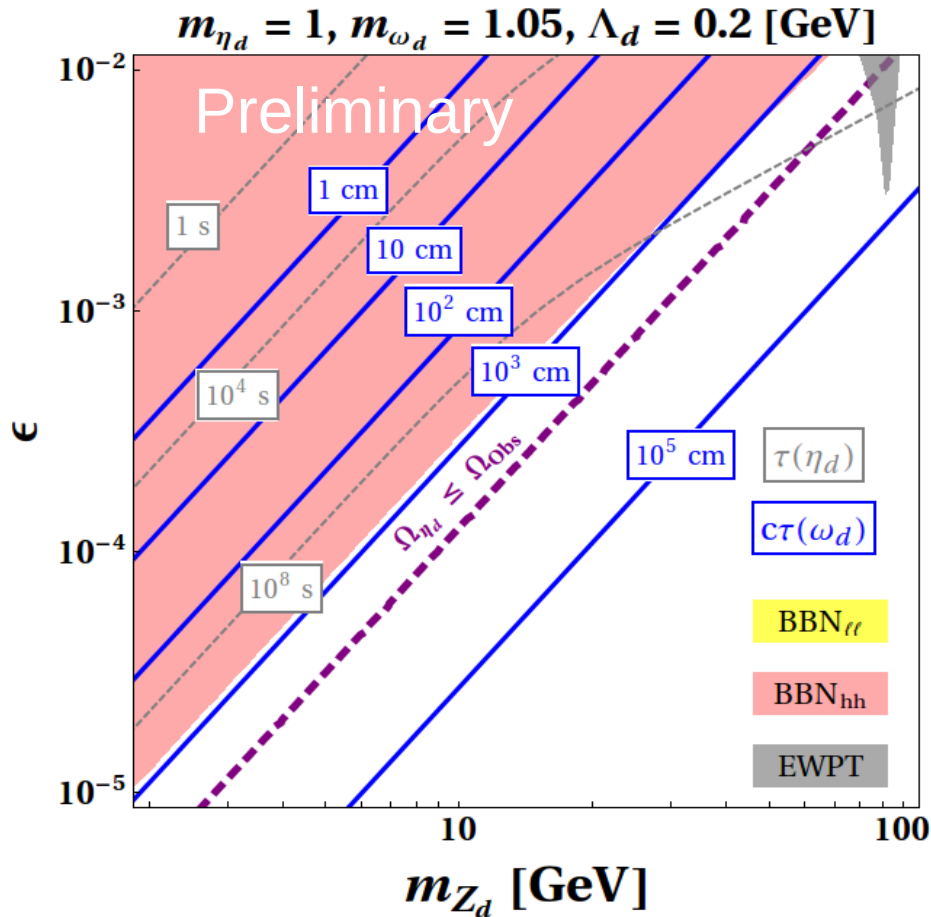
See [\[1512.02647 & 1805.09345\]](#)

Possible Decay Length



Wide range of possible lifetime

Dark Photon Portal Models

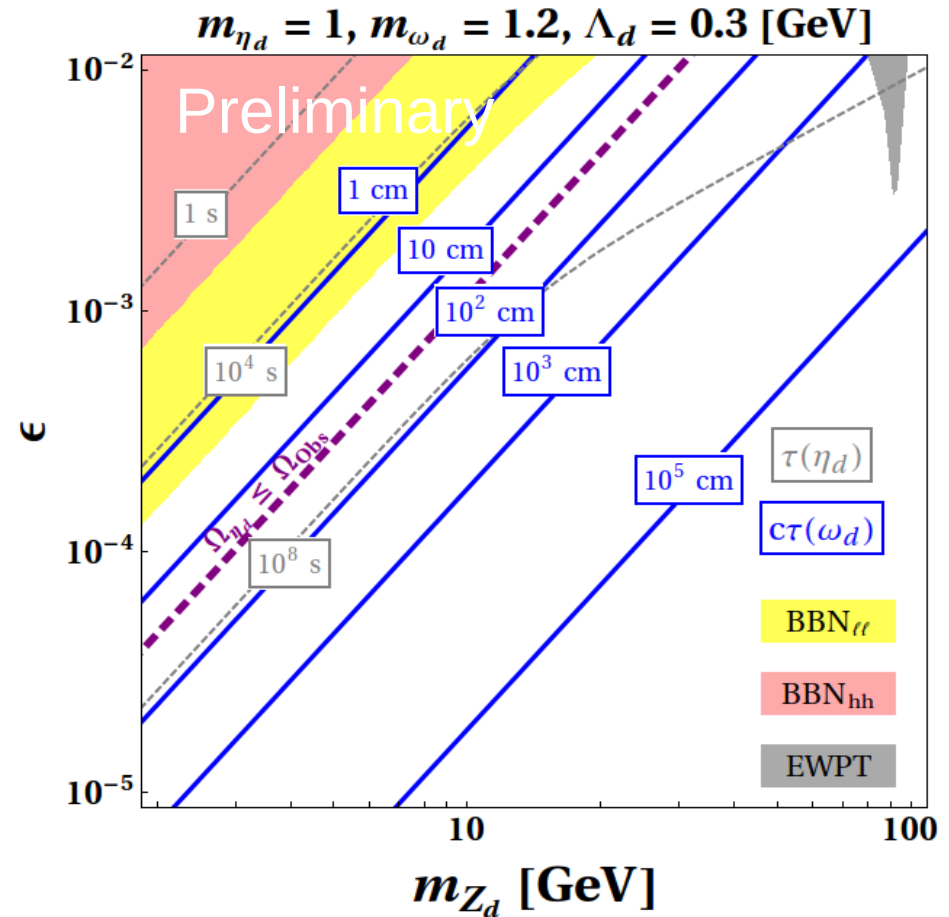


Small mass gap:

⇒ **weaker** $c\tau$ upper bounds

($< \sim 10$ m), within collider reach

For details about BBN, see [\[1709.01211\]](#)



