Chapter 8

Black Hole Gap Emission from Supermassive BHs

- §1 VHE observations of AGNs
- §2 The case of billion solar mass BHs
- §3 The case of million solar mass BHs
- §4 Conclusions

September 7, 2012 NCTS, NTHU HIROTANI, Kouichi

§1 VHE observations of Active Galactic Nuclei (AGNs)

§ 1 VHE observations of AGNs

So far, IACTs have detected 52 AGNs in VHE. Most of them are blazers, while five are non-blazers.



Fig.) Examples of Imaging **Atmospheric Cherenkov telescopes (IACTs).**

Upper: **MAGIC** (Major Atmospheric Gamma Imaging Cherenkov).

Lower: **CTA** (Cherenkov Telescope Array; *image*) So far, IACTs have detected 52 AGNs in VHE. Most of them are blazers, while five are non-blazers.

Blazer jets are strongly beamed towards us, masking the faint central region near the BH.

Non-blazer jets are beamed away from our L.O.S. Thus, the central regions are more easily investigated.

Among them, IC310 and M87 exhibit horizon-scale rapid variability in VHE (Abramowski+ 2012, ApJ, 746, 151; Aleksić+ 2014, Sci 346, 1080, 2014).

In this talk, I thus focus on the horizon-scale VHE emission from such low luminosity AGNs.

§ 1 VHE observations of AGNs

For SMBHs, HE & VHE emissions are expected in the shock-in-jet model.

Marscher & Gear 1985, ApJ 298, 114

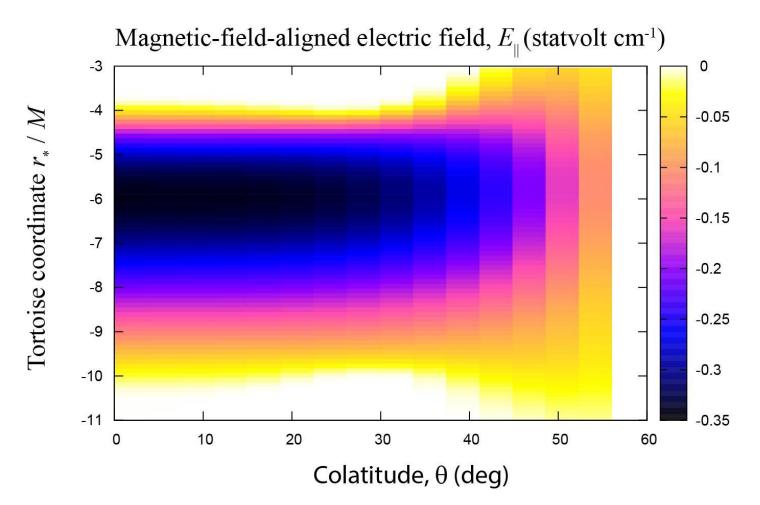
HE/VHE flux increases w/ increasing \dot{M} .

However, VHE emissions are also predicted to be emitted from the BH-gaps of SMBHs. KH & Pu 2016, ApJ 818, 50; KH + 2016, ApJ 833, 142

VHE flux increases w/ decreasing \dot{M} .

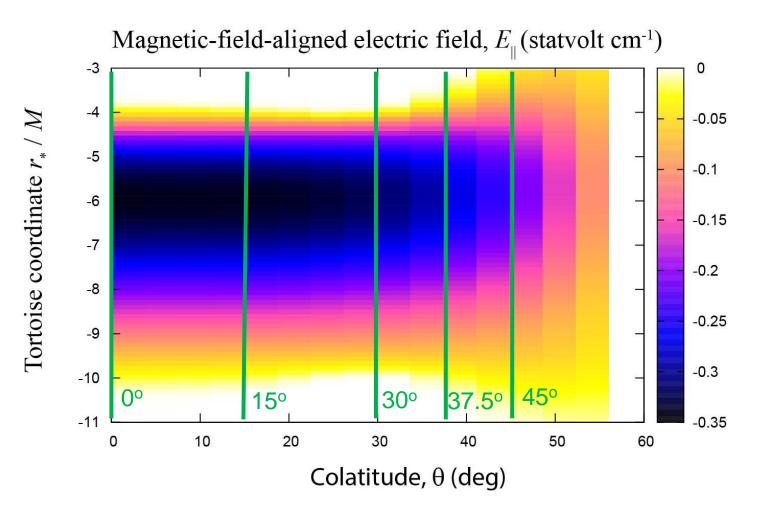
We will focus on the BH-gap model and discuss its theoretical predictions.

