Light Dark Matter Scattering in Gravitational Wave Detectors

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DM Scattering in GW Detectors

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Interferometer Used by GW Experiments



Figure: The interferometer with the arm length (distance between front mirror and end mirror) equals to L. [Taken from KAGRA PhD Thesis: D.Chen, 2016]

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The change in the interference of the light is proportional to

$$\Delta L = \Delta L_x - \Delta L_y$$

A gravitational wave will induce a strain, h_{GW}

$$h_{
m GW}\sim rac{\Delta L}{L}\leq 10^{-20}$$

- What about strain induced by DM ?
- Is it visible compared to the noises? What noises?

Strain Amplitude Budgets



Figure: Example of the strain amplitude of the noises. [https://gwcenter.icrr.u-tokyo.ac.jp/en/researcher/parameter]

Interferometer Isolation



Figure: Schematic of the interferometer isolation. [Taken from KAGRA PhD Thesis: D.Chen, 2016]

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Toy Model to Estimate DM Induced Strain

Simple pendulum describing the TM and its suspension system

$$m\ddot{x}_{c} + k_{c}(1 + \mathrm{i}\phi)x_{c} = \frac{F_{\mathrm{ext},c}}{L}$$

We model the total displacement [Moore, Cole, Berry '14]

$$x_{\text{tot},c}(t) = x_{\text{th},c}(t) + x_{\text{qu},c}(t) + x_{\text{DM},c}(t)$$

 The strain can be obtained from one-sided power spectral density (PSD)

$$h_i(\omega) = \sqrt{S_i(\omega)}$$

 According to fluctuation-dissipation theorem, the thermal noise PSD is given by [Callen, Welton '51]

$$S_{\text{th}}(\omega) = rac{4k_BT}{L^2\omega^2} \Re[Y(\omega)]$$

• The admittance for $F_{\text{ext}} \sim exp(i\omega t)$ in this case

$$Y(\omega) \equiv \frac{\dot{x}_{\text{th},c}}{F_{\text{ext}}/(mL)} = i\omega D_c^{-1}(\omega) = \frac{i\omega}{(-m\omega^2 + k_c(1 + i\phi))}$$

• $D_c(\omega)$ is the Fourier transform of the differential operator.

• The Thermal noise PSD of the Toy model ($\omega_c^2 \equiv k_c/m$)

$$S_{\text{th},c}(\omega) = \frac{4k_BT}{L^2} \frac{\phi \, \omega_c^2 / (m \, \omega)}{(\omega^2 - \omega_c^2)^2 + \omega_c^4 \phi^2}$$

The associate Thermal noise strain

$$h_{\rm th}(\omega) = \sqrt{S_{\rm th}(\omega)}$$

For quantum noise, we use the standard quantum limit (SQL)

$$S_{qu}(\omega) = rac{8\hbar}{m\omega^2 L^2}$$

The corresponding strain

$$h_{\mathsf{qu}}(\omega) = \sqrt{\mathcal{S}_{\mathsf{qu}}(\omega)}$$

The total noise reads

$$h_{\sf n}(\omega) = \sqrt{h_{\sf th}^2(\omega) + h_{\sf qu}^2(\omega)}$$

DM Signal

Recall the simple pendulum motion

$$m\ddot{x}_{\text{DM}} + k_c(1 + i\phi) x_{\text{DM}} = \frac{F_{\text{ext}}}{L}$$

Using Fourier expansion

$$m{g}(t) = \int_{-\infty}^{\infty} \, m{d}\omega \, ilde{m{g}}(\omega) m{e}^{{
m i}\omega t}$$

- We assume single DM hit at t = 0 [Lee, Nugroho, MS '20; Tsuchida et al.'19] $F_{\rm ext}(t) = q_R \, \delta(t)$
- The displacement induced by DM in frequency domain reads

$$\widetilde{x}_{\mathsf{DM}}(\omega) = rac{q_R}{mL} rac{1}{(-\omega^2 + \omega_c^2(1+\mathrm{i}\phi))}$$

The DM induced strain

$$h_{\mathsf{DM}}(\omega) = \sqrt{rac{2\,\omega}{\pi}} | ilde{x}_{\mathsf{DM}}(\omega)|$$

The Strain for the Toy Model



Figure: Thermal noise, SQL, DM Signal.