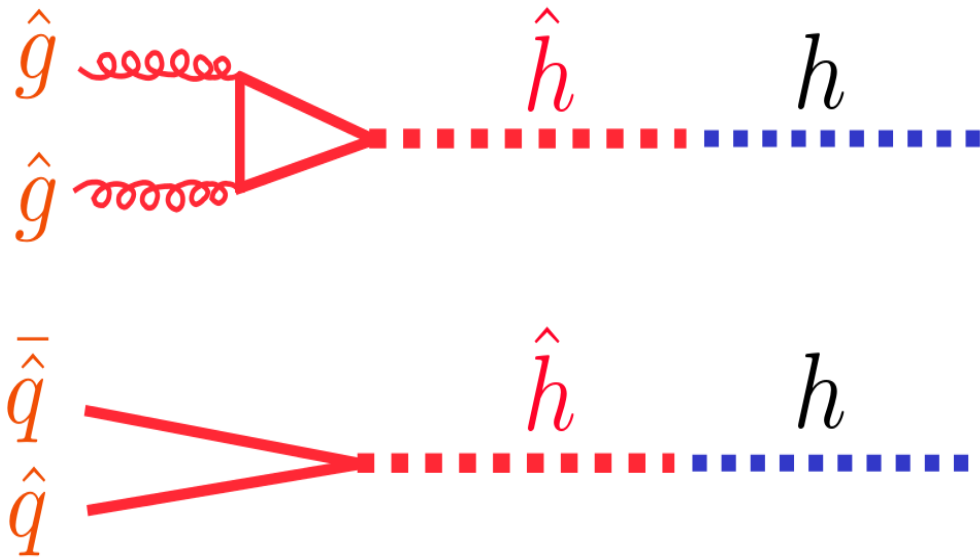
Moderate mass gap

⇒ larger Boltzmann suppression of the vector

⇒ $\omega_d - \eta_d$ freeze out happens earlier

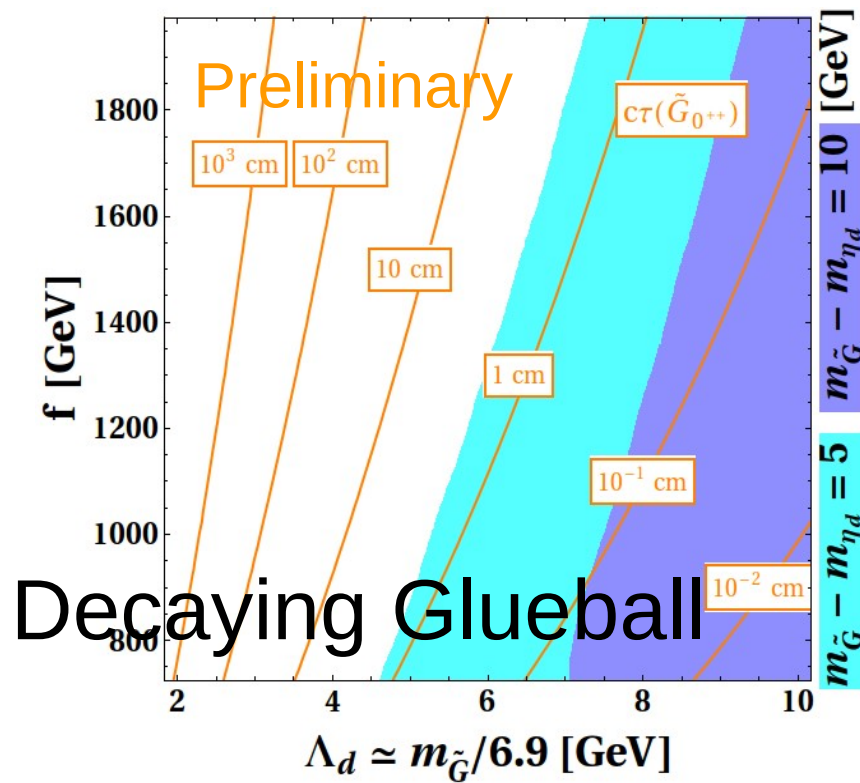
⇒ **stronger** $c\tau$ bounds (< 1 cm)

Higgs Portal Models: Scalar Hadrons Decay

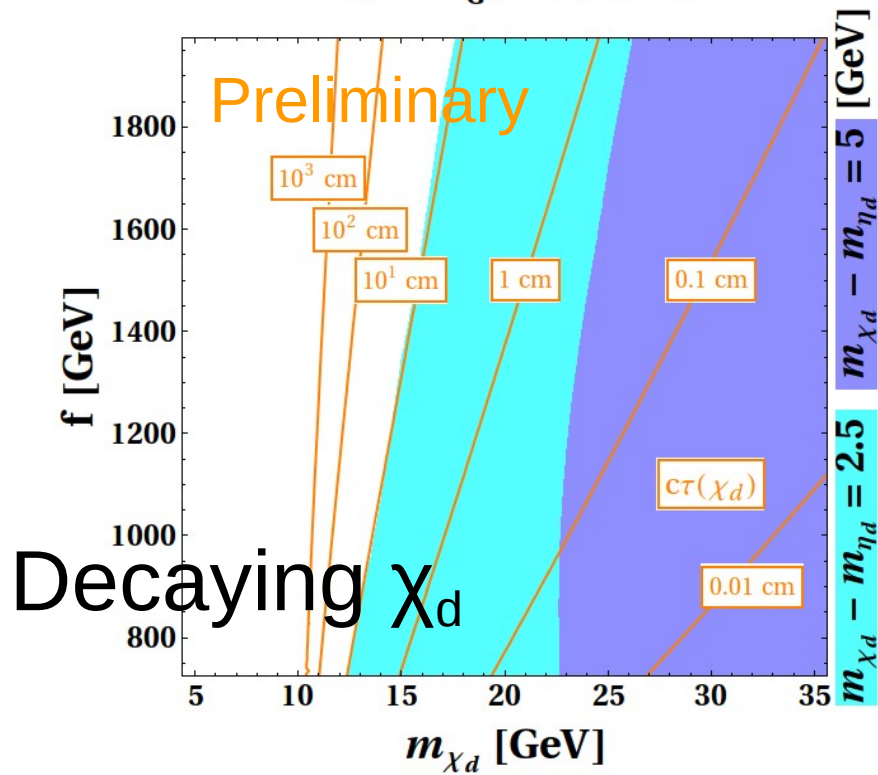


- Scalar Glueball:
Mass $\sim 7\Lambda$
Decay is loop induced at a higher scale (M), longer lifetime
- Scalar (χ_d) meson
Mass $\sim 2x$ fermion mass

