Is the decay of the Higgs boson to a photon and a dark photon currently observable at the LHC?

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General context I

The dark photon A' is an hypothetical Abelian gauge boson that can mix with the photon and has been the subject of extensive theoretical and experimental studies.

A potential discovery channel that has received considerable attention is the decay

$$h \to AA',$$

where A' is assumed to be (effectively) massless and invisible.

Signature: Photon + MET

General context II

An extensive phenomenological literature on the subject

- Gabrielli '14 [Phys. Rev. D 90 no. 5, (2014) 055032]
- Biswas '15 [JHEP 06 (2015) 102]
- Biswas '16 [Phys. Rev. D 93 no. 9, (2016) 093011]
- Biswas '17 [Phys. Rev. D 96 no. 5, (2017) 055012]
- ...

which predicted an upper limit $\mathrm{BR}(h \to AA') < 5\%$ that was claimed to be model-independent.

Motivations:

- Predicted by certain BSM models (flavour)
- Potentially observable
- Interesting experimental channel (subjective)

General context III

This motivated four experimental searches

- CMS '19 [JHEP 10 (2019) 139]
  Z-associated channel
- CMS '20 [JHEP 03 (2021) 011]
  VBF channel
- ATLAS '21 [Eur. Phys. J. C 82 no. 2, (2022) 105]
  VBF channel
- ATLAS '22 [ATLAS-CONF-2022-064]
  Z-associated channel

The strongest limit obtained was $\mathrm{BR}(h \to AA') < 1.8\%$ at 95% CL by ATLAS '21.

Overview

Goal:
Investigate experimental and theoretical constraints on $h \to AA'$

Improvements:

- Additional constraints
- More rigorous treatment
- More recent data

Constraints:

- Higgs signal strengths
- Oblique parameters
- Electric dipole moment (EDM) of the electron
- Unitarity

General form of the amplitude and what it tells us

Gauge invariance imposes

$$M^{h\to AA'} = \left[S^{h\to AA'}(p_1\cdot p_2 g_{\mu\nu} - p_{1\mu}p_{2\nu}) + i\tilde{S}^{h\to AA'}\epsilon_{\mu\nu\alpha\beta}p_1^{\alpha}p_2^{\beta}\right]\epsilon^{\nu}_{p_1}\epsilon^{\mu}_{p_2}.$$

This amplitude cannot be generated at tree-level with a renormalizable Lagrangian.

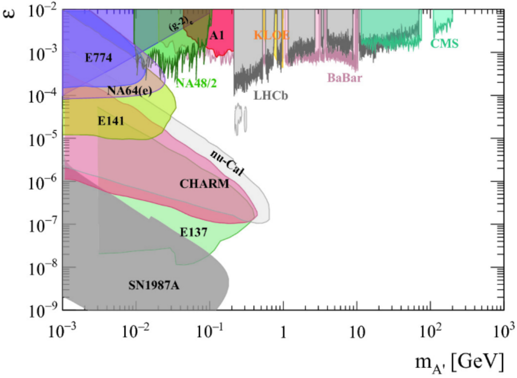
The process must take place at loop-level.

SM particles in the loop?

- There could be kinetic mixing between A and A'

$$\epsilon F^{\mu\nu} F'_{\mu\nu}$$

- SM particles in the loop could then contribute to $h \to AA'$.
- However, ϵ is constrained to be very small for a light A'.
- BR($h \to AA'$) would be negligible.



Fabbrichesi '20

Mediators I

Solution:

New mediators that:

- Interact with the Higgs
- Are charged under EM
- Are charged under a new $U(1)'$

Comments:

- It's not possible to be completely model independent.
- We will consider a very large set of mediators and explain why the bounds would be difficult to avoid.