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The optimal SNR [Moore, Cole, Berry '14]

$$arrho^{2} = \int_{f_{\min}}^{f_{\max}} df \, rac{4 \, | ilde{x}_{\mathsf{DM}}(2\pi f)|^{2}}{S_{n}(2\pi f)}$$

- ► For: $f_0 = 175.4 \text{ Hz}$, $\phi = 6.32 \times 10^{-12}$, m = 22.8 kg, T = 19 K, L = 3 km, $m_{\text{DM}} = 1 \text{ GeV}/c^2$, and $|\vec{v}_{DM}| = 220 \text{ km/s}$
- Around the peak at full width half maximum (FWHM)

$$arrho_{ ext{th}}^2 = rac{1}{2\pi} rac{q_R^2}{m \, k_B \, T} = rac{1}{2\pi} rac{E_R}{E_{ ext{th}}} = 4.09 imes 10^{-24}$$

- $E_R = 3.37 \times 10^{-45} \,\text{J}$ and $E_{\text{th}} = 1.31 \times 10^{-23} \,\text{J}$
- Lighter mirror and colder for better SNR.

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- KAGRA is GW detector currently running in Japan
- Equipped with cryogenic system with mirror temperature around 20 K
- The suspension system is modelled as triple pendulum consist of: IM, Blade Springs, and TM

Figure: [KAGRA PhD Thesis: D.Chen, 2016]

KAGRA Suspension Thermal Noise

Vertical suspension thermal noise [KAGRA Document, JGW-T1707038v9]

$$\left(Mrac{d^2}{dt^2}+K_{
m v}
ight)ec{x}_{
m v}(t)=rac{ec{F}_{{
m ext},
m v}(t)}{L}$$

$$M = \begin{pmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{pmatrix}, \quad K_v = \begin{pmatrix} K_{1v} + K_{2v} & -K_{2v} & 0 \\ -K_{2v} & K_{2v} + K_{3v} & -K_{3v} \\ 0 & -K_{3v} & K_{3v} \end{pmatrix}$$
$$\vec{x}_v(t) = \begin{pmatrix} x_{1v}(t) \\ x_{2v}(t) \\ x_{3v}(t) \end{pmatrix} \text{ and } \vec{F}_{\text{ext},v}(t) = \begin{pmatrix} F_{\text{ext},1v}(t) \\ F_{\text{ext},2v}(t) \\ F_{\text{ext},3v}(t) \end{pmatrix}.$$

$$K_{iv} \equiv k_{iv} (1 + \mathrm{i} \, \phi_{iv})$$

The index i = (1, 2, 3) stands for (IM, BS, and TM)

KAGRA Suspension Thermal Noise

The equation in frequency domain

$$D_{\nu}(\omega)\vec{\tilde{x}}_{\nu} \equiv (-\omega^2 M + K_{\nu})\vec{\tilde{x}}_{\nu} = \frac{\vec{\tilde{F}}_{\mathrm{th},\nu}}{L}$$

The vertical admittance reads

$$Y_{v}(\omega) \equiv i\omega D_{v}^{-1}(\omega)$$

Thermal noise PSD of the vertical thermal noise

$$S_{\mathrm{th},\nu}(\omega) = rac{4 \, k_B \, T \Re(Y_{
u}(\omega))_{33}}{L^2 \omega^2}$$

The strain amplitude of vertical suspension thermal noise

$$h_{\mathsf{th}, v}(\omega) = \mathsf{VHC}\sqrt{4|\mathcal{S}_{\mathsf{th}, v}(\omega)|}$$