 E_{\parallel} distribution on (r, θ) plane. Polar funnel is bounded with ADAF at $\theta = 60^{\circ}$. $M = 10^{9} M_{\odot}$ and a = 0.90M. $\dot{m} = 1.8 \times 10^{-5}$.

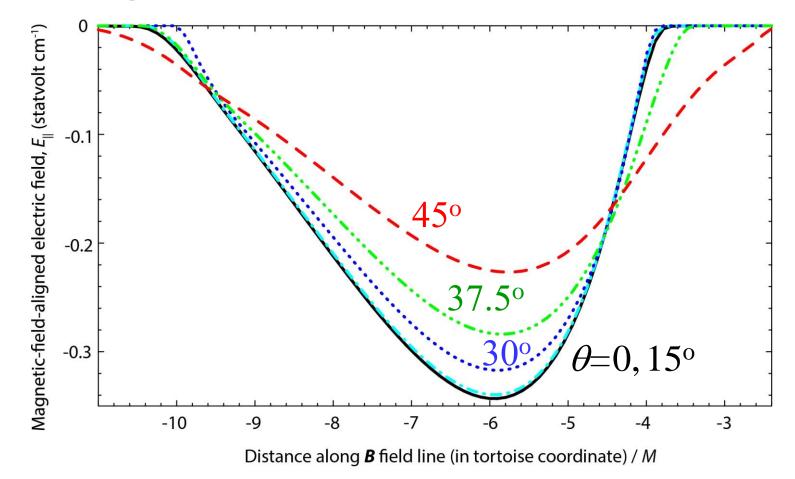


 E_{\parallel} distribution on (r, θ) plane.

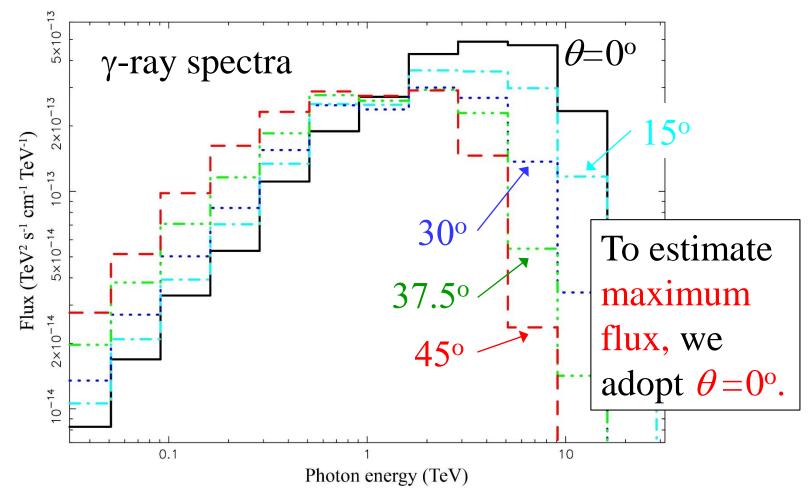
Slice at five discrete colatitude, θ .



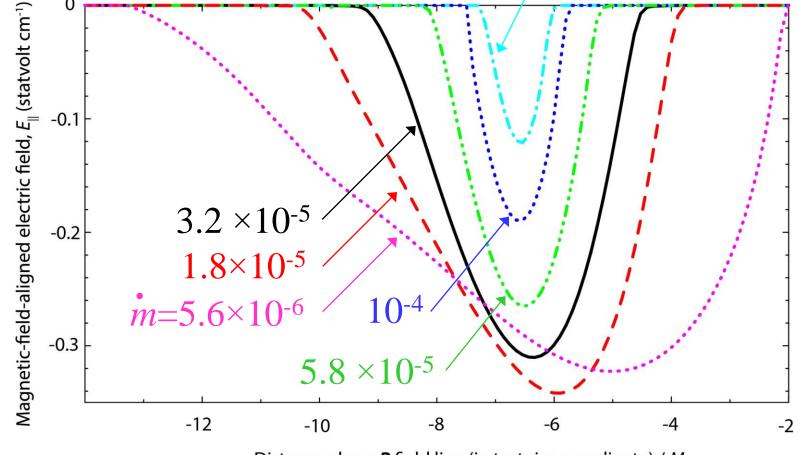
 E_{\parallel} at five discrete θ 's. Within $\theta < 15-30^{\circ}$, maximum potential drop is realized. $M=10^{9}M_{\odot}$ and a=0.90M. $\dot{m}=1.8 \times 10^{-5}$.