Mediators II

We will consider mediators that:

1. Have a renormalizable Lagrangian that preserves all gauge symmetries
2. Lead to the $h \to AA'$ decay at one loop
3. Contain no mediators charged under QCD
4. Contain only mediators that are complex scalars or vector-like fermions
5. Contain no more than two new fields
6. Contain no mediators that mix with SM fields or have a non-zero expectation value

Mediators III

All mediators that satisfy these criterias fall into a finite number of categories:

$$\begin{aligned}
&\text{Fermion:} && \bar{\psi}_1(A_L P_L + A_R P_R)\psi_2 H + \text{h.c.}, \\
&\text{Scalar:} && \text{I:} \quad \mu\phi_1^\dagger\phi_2 H + \text{h.c.}, && \text{II:} \quad \lambda H^\dagger H\phi^\dagger\phi, \\
& && \text{III:} \quad \lambda H^\dagger H\phi_1^\dagger\phi_2 + \text{h.c.}, && \text{IV:} \quad \lambda HH\phi_1^\dagger\phi_2 + \text{h.c.}
\end{aligned}$$

Example: Fermion mediators

Consider the Lagrangian

$$\mathcal{L}_m = -\left[\sum_{a,b,c} \hat{d}^{pn}_{abc} \bar{\psi}^a_1 (A_L P_L + A_R P_R) \psi^b_2 H^c + \text{h.c.}\right] - m_1 \bar{\psi}_1 \psi_1 - m_2 \bar{\psi}_2 \psi_2 .$$

where a, b and c are $SU(2)_L$ indices and are summed from 1 to the size of the corresponding multiplet and with

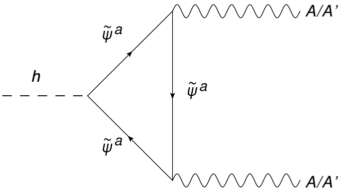
$$\hat{d}^{pn}_{abc} = C^{JM}_{j_1 m_1 j_2 m_2} = \langle j_1 j_2 m_1 m_2 | JM \rangle ,$$

where

$$J = \frac{p-1}{2}, \qquad j_1 = \frac{n-1}{2}, \qquad j_2 = \frac{1}{2},$$
$$M = \frac{p+1-2a}{2}, \quad m_1 = \frac{n+1-2b}{2}, \quad m_2 = \frac{3-2c}{2}.$$

There is a phase that cannot be reabsorbed.

Higgs decay I



+ arrows reversed

The amplitudes are

$$M^{h\to AA} = \left[S^{h\to AA}(p_1 \cdot p_2 g_{\mu\nu} - p_{1\mu}p_{2\nu}) + i\tilde{S}^{h\to AA}\epsilon_{\mu\nu\alpha\beta}p_1^\alpha p_2^\beta\right]\epsilon_{p_1}^\nu \epsilon_{p_2}^\mu,$$
$$M^{h\to AA'} = \left[S^{h\to AA'}(p_1 \cdot p_2 g_{\mu\nu} - p_{1\mu}p_{2\nu}) + i\tilde{S}^{h\to AA'}\epsilon_{\mu\nu\alpha\beta}p_1^\alpha p_2^\beta\right]\epsilon_{p_1}^\nu \epsilon_{p_2}^\mu,$$
$$M^{h\to A'A'} = \left[S^{h\to A'A'}(p_1 \cdot p_2 g_{\mu\nu} - p_{1\mu}p_{2\nu}) + i\tilde{S}^{h\to A'A'}\epsilon_{\mu\nu\alpha\beta}p_1^\alpha p_2^\beta\right]\epsilon_{p_1}^\nu \epsilon_{p_2}^\mu,$$

