

# Quantum Mirror Map

for

## Del Pezzo Geometries

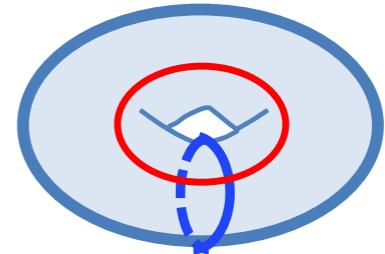
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# Susy. Gauge Theory



encoded in algebraic curve

e.g.) D5 del Pezzo

$$0 = \left[ \left( Q^{1/2} + Q^{-1/2} \right) \left( P^{1/2} + P^{-1/2} \right) \right]^2 - z/\alpha$$

A-period

Mirror Map  $\Pi_A(z) \sim c_1 z + c_2 z + \dots$   
redefining the variables

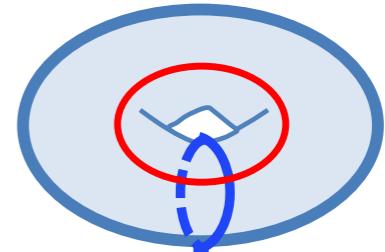
B-period

Free energy relating to  
BPS indices

not well studied

well studied

# Susy. Gauge Theory



encoded in algebraic curve

e.g.) D5 del Pezzo  $0 = \left[ \left( Q^{1/2} + Q^{-1/2} \right) \left( P^{1/2} + P^{-1/2} \right) \right]^2 - z/\alpha$

A-period

Mirror Map  $\Pi_A(z) \sim c_1 z + c_2 z + \dots$

redefining the variables

We need this to reach  
important results.

e.g. Free energy of ABJM

$$F = (\text{pert.}) + (\text{non-pert.}) \quad \text{for CS level}$$

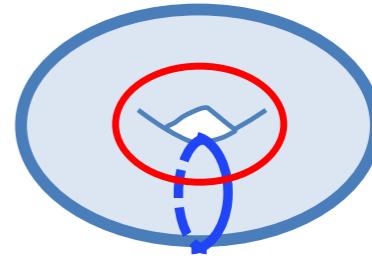


A-period



No pole

# Susy. Gauge Theory



- Some kinds of ABJM has group theoretical structure coming from symmetry of curve.  
(Kubo-Moriyama-Nosaka(2018))
- Free energy corresponds to B-period.

Then, what about A-period ?

e.g. Free energy of ABJM

$$F = (\text{pert. poles}) + (\text{non-pert. poles}) \quad \text{for CS level}$$

A-period

No pole

Generalized ABJM theory [Honda, Moriyama, 2014]  
 $(U(N)_k \times U(N)_{-k} \times U(N)_k \times U(N)_{-k})$

encoded in  $D_5$  del Pezzo with quantization

$$\hat{H} = \left[ \left( \hat{Q}^{1/2} + \hat{Q}^{-1/2} \right) \left( \hat{P}^{1/2} + \hat{P}^{-1/2} \right) \right]^2 - z/\alpha$$
$$\hat{Q}\hat{P} = q\hat{P}\hat{Q}$$

$$\Pi_A(z) \longrightarrow \Pi_A(z, \hbar)$$
$$(q = e^{i\hbar})$$

Quantum Mirror Map

Generalized ABJM theory [Honda, Moriyama, 2014]  
 $(U(N)_k \times U(N)_{-k} \times U(N)_k \times U(N)_{-k})$

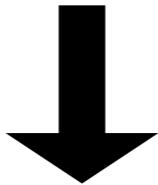
encoded in  $D_5$  del Pezzo with quantization

$$\hat{H} = \left[ \left( \hat{Q}^{1/2} + \hat{Q}^{-1/2} \right) \left( \hat{P}^{1/2} + \hat{P}^{-1/2} \right) \right]^2 - z/\alpha$$

$$\hat{Q}\hat{P} = q\hat{P}\hat{Q}$$

$[SU(2)]^3$  sym.