SED (30GeV-30TeV) vs. viewing angle, θ . Gap emission is hardest along rotation axis, $\theta = 0^{\circ}$. $M = 10^{9} M_{\odot}$ and a = 0.90M. $m = 1.8 \times 10^{-5}$.



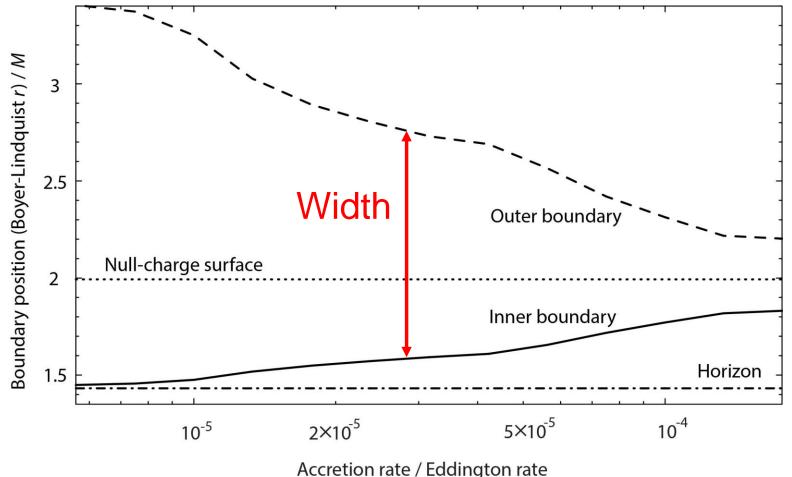
 E_{\parallel} at five six discrete accretion rates, \dot{m} 's. Gap enlarges as \dot{m} approaches the lower bound, $\sim 3 \times 10^{-5}$. $M = 10^9 M_{\odot}$ and a = 0.90M. $\theta = 0^{\circ}$. 1.8×10^{-4}



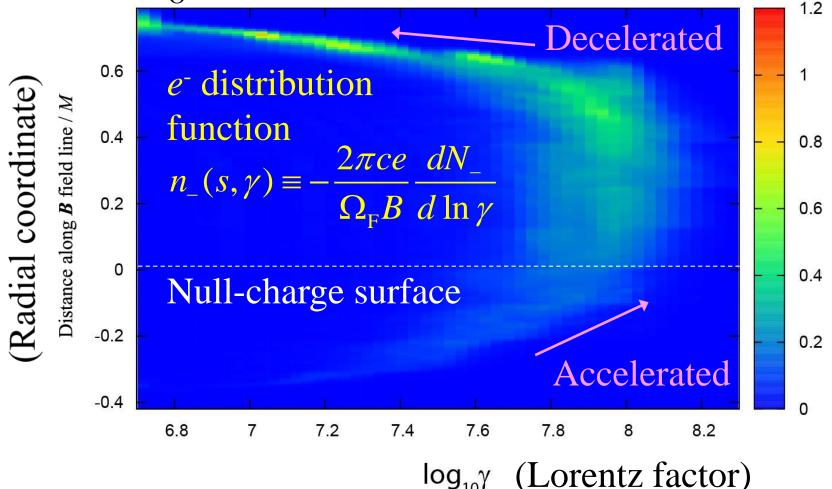
Distance along **B** field line (in tortoise coordinate) / M

Gap extent vs. accretion rate, \dot{m} .