Higgs decay II

With

$$S^{h\to AA} = e^2 \sum_a \mathrm{Re}(\Omega_{aa}) \tilde{Q}_{aa}^2 S_a + S_{\mathrm{SM}}^{h\to AA}, \qquad \tilde{S}^{h\to AA} = e^2 \sum_a \mathrm{Im}(\Omega_{aa}) \tilde{Q}_{aa}^2 \tilde{S}_a + \tilde{S}_{\mathrm{SM}}^{h\to AA},$$

$$S^{h\to AA'} = ee' \sum_a \mathrm{Re}(\Omega_{aa}) \tilde{Q}_{aa} Q' S_a, \qquad \tilde{S}^{h\to AA'} = ee' \sum_a \mathrm{Im}(\Omega_{aa}) \tilde{Q}_{aa} Q' \tilde{S}_a,$$

$$S^{h\to A'A'} = e'^2 \sum_a \mathrm{Re}(\Omega_{aa}) Q'^2 S_a, \qquad \tilde{S}^{h\to A'A'} = e'^2 \sum_a \mathrm{Im}(\Omega_{aa}) Q'^2 \tilde{S}_a,$$

where in the example of the fermion case

$$S_a = -\frac{m_a}{2\pi^2 m_h^2}\left(2 + (4m_a^2 - m_H^2) C_0(0, 0, m_H^2; m_a, m_a, m_a)\right),$$

$$\tilde{S}_a = -i\frac{m_a}{2\pi^2} C_0(0, 0, m_H^2; m_a, m_a, m_a),$$

and

$$\Gamma^{h\to AA} = \frac{|S^{h\to AA}|^2 + |\tilde{S}^{h\to AA}|^2}{64\pi} m_h^3, \qquad \Gamma^{h\to AA'} = \frac{|S^{h\to AA'}|^2 + |\tilde{S}^{h\to AA'}|^2}{32\pi} m_h^3,$$

$$\Gamma^{h\to A'A'} = \frac{|S^{h\to A'A'}|^2 + |\tilde{S}^{h\to A'A'}|^2}{64\pi} m_h^3.$$

Higgs decay III

Comments:

- The presence of the Levi-Civita symbols comes from γ^5 in the Higgs/fermion vertex.
- The amplitudes are correlated. The constraints on $h \to AA$ and $h \to A'A'$ will constrain $h \to AA'$.
- For $h \to AA$, the interference with the SM contributions will typically dominate. This can be avoided in two ways:
  - Purely imaginary contribution to $S^{h \to AA}$
    Not possible because:
    - Coefficient is real.
    - Diagram cannot be cut with charged particles below $m_h/2$.
  - Contibution only to $\tilde{S}^{h \to AA}$
    - Will be constrained by the electron EDM

Higgs signal strengths

κ formalism

$$\kappa_i^2 = \frac{\sigma_i}{\sigma_i^{\text{SM}}} \quad \text{or} \quad \kappa_i^2 = \frac{\Gamma_i}{\Gamma_i^{\text{SM}}},$$

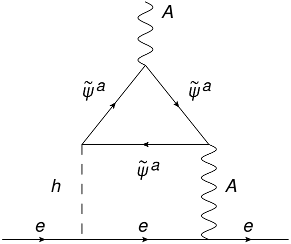
and

$$\kappa_{AA}^2 = \frac{|S^{h\to AA}|^2 + |\tilde{S}^{h\to AA}|^2}{|S_{\text{SM}}^{h\to AA}|^2 + |\tilde{S}_{\text{SM}}^{h\to AA}|^2}, \qquad \kappa_{ZA}^2 = \frac{|S^{h\to ZA}|^2 + |\tilde{S}^{h\to ZA}|^2}{|S_{\text{SM}}^{h\to ZA}|^2 + |\tilde{S}_{\text{SM}}^{h\to ZA}|^2}.$$

Do χ^2 fit using results from

- CMS-PAS-HIG-19-005,
- ATLAS-CONF-2021-053.