Decay much faster

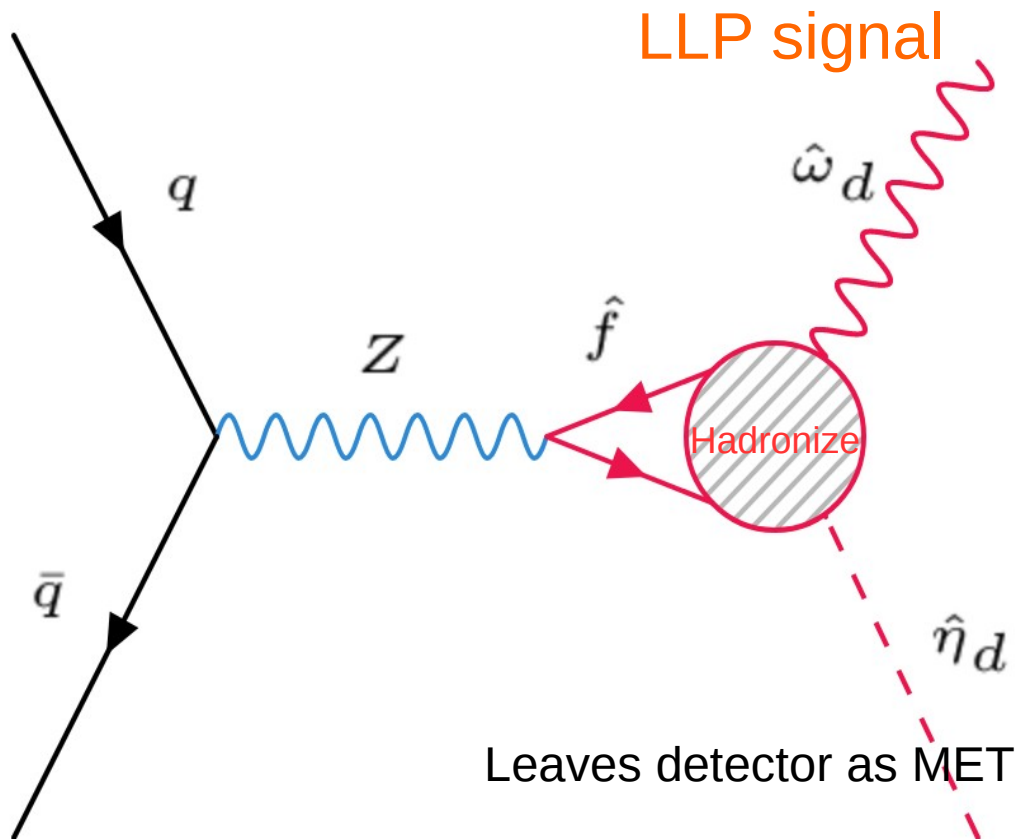


- Short $c\tau$ (< 10 cm)
- Heavier mass ($> \sim 15$ GeV)
- Hadronic Decays



LLP Detection at Hadron Colliders

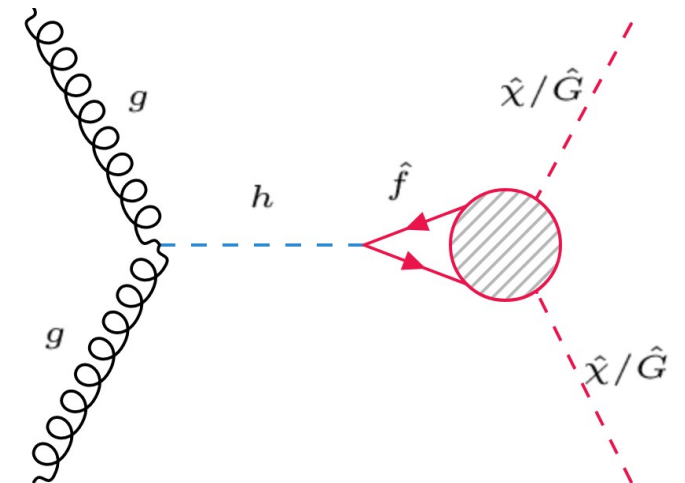
Exotic Z decay from Z_D mixing:



(Not thoroughly discussed, see [\[1710.07635 & 1712.07237\]](#))

Higgs portal:

Exotic Higgs decay



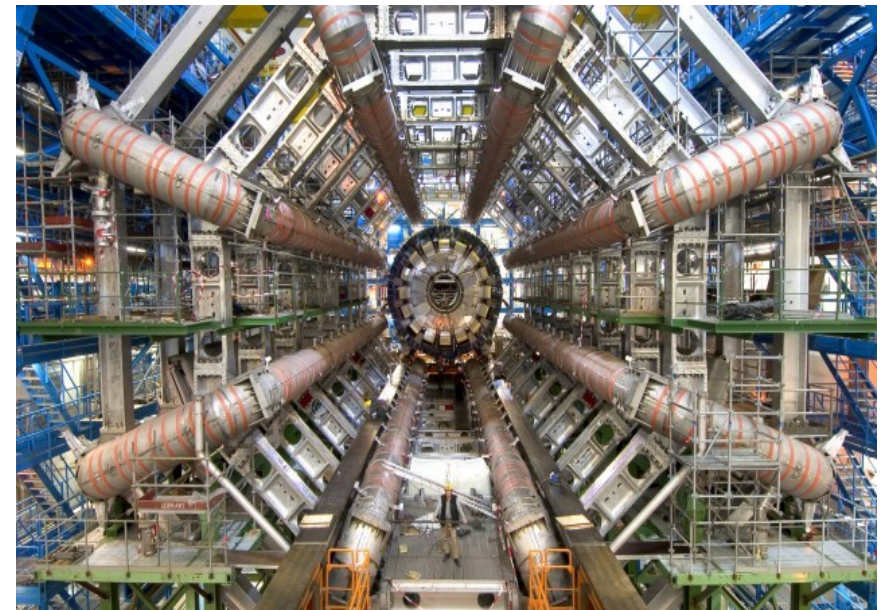
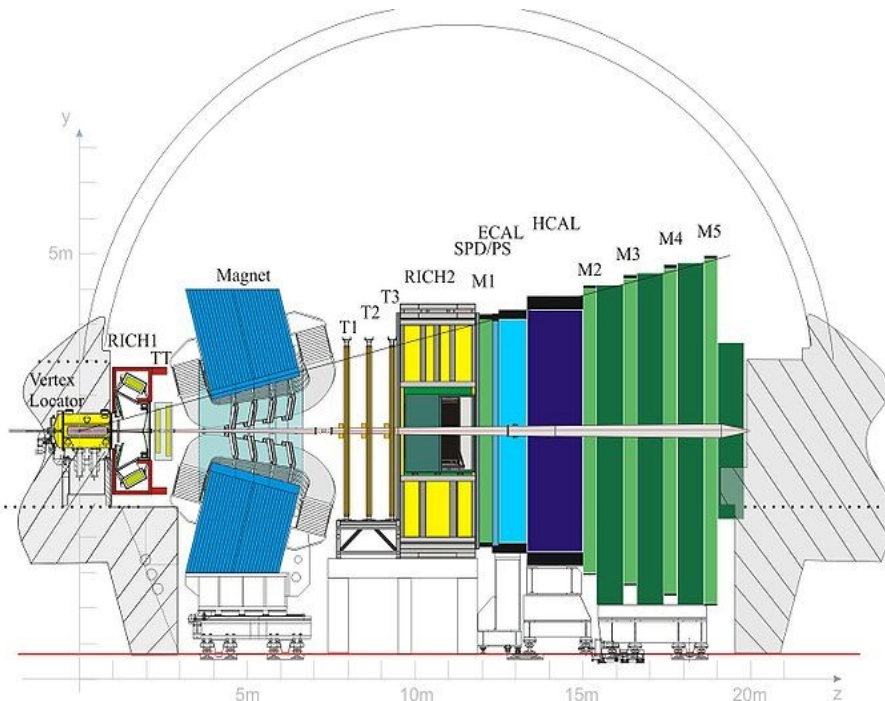
LLP Detection at Hadron Colliders

LHCb:

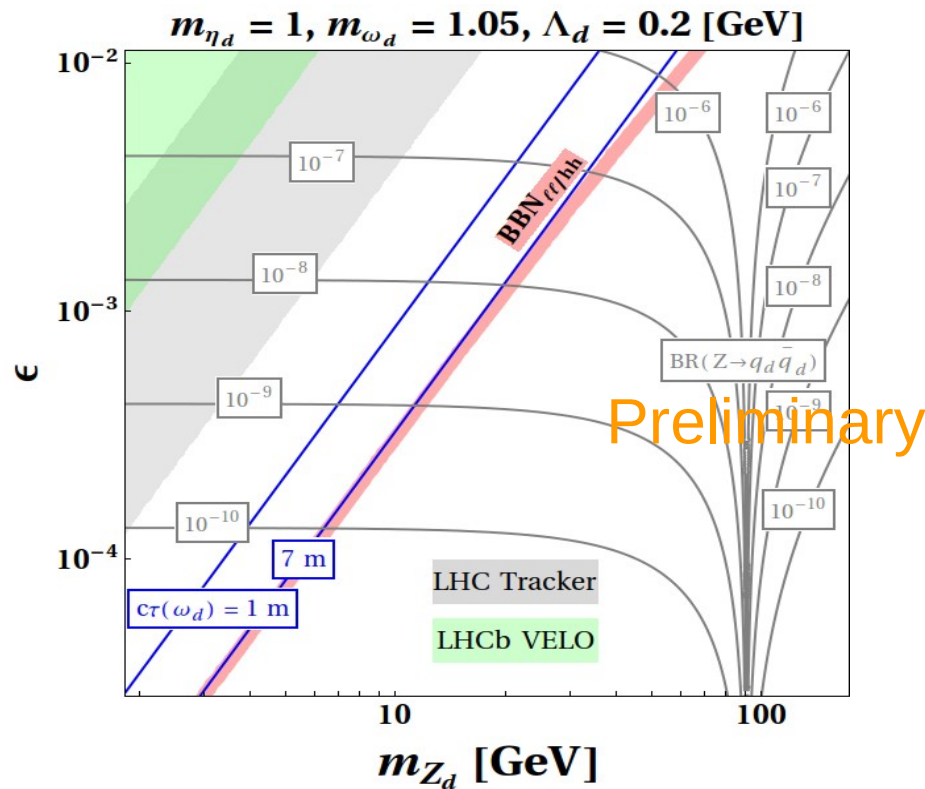
- Sensitive to small length scale ($< \sim 1$ cm)
- Good at LLP $\rightarrow \mu\mu/\pi\pi/KK$, etc.