VHC due to the tilt of the baseline. Its value is ¹/₂₀₀

- Horizontal thermal noise proceed in similar manner as the verctical one
- Mirror thermal noise and quantum noise
- Sum over the total noise

$$h_{\text{tot}}(\omega) = \sqrt{h_{\text{th},v}^2 + h_{\text{th},h}^2 + h_{\text{qu}}^2(\omega) + h_{\text{mir}}^2(\omega)}$$



Figure: total noise: vert.susp., horz.susp., mirror noise, quantum noise (solid blue), and SQL (dashed blue) [Lee, Nugroho, MS '20; JGW-T1707038v9]

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DM hit the j-th component from vertical direction [Lee, Nugroho, MS '20]

$$F_{\mathsf{DM},jv}(t) = q_{\mathsf{R},jv}\delta(t)$$

The Fourier transform of the displacement of the i-th component

$$ilde{x}_{i
u}(\omega) = \sum_{j=1}^{3} (D_{v}^{-1}(\omega))_{ij} rac{ ilde{F}_{j,v}(\omega)}{L}$$

Since we probe the third component, we have

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$$|\tilde{x}_{\mathsf{DM},\nu}(\omega)|^2 \equiv |\tilde{x}_{3\nu}(\omega)|^2 = \left|\sum_{j=1}^3 (D_{\nu}^{-1}(\omega))_{3j} \frac{\tilde{F}_{j,\nu}(\omega)}{L}\right|^2$$

The strain is given by

$$h_{\mathsf{DM}, v}(\omega) = \mathsf{VHC}\sqrt{rac{2\,\omega}{\pi}} | ilde{x}_{\mathsf{DM}, v}(\omega)|^2$$

The DM induced strain in horizontal direction

$$h_{\mathsf{DM}, \mathbf{v}}(\omega) = \sqrt{rac{2\omega}{\pi} | ilde{x}_{\mathsf{DM}, h}(\omega)|^2}$$

The total DM signal

$$|\tilde{x}_{\mathsf{DM}}(\omega)|^2 = \mathsf{VHC}^2 |\tilde{x}_{\mathsf{DM}, v}(\omega)|^2 + |\tilde{x}_{\mathsf{DM}, h}(\omega)|^2$$



Figure: total noise, DM vertical hit, and DM horizontal hit

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For horizontal DM hit

$$\varrho_h^2 = \int_{f_{\min,h}}^{f_{\max,h}} df \frac{4|\tilde{x}_{\mathsf{DM}}(2\pi f)|^2}{S_{\mathsf{tot}}(2\pi f)}$$

We integrate around the peak which is located at 175.4 Hz

$$f_{\min,h} = 170$$
Hz and $f_{\max,h} = 180$ Hz

This gives the SNR

$$\varrho_h^2 = 6.94 \times 10^{-18}$$

For vertical hit around 31.4 Hz

$$f_{\min,h} = 30.4$$
Hz and $f_{\max,h} = 32.4$ Hz

The corresponding SNR

$$\varrho_v^2 = 1.89 \times 10^{-21}$$

► LISA path finder sensitivity within 0.1 mHz $\leq f \leq$ 30 Hz [LISA Pathfinder '16]

$$\sqrt{S_{\Delta g}} \leq 3\sqrt{2} \,\,\mathrm{fm}\,\mathrm{s}^{-2}/\sqrt{\mathrm{Hz}} imes \sqrt{1+(f/8\,\mathrm{mHz})^4}$$

DM induced strain

$$\sqrt{S_{\Delta g, \text{DM}}} \sim 4.1 imes 10^{-7} \sqrt{rac{f}{\text{Hz}}} \, \text{fm} \, \text{s}^{-2} / \sqrt{\text{Hz}}$$

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- The current interferometer used in GW experiments are very sensitive.
- DM can excite mechanical resonance as suspension thermal noise does.
- ► For DM detection, lighter and colder mirror are needed.
- ► LPF is quite sensitive, but not sensitive enough to detect DM.