- $[SU(2)]^3$  comes from breaking of  $D_5$  sym.
- The sym. breaking is observed in B-period  
(Kubo-Moriyama-Nosaka(2018))



We consider quantum mirror map of  $D_5$  del Pezzo

# Result

A-period is similar to B-period

- expressed by Weyl characters
- coef. of characters = Integer
- has same reps. as those in B-period

# Plan

A-period is similar to B-period

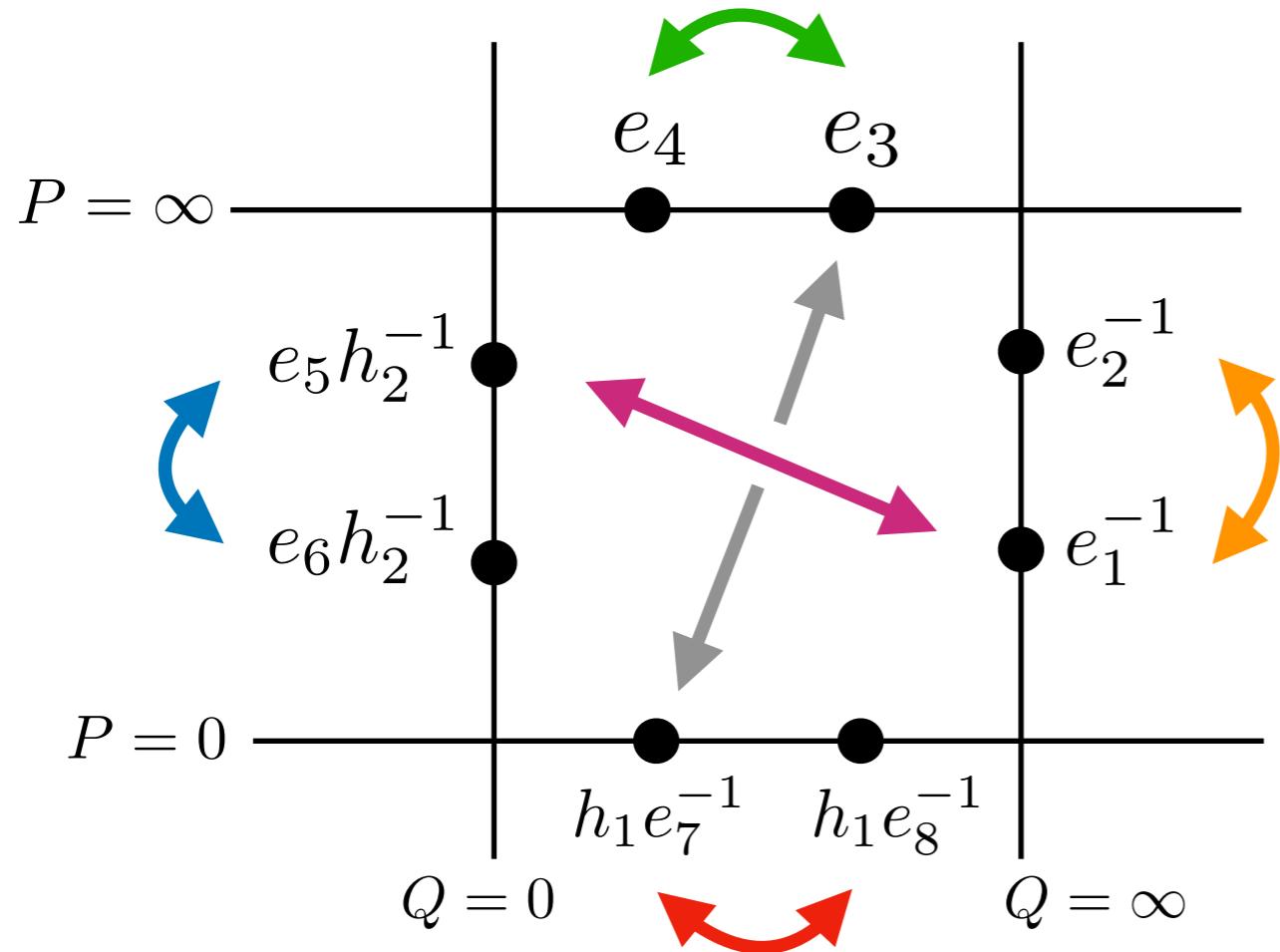
- ① ● expressed by Weyl characters
- ② [ ● coef. of characters = Integer
- has same reps. as those in B-period