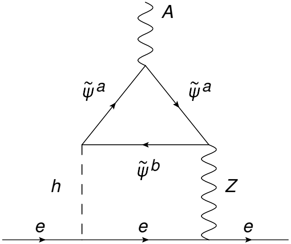
EDM

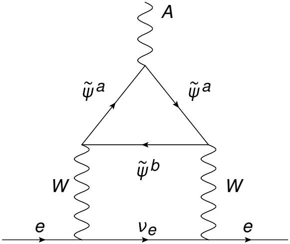


hA diagrams are dominant near limits

hA contribution $\propto \mathrm{Im}(\Omega_{aa})$

Effectively forces Ω_{aa} to be purely real or tiny





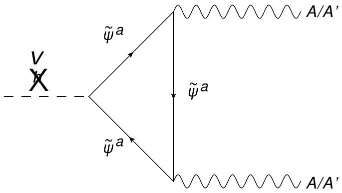
Oblique parameters

The Peskin-Takeuchi parameters are defined as

$$\alpha S = 4s_W^2 c_W^2 \left[\Pi'_{ZZ}(0) - \frac{c_W^2 - s_W^2}{s_W c_W}\Pi'_{ZA}(0) - \Pi'_{AA}(0)\right]$$

$$\alpha T = \frac{\Pi_{WW}(0)}{m_W^2} - \frac{\Pi_{ZZ}(0)}{m_Z^2}.$$

Unsurprising contribution



Unitarity

An amplitude can be expanded as

$$\mathcal{M} = 16\pi \sum_l (2l+1) a_l P_l(\cos\theta), \qquad \max\left(\left|\mathrm{Re}\left(a_0^{\mathrm{eig}}\right)\right|\right) < \frac{1}{2}.$$

Different cases:

$$\text{Fermion:} \quad |A_R|^2 + |A_L|^2 < \frac{32\pi}{p}$$

$$\text{Scalar II:} \quad \frac{1}{16\sqrt{2}\pi}\left[\sum_{i,j}\left|\sum_r \lambda^r \hat{d}_{22ij}^{nr}\right|^2\right]^{\frac{1}{2}} < \frac{1}{2},$$

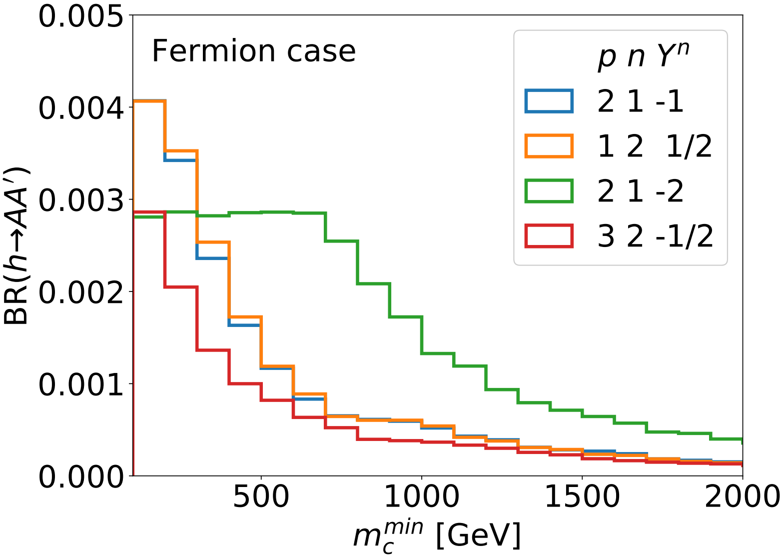
$$\text{Scalar III:} \quad \frac{1}{16\sqrt{2}\pi}\left[\sum_{i,j}\left|\sum_r \lambda^r \hat{d}_{22ij}^{pnr}\right|^2\right]^{\frac{1}{2}} < \frac{1}{2},$$

$$\text{Scalar IV:} \quad \frac{|\lambda|}{16\sqrt{2}\pi}\left[\sum_{i,j}\left|\hat{d}_{22ij}^{pn}\right|^2\right]^{\frac{1}{2}} < \frac{1}{2}.$$