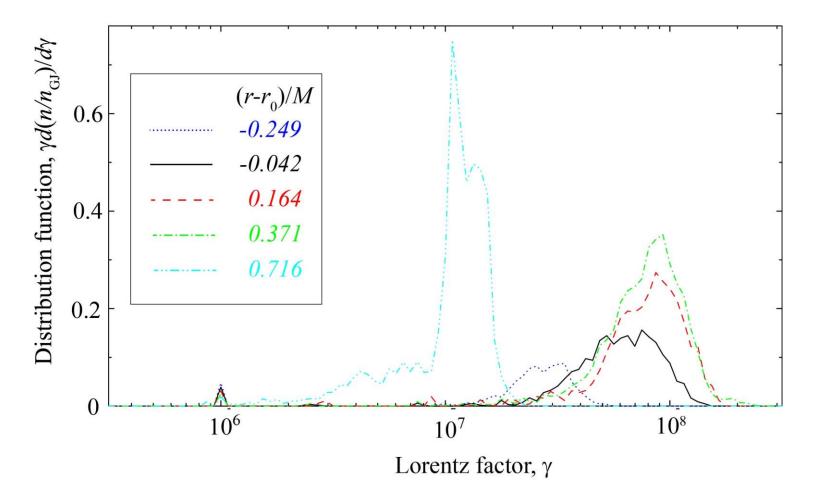
Gap enlarges as *m* approaches the lower bound, $\sim 3 \times 10^{-6}$. $M = 10^9 M_{\odot}$ and a = 0.90M. $\theta = 0^{\circ}$.



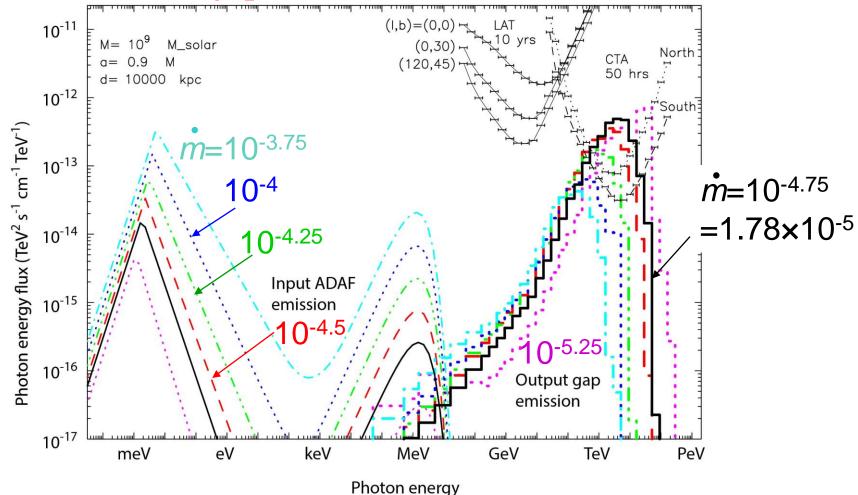
 $E_{\parallel} < 0 \rightarrow e^{-1}$'s (e^{+1} 's) are accelerated outwards (inwards). e^{-1} 's saturate between $4 \times 10^7 < \gamma < 1.5 \times 10^8$ via ICS. $M = 10^9 M_{\odot}$ and a = 0.90M. $\theta = 0^{\circ}$.



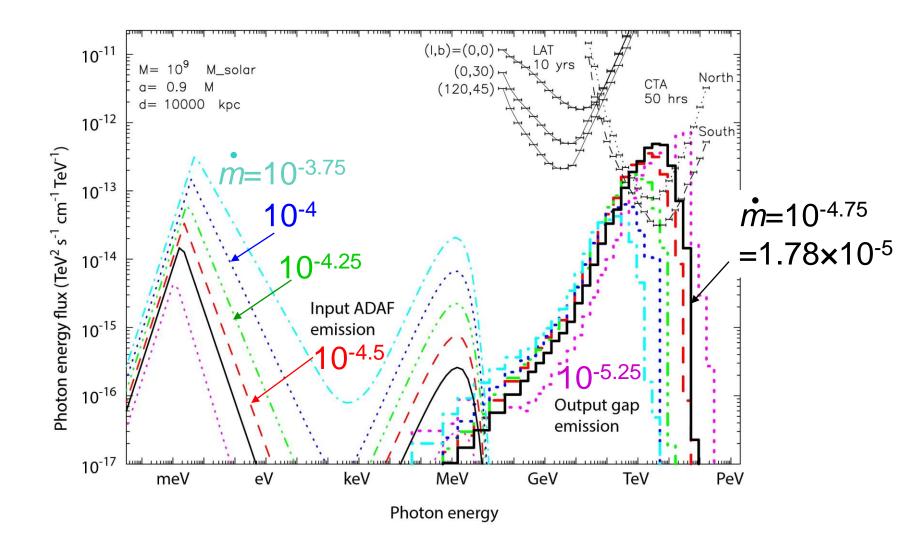
 e^{-} distribution function @ five discrete radial coordinates. e^{-} 's saturate between $4 \times 10^{7} < \gamma < 1.5 \times 10^{8}$ via ICS. $M = 10^{9} M_{\odot}$ and a = 0.90M. $\theta = 0^{\circ}$.