# D<sub>5</sub> Del Pezzo geometry

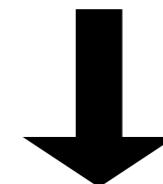
Kubo-Moriyama-Nosaka(2018)

$$0 = \begin{aligned} & e_3 e_4 Q^{-1} P & - (e_3 + e_4) P & + Q P \\ & - h_2^{-1} e_3 e_4 (e_5 + e_6) Q^{-1} & - \frac{z}{\alpha} & - (e_1^{-1} + e_2^{-1}) Q \\ & + h_1^2 (e_1 e_2 e_7 e_8)^{-1} Q^{-1} P^{-1} & - h_1 (e_1 e_2)^{-1} (e_7^{-1} + e_8^{-1}) P^{-1} & + (e_1 e_2)^{-1} Q P^{-1} \end{aligned}$$

- $h_1^2 h_2^2 = \prod_{i=1}^8 e_i$  classical curve



invariant under exchanging of  
asymptotic values



D5 Weyl sym.

# D<sub>5</sub> Del Pezzo geometry Kubo-Moriyama-Nosaka(2018)

$$\begin{aligned}
 \frac{\hat{H}}{\alpha} + \frac{z}{\alpha} = & e_3 e_4 \hat{Q}^{-1} \hat{P} & - (e_3 + e_4) \hat{P} & + \hat{Q} \hat{P} \\
 & - h_2^{-1} e_3 e_4 (e_5 + e_6) \hat{Q}^{-1} & + \frac{E}{\alpha} & - (e_1^{-1} + e_2^{-1}) \hat{Q} \\
 & + h_1^2 (e_1 e_2 e_7 e_8)^{-1} \hat{Q}^{-1} \hat{P}^{-1} & - h_1 (e_1 e_2)^{-1} (e_7^{-1} + e_8^{-1}) \hat{P}^{-1} & + (e_1 e_2)^{-1} \hat{Q} \hat{P}^{-1}
 \end{aligned}$$

quantum curve

10 parameters - (2 + 2 + 1) parameters

- 8 asymptotic values determine curve 2
- $(\hat{P}, \hat{Q}) \sim (A\hat{P}, B\hat{Q})$  2
- $h_1^2 h_2^2 = \prod_{i=1}^8 e_i$  1

→ 5 parameters  $(e_1, e_3, e_5, \bar{h}_1, \bar{h}_2)$

$$(\bar{h}_1 = qh_1, \bar{h}_2 = q^{-1}h_2)$$

# Quantum mirror map

Aganagic, Cheng, Dijkgraaf Krefl, Vafa(2011)

$$\Pi_A(z, \hbar) \sim \oint \frac{\log P[X]}{X} dX = Ez^{-1} + \left( -\frac{E^2}{2} - A_2 \right) z^{-2} + \dots$$

$$\left[ P[X] = \frac{\Psi[q^{-1}X]}{\Psi[X]}, \quad \hat{Q}\Psi[X] = X\Psi[X] \quad \left[ \frac{\hat{H}}{\alpha} + \frac{z}{\alpha} \right] \Psi[X] = 0 \right]$$

solve Schrödinger eq. order by order

$$\frac{A_2}{\alpha^2 q^{-1}} = \frac{e_3}{e_1} + \frac{\bar{h}_1}{e_1} + \frac{\bar{h}_1 e_3}{e_1} + \frac{e_3 e_5}{\bar{h}_2^2} + \frac{e_3 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3^2 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3}{\bar{h}_2} + \frac{e_3}{\bar{h}_2 e_1} + \frac{e_3 e_5}{\bar{h}_2} + \frac{e_3 e_5}{\bar{h}_2 e_1}$$

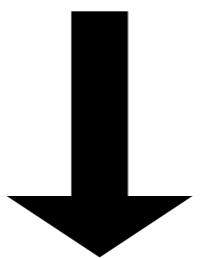
- This is invariant under Weyl transf.
- It contains 10 terms.