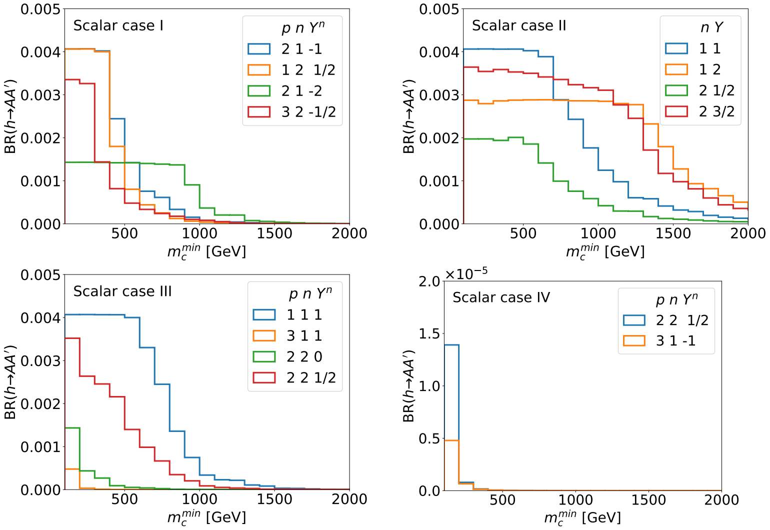
Assembling everything

- Consider multiple benchmark models
- Scan entire parameter space with Markov chain using Metropolis-Hasting algorithm
- Prior proportional to $\mathrm{BR}(h \to AA')^2$ to speed up simulations
- Impose $|Q'e'| < \sqrt{4\pi}$ for fermions or $|Q'e'| < \frac{\sqrt{4\pi}}{q^{1/4}}$ for scalars.

Constraints: Fermion mediators



Constraints: Scalar mediators



Caveats

Three caveats:

- There should be a lower limit on the mass of the charged mediators, but there does not exist sufficiently general searches. The limits would probably be stronger.
- It could have been possible that $S_{\text{BSM}}^{h\to AA} \approx -2S_{\text{SM}}^{h\to AA}$. That would avoid the Higgs signal strength constraints, but it does not happen for the fermion cases considered. It can technically happen for some scalar cases, but only for rare cases and extreme fine-tuning.
- The bounds are model dependent.

Conclusion

Goal:
Investigate experimental and theoretical constraints on $h \to AA'$

Conclusions:

- $\mathrm{BR}(h \to AA') \lesssim 0.4\%$ at best*
- Difficult to even get this large
- Far more constrained than previous phenomenology papers (5%) and experimental searches (1.8%)
- Seems unlikely that we will be able to probe this channel at the LHC

Thanks!

Higgs decay IV

- If the loops are dominated by one particle, $S^{h\to AA}$ is purely real and $\tilde{S}^{h\to AA} = 0$, we have

$$\mathrm{BR}(h \to AA') \approx \sqrt{\mathrm{BR}(h \to A'A')\mathrm{BR}(h \to AA)} \left| \frac{\Delta \mathrm{BR}(h \to AA)}{\mathrm{BR}(h \to AA)} \right| ,$$

where $\mathrm{BR}(h \to AA) = \mathrm{BR}(h \to AA)_{\mathrm{SM}} + \Delta\mathrm{BR}(h \to AA)$.

Considering
$\mathrm{BR}(h \to AA) \sim 0.23\%$, $\left| \frac{\Delta \mathrm{BR}(h\to AA)}{\mathrm{BR}(h\to AA)} \right| \lesssim 25\%$,
$\mathrm{BR}(h \to A'A') \lesssim 10\%$.

we get a rough estimate of $\mathrm{BR}(h \to AA') \lesssim 0.4\%$.