CMS & ATLAS:

- Sensitive to wider range of decay length
- Larger luminosity and coverage
- Larger SM backgrounds

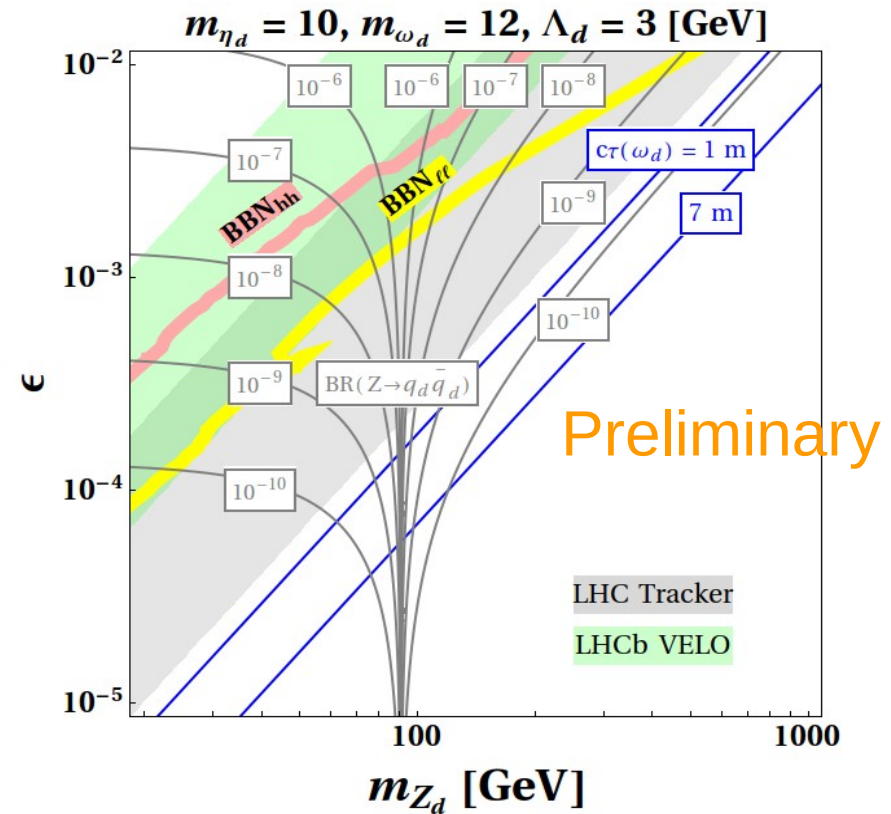


LLPs from Z_D -Z mixing: Exotic $Z \rightarrow$ dark hadron BR.



Green Shade: Chance to decay inside the LHCb VELO $>10\%$

Gray Shade: Chance to decay inside the LHC Trackers $>10\%$

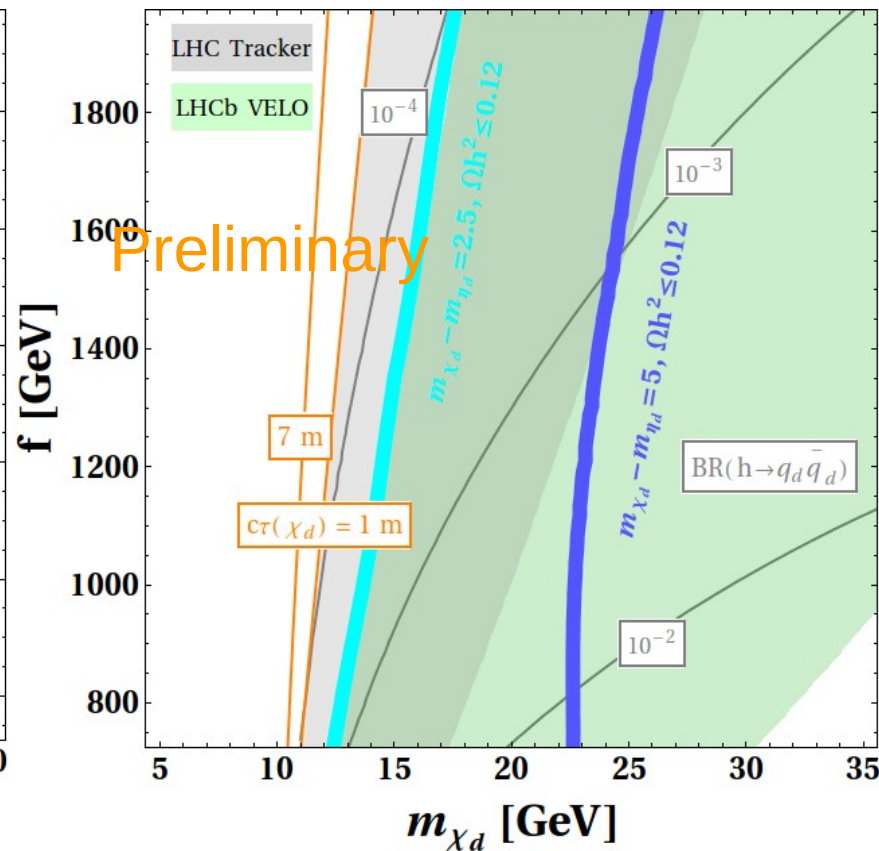
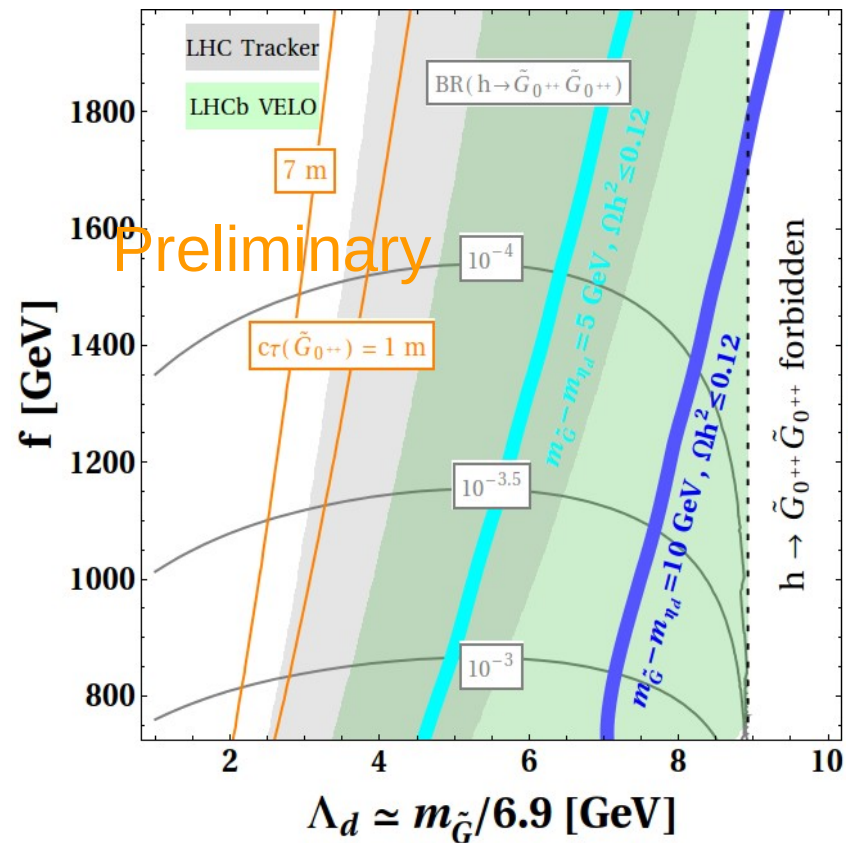