# Quantum mirror map

$$\Pi_A(z, \hbar) \sim \oint \frac{\log P[X]}{X} dX = Ez^{-1} + \left( -\frac{E^2}{2} - A_2 \right) z^{-2} + \dots$$

$$\left[ \begin{array}{l} P[X] = \frac{\Psi[q^{-1}X]}{\Psi[X]}, \quad \widehat{Q}\Psi[X] = X\Psi[X] \\ \left[ \frac{\widehat{H}}{\alpha} + \frac{z}{\alpha} \right] \Psi[X] = 0 \\ \text{solve Schrödinger eq. order by order} \end{array} \right]$$

$$\frac{A_2}{\alpha^2 q^{-1}} = \frac{e_3}{e_1} + \frac{\bar{h}_1}{e_1} + \frac{\bar{h}_1 e_3}{e_1} + \frac{e_3 e_5}{\bar{h}_2^2} + \frac{e_3 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3^2 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3}{\bar{h}_2} + \frac{e_3}{\bar{h}_2 e_1} + \frac{e_3 e_5}{\bar{h}_2} + \frac{e_3 e_5}{\bar{h}_2 e_1}$$



$$\alpha = q^{1/2} \bar{h}_2^{1/2} e_1^{1/4} e_3^{-1/2} e_5^{-1/4}$$

$$A_2 = \chi_{10}$$

# Plan

A-period is similar to B-period



- expressed by Weyl characters
- coef. of characters = Integer
- has same reps. as those in B-period

②

# Multi-covering & BPS

$$\Pi_A(z, \hbar) \sim \sum_{l=1}^{\infty} (-1)^{l+1} A_l z^{-l} \quad A_4 \ni \frac{3}{2}\chi_{54} + \frac{5}{2}\chi_{45} + \frac{11}{2}\chi_1$$

fractional

→ Does this have multi-covering structure ?

- A-period of  $A_1$  geometry has **multi-covering structure.**  
(ABJM theory)
- By this structure, coefficients are **integers.**

[Hatsuda-Marino-Moriyama-Okuyama(2013)]

multi-covering structure

$$(\text{coef. of } n\text{-th order}) = \sum_{j \leq n} (\text{coef. of } j\text{-th order})$$

# Multi-covering & BPS

$$\Pi_A(z, \hbar) =: \log z_{\text{eff}} - \log z$$

$$\rightarrow \log z \sim -A_1 z_{\text{eff}}^{-1} + (A_2 - A_1^2) z_{\text{eff}}^{-2} - \left( A_3 - 3A_1 A_2 + \frac{3}{2} A_1^3 \right) z_{\text{eff}}^{-3}$$
$$+ \left( A_4 - 2A_2^2 - 4A_1 A_3 - 8A_1^2 A_2 - \frac{8}{3} A_1^4 \right) z_{\text{eff}}^{-4} + \dots$$

consider inverse function

$$\begin{cases} A_1 = 0 \\ A_2 = \chi_{\mathbf{10}} \\ A_3 = (q^{1/2} + q^{-1/2}) \chi_{\mathbf{16}} \\ A_4 = (q^2 + q^{-2}) \chi_{\mathbf{1}} + (q^{3/2} + q^{-3/2}) (\chi_{\mathbf{45}} + 3\chi_{\mathbf{1}}) + \frac{3\chi_{\mathbf{54}} + 5\chi_{\mathbf{45}} + 11\chi_{\mathbf{1}}}{2} \end{cases}$$