EDM I

$$\frac{d_e^{Ah}}{e} = -\sum_a \frac{\alpha \tilde{Q}_{aa}^2 m_a m_e}{16\pi^3 m_h^2 v} \mathrm{Im}(\Omega_{aa}) \int_0^1 dx \frac{1}{x(1-x)} j\left(0, \frac{m_a^2}{x(1-x)m_h^2}\right),$$

where

$$j(r,s) = \frac{1}{r-s}\left(\frac{r\ln r}{r-1} - \frac{s\ln s}{s-1}\right).$$

EDM II

$$\frac{d_e^{Zh}}{e} = \sum_{a,b} \frac{\tilde{Q}_{bb}}{32\pi^4 m_h^2} g_{ee}^V g_{ee}^S \left(m_a C_{ab}^1 f_1(m_a, m_b) + m_b C_{ab}^2 f_2(m_a, m_b) \right),$$

where

$$C_{ab}^1 = \mathrm{Re}\left(i g_{ba}^S g_{ab}^{A*} - g_{ba}^P g_{ab}^{V*} \right), \qquad C_{ab}^2 = -\mathrm{Re}\left(i g_{ba}^S g_{ab}^{A*} + g_{ba}^P g_{ab}^{V*} \right),$$

with

$$g_{ee}^S = -\frac{m_e}{v}, \qquad g_{ee}^V = -\frac{\sqrt{g^2 + g'^2}}{2} \left(-\frac{1}{2} + 2 s_W^2 \right),$$

$$g_{ab}^S = -\frac{(\Omega_{ba} + \Omega_{ab}^*)}{2}, \qquad g_{ab}^P = -i \frac{(\Omega_{ba} - \Omega_{ab}^*)}{2},$$

$$g_{ab}^V = -\frac{\sqrt{g^2 + g'^2}}{2} (B_{Rba} + B_{Lba}), \qquad g_{ab}^A = -\frac{\sqrt{g^2 + g'^2}}{2} (B_{Rba} - B_{Lba}),$$

and

$$f_1(m_a, m_b) = \int_0^1 dx j\left(\frac{m_Z^2}{m_h^2}, \frac{\tilde{\Delta}_{ab}}{m_h^2} \right), \qquad f_2(m_a, m_b) = \int_0^1 dx j\left(\frac{m_Z^2}{m_h^2}, \frac{\tilde{\Delta}_{ab}}{m_h^2} \right) \frac{(1-x)}{x},$$

EDM III

$$\frac{d_e^{WW}}{e} = -\frac{\alpha^2 m_e}{8\pi^2 s_W^4 m_W^2} \sum_{a,b} \frac{m_a m_b}{m_W^2} \mathrm{Im}\left(\hat{A}_{Lba}\hat{A}^*_{Rba}\right)\left[\tilde{Q}_{bb}\mathcal{G}(r_a, r_b, 0) + \tilde{Q}_{aa}\mathcal{G}(r_b, r_a, 0)\right],$$

where $r_a = m_a^2/m_W^2$, $r_b = m_b^2/m_W^2$ and

$$\mathcal{G}(r_a, r_b, r_c) = \int_0^1 \frac{d\gamma}{\gamma}\int_0^1 dyy\left[\frac{(R-3K_{ab})R + 2(K_{ab}+R)y}{4R(K_{ab}-R)^2} + \frac{K_{ab}(K_{ab}-2y)}{2(K_{ab}-R)^3}\ln\frac{K_{ab}}{R}\right],$$

where

$$R = y + (1-y)r_c \qquad K_{ab}, = \frac{r_a}{1-\gamma} + \frac{r_b}{\gamma}.$$

Oblique parameters fermions

$$S = \frac{1}{2\pi} \sum_{a,b} \Big\{ \left(|\hat{A}_{Lab}|^2 + |\hat{A}_{Rab}|^2 \right) \psi_+(y_a, y_b) + 2\mathrm{Re} \left(\hat{A}_{Lab} \hat{A}^*_{Rab} \right) \psi_-(y_a, y_b)$$

$$- \frac{1}{2} \left[\left(|X_{ab}|^2 + |X_{Rab}|^2 \right) \chi_+(y_a, y_b) + 2\mathrm{Re} \left(X_{Lab} X^*_{Rab} \right) \chi_-(y_a, y_b) \right] \Big\},$$