DV($\mu\mu$) + X: probe exotic BR down to $O(10^{-6})$ at LHCb and comparable bounds ATLAS/CMS after run III, for $ct \sim 1 \text{ cm}$

See [\[1808.03057 & 1806.07396\]](#)

LLPs from Higgs Portal Models: Higg Exotic Decays

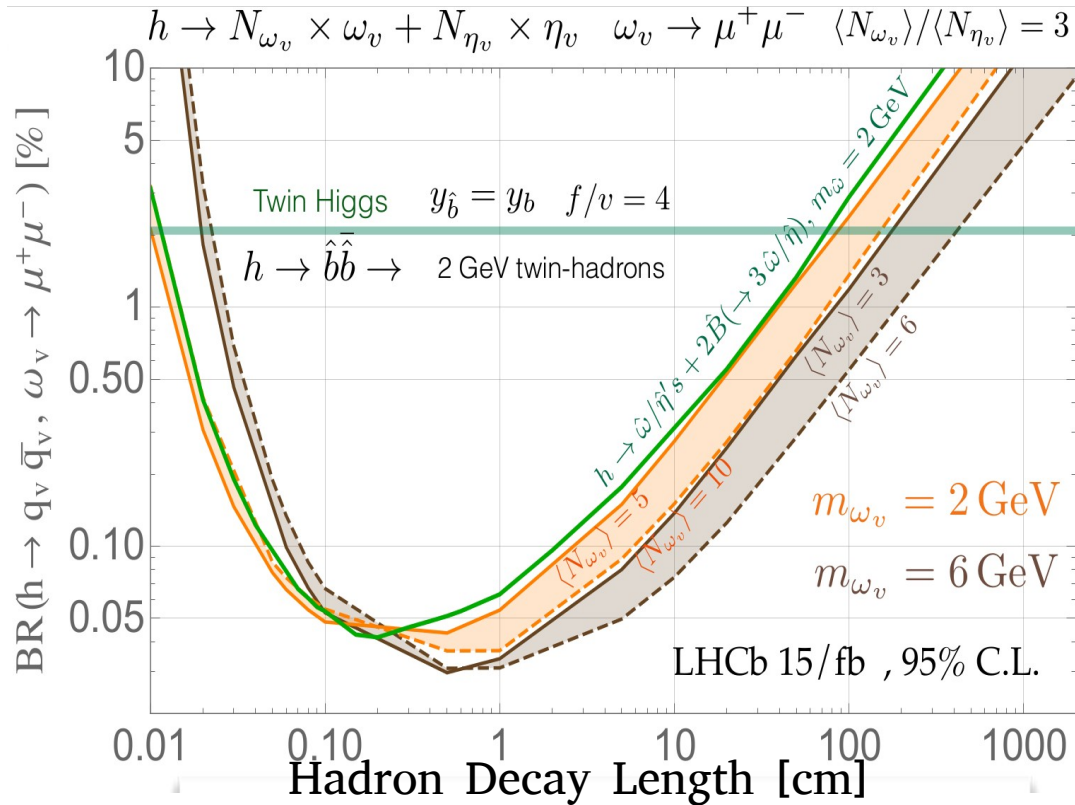
Displaced vertexes of **dijets** and **short** decay length.
Preliminary



Higgs exotic
BR: $\mathcal{O}(10^{-3})$ at
ATLAS/CMS
and $\mathcal{O}(10^{-2})$ at
LHCb during
HL-LHC era

See
[1806.07396 &
1808.08865]

Hidden Strong Dynamics



Some numerical results for this scenario when lifetime is short (still in work)

An related alternative scenario in Hsin-Chia's talk on Thursday.

↑↑ Some results from [\[1708.05389\]](#)

see also [\[1502.05409 & 1503.00009\]](#)

STAY TUNED!

LHCb can give the leading bound for short length ($\sim 1\text{cm}$) cases

Dinner Time!

Take Home Message:

- Cosmological constraints for hidden valley models can be stronger than often considered.
- Interesting predictions about future LLP searches, in the short length scale.