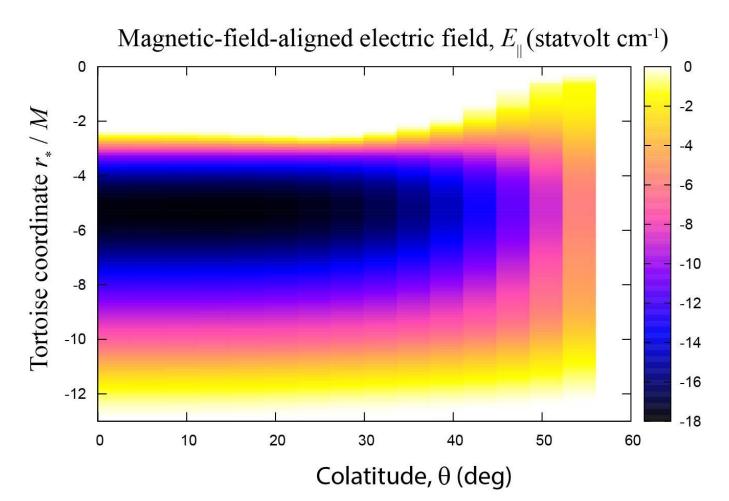
SEDs @ six discrete accretion rates, \dot{m} . If d < 10Mpc, a hyper-massive BH with $M = 10^9 M_{\odot}$ can emit detectable gap emission in 1-30TeV.



Gap luminosity increases with decreasing \dot{m} . \rightarrow Anti-correlation between mm & VHE fluxes.

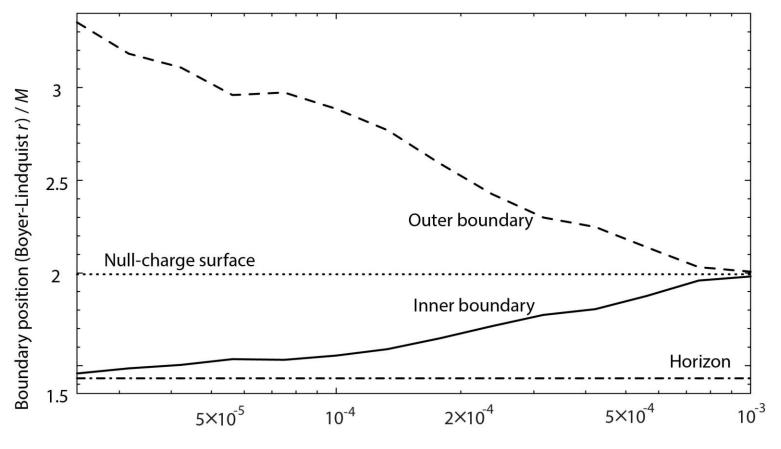


 E_{\parallel} distribution on (r, θ) plane. Polar funnel is bounded with ADAF at $\theta = 60^{\circ}$. $M = 10^{6}M_{\odot}$ and a = 0.90M. $\dot{m} = 1.8 \times 10^{-5}$.