$$T = \frac{1}{16\pi s_W^2 c_W^2} \sum_{a,b} \Big\{ \left(|\hat{A}_{Lab}|^2 + |\hat{A}_{Rab}|^2 \right) \theta_+(y_a, y_b) + 2\mathrm{Re} \left(\hat{A}_{Lab} \hat{A}^*_{Rab} \right) \theta_-(y_a, y_b)$$

$$- \frac{1}{2} \left[\left(|X_{Lab}|^2 + |X_{Rab}|^2 \right) \theta_+(y_a, y_b) + 2\mathrm{Re} \left(X_{Lab} X^*_{Rab} \right) \theta_-(y_a, y_b) \right] \Big\},$$

with $y_a = m_a^2/m_Z^2$, $X_{L/R} = -2B_{L/R} + 2\tilde{Q} s_W^2$ and

$$\psi_+(y_1, y_2) = \frac{1}{3} - \frac{1}{9} \ln \frac{y_1}{y_2},$$

$$\psi_-(y_1, y_2) = -\frac{y_1 + y_2}{6\sqrt{y_1 y_2}},$$

$$\chi_+(y_1, y_2) = \frac{5(y_1^2 + y_2^2) - 22 y_1 y_2}{9(y_1 - y_2)^2} + \frac{3 y_1 y_2 (y_1 + y_2) - y_1^3 - y_2^3}{3(y_1 - y_2)^3} \ln \frac{y_1}{y_2},$$

$$\chi_-(y_1, y_2) = -\sqrt{y_1 y_2} \left[\frac{y_1 + y_2}{6 y_1 y_2} - \frac{y_1 + y_2}{(y_1 - y_2)^2} + \frac{2 y_1 y_2}{(y_1 - y_2)^3} \ln \frac{y_1}{y_2} \right],$$

$$\theta_+(y_1, y_2) = y_1 + y_2 - \frac{2 y_1 y_2}{y_1 - y_2} \ln \frac{y_1}{y_2},$$

$$\theta_-(y_1, y_2) = 2\sqrt{y_1 y_2} \left[\frac{y_1 + y_2}{y_1 - y_2} \ln \frac{y_1}{y_2} - 2 \right].$$

Oblique parameters scalars

$$S = \frac{1}{2\pi} \sum_{a,b} \left[|B_{ab}|^2 - (c_W^2 - s_W^2) B_{ab} \tilde{Q}_{ab} - c_W^2 s_W^2 \tilde{Q}_{ab}^2 \right] F_1(y_a, y_b),$$

$$T = \frac{1}{16\pi c_W^2 s_W^2} \left[\sum_{a,b} |\hat{A}_{ab}|^2 F_2(y_a, y_b) - \sum_a \left[\hat{A}\hat{A}^\dagger + \hat{A}^\dagger \hat{A} \right]_{aa} F_3(y_a) \right.$$
$$\left. -2 \sum_{a,b} |B_{ab}|^2 F_2(y_a, y_b) + 4 \sum_a \left[B^2 \right]_{aa} F_3(y_a) \right],$$

where $y_a = m_a^2 / m_Z^2$ and

$$F_1(y_1, y_2) = -\frac{5y_1^2 - 22 y_1 y_2 + 5 y_2^2}{9(y_1 - y_2)^2} + \frac{2\left(y_1^2 (y_1 - 3y_2) \ln y_1 - y_2^2 (y_2 - 3y_1) \ln y_2 \right)}{3(y_1 - y_2)^3},$$

$$F_2(y_1, y_2) = 3\,(y_1 + y_2) - \frac{2\left(y_1^2 \ln y_1 - y_2^2 \ln y_2 \right)}{y_1 - y_2},$$

$$F_3(y_1) = 2y_1 - 2y_1 \ln y_1 .$$

Example: Fermion mediators II

Once the Higgs is replaced by its expectation value

$$\mathcal{L}_m \supset -\sum_{a,b} \frac{A_L v}{\sqrt{2}} \hat{d}^{pn}_{ab2} \bar{\psi}^a_1 P_L \psi^b_2 - \sum_{a,b} \frac{A_R^* v}{\sqrt{2}} \hat{d}^{pn}_{ba2} \bar{\psi}^a_2 P_L \psi^b_1 - m_1 \bar{\psi}_1 P_L \psi_1 - m_2 \bar{\psi}_2 P_L \psi_2 + \text{h.c.}$$