Backup Slides

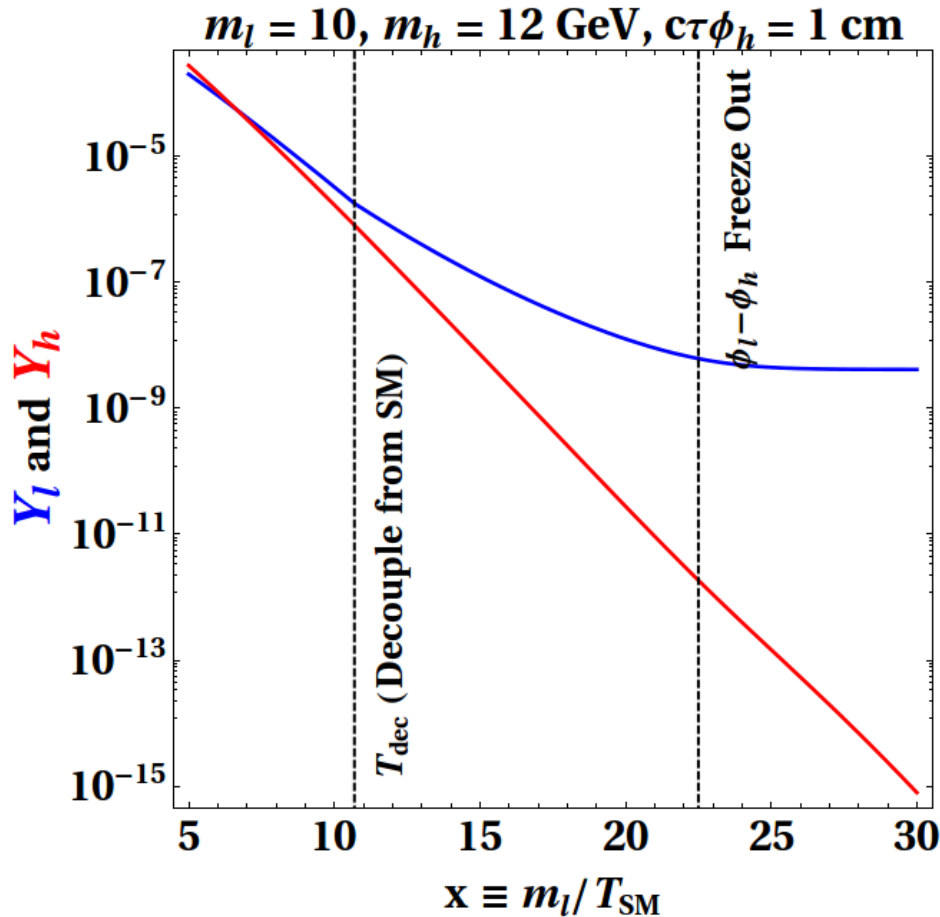
Thermal History

$$\frac{dY_h}{dx} = \frac{-1}{3H(x)} \frac{ds}{dx} \left[\langle \sigma_{+2h\nu} \rangle Y_l^2 - \langle \sigma_{-2h\nu} \rangle Y_h^2 - \langle \sigma_{-h\nu} \rangle Y_h Y_l + \langle \sigma_{+h\nu} \rangle Y_l^2 - \frac{\langle \Gamma_{\phi_h \rightarrow \text{SM}} \rangle \hat{T}}{s} Y_h + \frac{\langle \Gamma_{\phi_h \rightarrow \text{SM}} \rangle T}{s} Y_h^{\text{eq}}(T) \right],$$

$$\frac{dY_l}{dx} = \frac{-1}{3H(x)} \frac{ds}{dx} \left[\langle \sigma_{-2h\nu} \rangle Y_h^2 - \langle \sigma_{+2h\nu} \rangle Y_l^2 + \langle \sigma_{-h\nu} \rangle Y_h Y_l - \langle \sigma_{+h\nu} \rangle Y_l^2 \right].$$

↑↑ Solve the Boltzmann Eq. Above, we'll be all good...

Thermal History



$$Y_l(T) \simeq \frac{Y_l^{\text{eq}}(\hat{T})}{Y_h^{\text{eq}}(\hat{T})} Y_h^{\text{eq}}(T)$$

$$\propto \exp\left(\frac{\Delta m T_{\text{dec}}}{T^2}\right) \exp\left(-\frac{m_h}{T}\right)$$

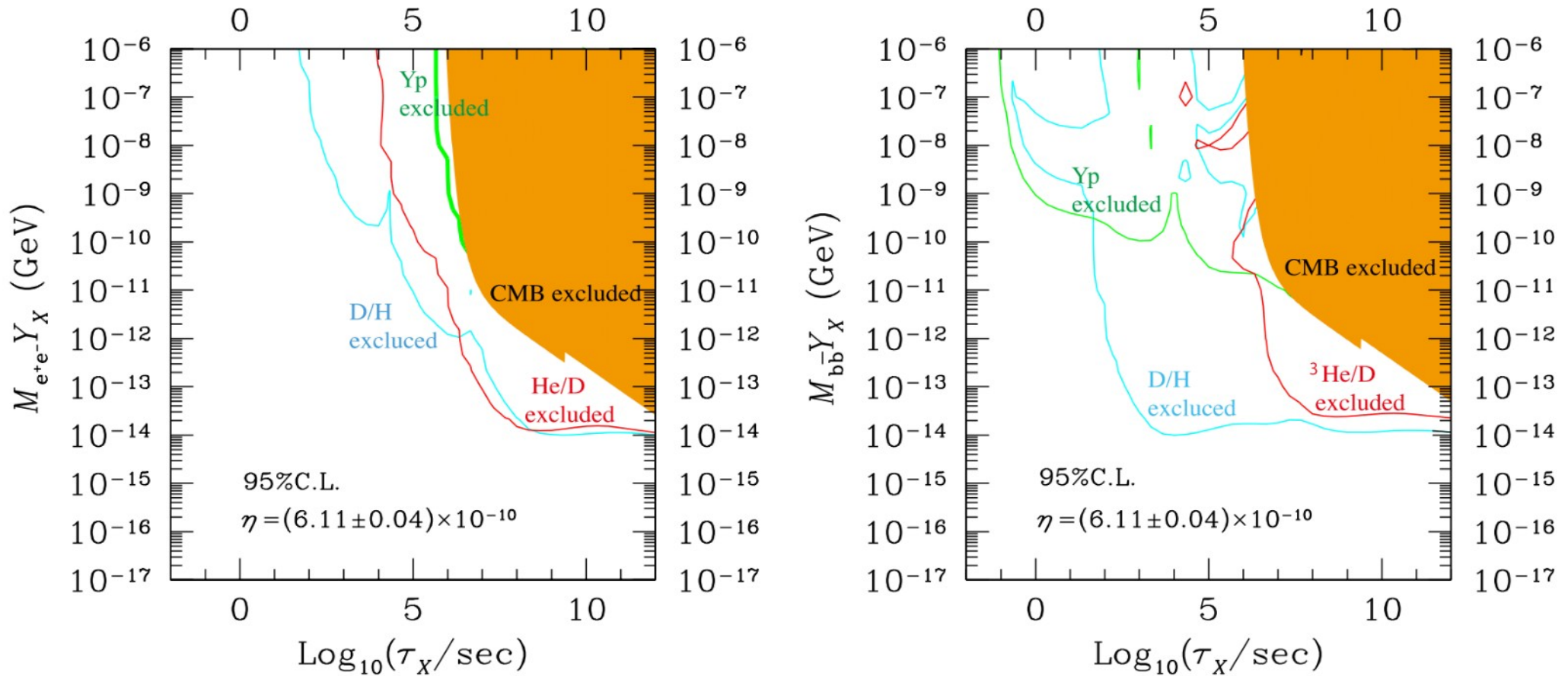
Sensitive to h decay length Γ (related to T_{dec}) and Δm (related to freeze out temperature).