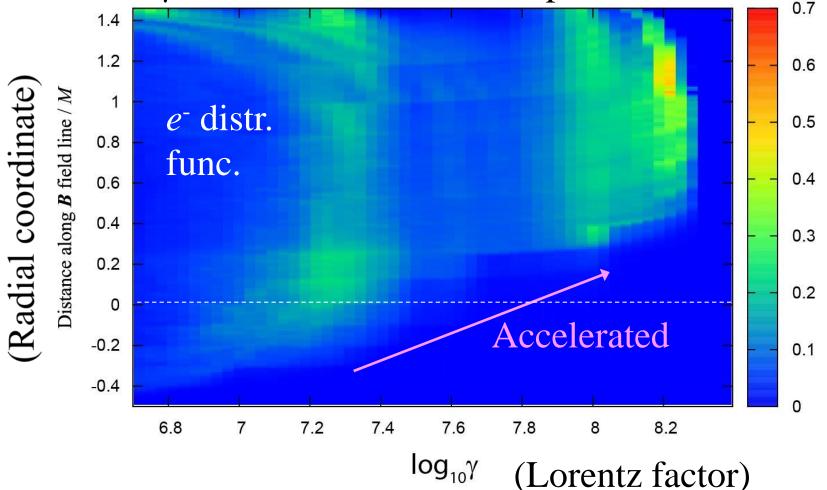
Gap extent vs. accretion rate, \dot{m} .

Gap enlarges as \dot{m} approaches the lower bound, ~1.8×10⁻⁵. $M=10^{6}M_{\odot}$ and a=0.90M. $\theta=0^{\circ}$.

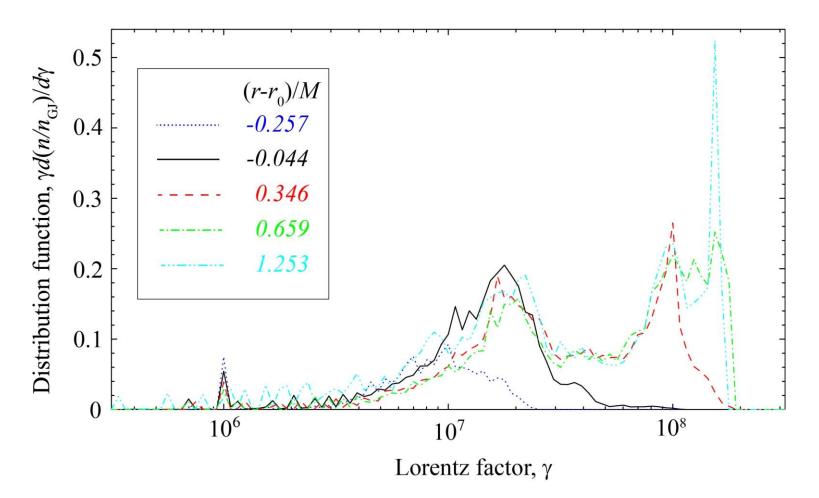


Accretion rate / Eddington rate

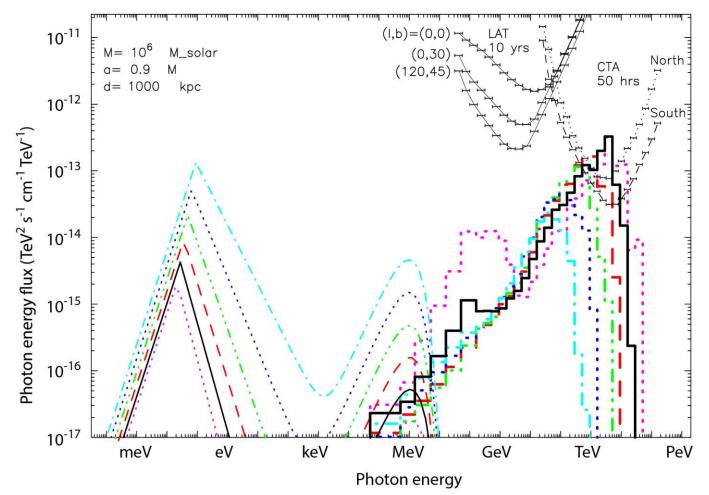
 $E_{\parallel} < 0 \rightarrow e^{-1}$'s (e^{+1} 's) are accelerated outwards (inwards). e^{-1} 's saturate between $10^7 < \gamma < 2.5 \times 10^7$ via ICS, and $9 \times 10^7 < \gamma < 1.8 \times 10^8$ via curvature process.