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still fractional

# Multi-covering & BPS

$$\Pi_A(z, \hbar) =: \log z_{\text{eff}} - \log z$$

$$\begin{aligned}
 \rightarrow \log z \sim & \underbrace{-A_1 z_{\text{eff}}^{-1} + (A_2 - A_1^2) z_{\text{eff}}^{-2}}_{+ \underbrace{\left( A_4 - 2A_2^2 - 4A_1 A_3 - 8A_1^2 A_2 - \frac{8}{3} A_1^4 \right) z_{\text{eff}}^{-4}}_{= \epsilon_1(q, e_i, \bar{h}_i)} + \dots}_{\text{multi-covering}} \\
 & + \underbrace{\left( A_3 - 3A_1 A_2 + \frac{3}{2} A_1^3 \right) z_{\text{eff}}^{-3}}_{= -\epsilon_2(q, e_i, \bar{h}_i) + \frac{\epsilon_1(q^2, e_i^2, \bar{h}_i^2)}{2}} \\
 & + \underbrace{\epsilon_3(q, e_i, \bar{h}_i) + \frac{\epsilon_1(q^3, e_i^3, \bar{h}_i^3)}{3}}_{= -\epsilon_4(q, e_i, \bar{h}_i) + \frac{\epsilon_2(q^2, e_i^2, \bar{h}_i^2)}{2} + \frac{\epsilon_1(q^4, e_i^4, \bar{h}_i^4)}{4}} \\
 & \qquad \qquad \qquad \text{structure}
 \end{aligned}$$

$$\epsilon_1 = 0, \quad -\epsilon_2 = \chi_{\mathbf{10}},$$

$$\epsilon_3 = \left( q^{1/2} + q^{-1/2} \right) \chi_{\mathbf{16}}, \quad -\epsilon_4 = \left( q^2 + q^{-2} \right) \chi_{\mathbf{1}} + \left( q + q^{-1} \right) (\chi_{\mathbf{45}} + 3\chi_{\mathbf{1}}) + 4\chi_{\mathbf{1}}$$

integer !

# Multi-covering & BPS

[Moriyama-Nosaka-Yano(2017)]

$d$	$(j_L, j_R)$	BPS	$(-1)^{d-1} \sum_{ d =1} (N_{j_L, j_R}^d)_{d+ -d-}$	representations
1	$(0; 0)$	16	$8_{+1} + 8_{-1}$	<b>16</b>
2	$(0, \frac{1}{2})$	10	$1_{+2} + 8_0 + 1_{-2}$	<b>10</b>
3	$(0, 1)$	16	$8_{+1} + 8_{-1}$	<b>16</b>
4	$(0, \frac{1}{2})$	1	$1_0$	<b>1</b>
	$(0, \frac{3}{2})$	45	$8_{+2} + 29_0 + 8_{-2}$	<b>45</b>
	$(\frac{1}{2}, 2)$	1	$1_0$	<b>1</b>

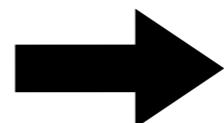
A-period

$$\left[ \begin{array}{l} \epsilon_1 = 0, \\ -\epsilon_2 = \underline{\chi_{10}}, \\ \epsilon_3 = (q^{\frac{1}{2}} + q^{-\frac{1}{2}})\underline{\chi_{16}}, \\ -\epsilon_4 = (q^2 + q^{-2})\underline{\chi_1} + (q + q^{-1})(\underline{\chi_{45}} + 3\underline{\chi_1}) + 4\underline{\chi_1}, \end{array} \right]$$

B-period

same reps. as B-period for each order

Physical meaning ??



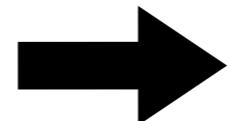
Future work

# Summary

A-period is similar to B-period

- expressed by Weyl characters
- coef. of characters = Integer
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Physical meaning ??



Future work