Analytical Spoiler:

$$Y_l|_{\text{BBN}} \simeq 10^{-4} \left(\frac{1000 m_h^2}{M_P \Gamma_h} \right)^{\frac{\text{const} \times m_h}{\Delta m}} \exp\left(-\text{const} \frac{m_h}{\Delta m}\right)$$

Reasonable form but **not always precise** $\Rightarrow$ Numerical results

BBN Constraint



Taken from [\[1709.01211\]](#)

Estimate the Decay Chance in Valid Detector Volumes

LHCb VELO:

- Distance: 0.1-22 mm
- Pseudorapidity coverage: 2-5

LHC Tracking system:

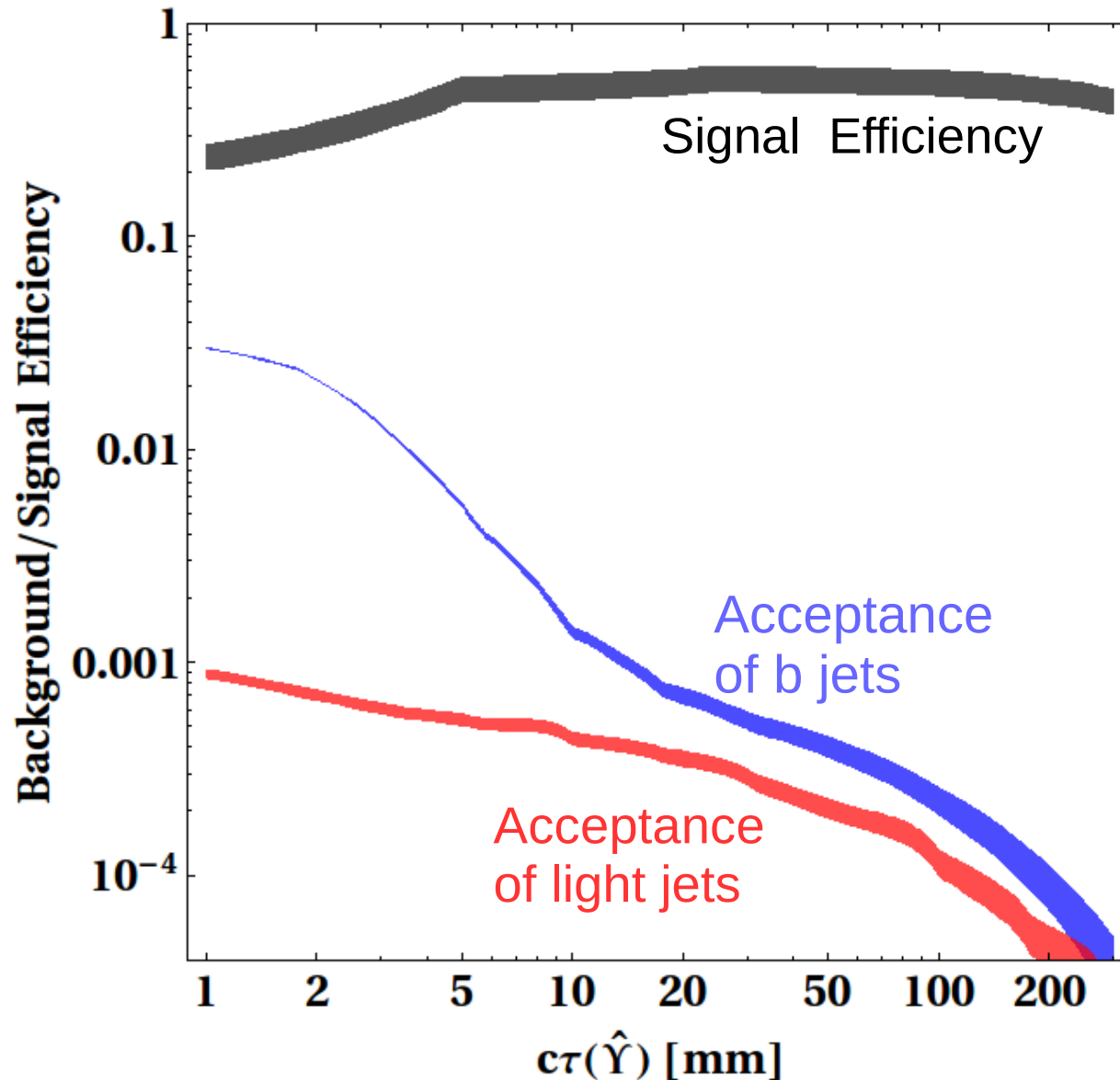
- Distance: 1-30 cm
 - Pseudorapidity coverage $[-3,3]$
-
- Z/h exotic decay estimated from MG5 simulations.

Estimate the di-muon LLP reach

- Z production rate at LHC: ~ 50 nb
- Luminosity: 15 fb^{-1} at LHCb and 300 fb^{-1} at ATLAS/CMS
- Background for $DV(\mu\mu)$ estimated as ~ 25
- Assume the $DV(\mu\mu)$ reconstruction efficiency ~ 0.5
- The $\omega_d \rightarrow \mu\mu$ BR $\sim 10\%$
- Needed #signal ~ 10
 $\Rightarrow \text{BR}(Z \rightarrow 2q_d) \sim \mathcal{O}(10^{-6})$

At LHC main detectors, the situation is more complicated due to MET triggering efficiency and background estimation. Reprojected from references.

Tell Emerging Jets from Ordinary Jets (Preliminary)



In collaboration with
H.C. Cheng, E.
Salvioni and C.
Verhaaren