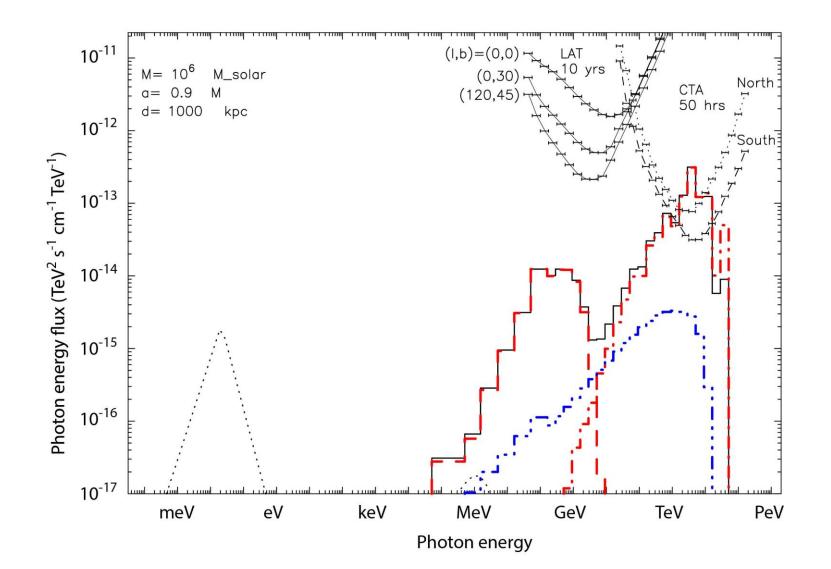


 e^{-} distribution function @ five discrete radial coordinates. Energy distribution is broadened because both IC & curvature processes contribute for $M=10^{6}M_{\odot}$ case.

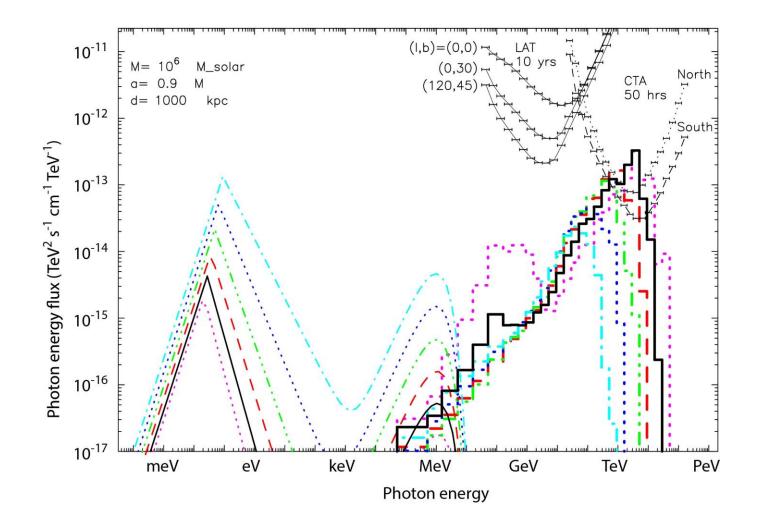


SEDs @ six discrete accretion rates, \dot{m} . If d < 10Mpc, a hyper-massive BH with $M = 10^9 M_{\odot}$ can emit detectable gap emission in 1-30TeV.





Gap luminosity increases with decreasing \dot{m} . \rightarrow Anti-correlation betw. sub-mm & VHE fluxes.



Summary on SMBH gap

When the mass accretion rate is highly sub-Eddington, the polar funnel becomes charge-starved, because the ADAF-emitted MeV photons materialize less efficiently.

Owing to this charge starvation, particle accelerators (=gaps) appear within a few GM/c^2 above the event horizon.

For low-luminosity AGNs (e.g., $M \sim 10^9 M_{\odot}$), the BH gap emits IC photons in 0.5~30TeV along $\theta \sim 0^{\circ}$. Such a flare is detectable with CTA if d < 30 Mpc when the accretion rate resides in $6 \times 10^{-6} < \dot{M} / \dot{M}_{Edd} < 2 \times 10^{-5}$. We can discriminate gap vs. jet emissions by anticorrelation vs. correlations at submillimeter (or near-IR) & VHE for a SMBH (or a stellar-mass BH).

END OF THE LECTURES