Introduce the convenient notation

$$\hat{\psi} = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

and

$$d^{pn}_{ab} = \begin{cases} \hat{d}^{pn}_{a(b-p)2}, & \text{if } a \in [1, p] \text{ and } b \in [p+1, n+p], \\ 0, & \text{otherwise.} \end{cases}$$

Fermion mediators: Higgs interactions and mass IV

The mass Lagrangian can be written as

$$\mathcal{L}_m \supset -\sum_{a,b} M_{ab}\bar{\hat{\psi}}^a P_L \hat{\psi}^b + \text{h.c.},$$

where the mass matrix is

$$M = m_1 \begin{pmatrix} \mathbb{I}_{p\times p} & 0_{p\times n} \\ 0_{n\times p} & 0_{n\times n} \end{pmatrix} + m_2 \begin{pmatrix} 0_{p\times p} & 0_{p\times n} \\ 0_{n\times p} & \mathbb{I}_{n\times n} \end{pmatrix} + \frac{A_L v}{\sqrt{2}} d^{pn} + \frac{A_R^* v}{\sqrt{2}} d^{pnT}.$$

The mass matrix can then be diagonalized by introducing $\tilde{\psi}^a$

$$P_L\hat{\psi} = R_L P_L \tilde{\psi}, \quad P_R\hat{\psi} = R_R P_R \tilde{\psi},$$

where R_L and R_R are unitary matrices that diagonalize $M^\dagger M$ and $MM^\dagger$ respectively.

Fermion mediators: Higgs interactions and mass V

The interactions of the Higgs with the mass eigenstates are

$$\mathcal{L}_m \supset -\sum_{a,b} \Omega_{ab} h \bar{\tilde{\psi}}^a P_L \tilde{\psi}^b + \text{h.c.},$$

where Ω is

$$\Omega = \frac{A_L}{\sqrt{2}} R_R^\dagger d^{pn} R_L + \frac{A_R^*}{\sqrt{2}} R_R^\dagger d^{pnT} R_L.$$

Fermion mediators: Gauge interactions I

The interactions of the A/A' with $\tilde{\psi}^a$ are

$$\mathcal{L}_g \supset -eA_\mu \bar{\tilde{\psi}}\gamma^\mu \tilde{Q}\tilde{\psi} - Q'e'A'_\mu \bar{\tilde{\psi}}\gamma^\mu \tilde{\psi},$$

where e' is the $U(1)'$ gauge coupling and

$$\tilde{Q} = R_L^\dagger \hat{Q} R_L = R_R^\dagger \hat{Q} R_R,$$

where

$$\hat{Q} = \begin{pmatrix} Y^p + T_3^p & 0_{p\times n} \\ 0_{n\times p} & Y^n + T_3^n \end{pmatrix},$$

where $(T_3^p)_{ab} = (p+1-2a)\delta_{ab}/2$ and similarly for T_3^n.

Fermion mediators: Gauge interactions II

The interactions between the Z boson and $\tilde{\psi}^a$ are

$$\mathcal{L}_g \supset -\sqrt{g^2 + g'^2} Z_\mu \bar{\tilde{\psi}} \gamma^\mu (B_L P_L + B_R P_R) \tilde{\psi},$$

where B_L and B_R are

$$B_L = R_L^\dagger \begin{pmatrix} -s_W^2 Y^p + c_W^2 T_3^p & 0_{p\times n} \\ 0_{n\times p} & -s_W^2 Y^n + c_W^2 T_3^n \end{pmatrix} R_L,$$
$$B_R = R_R^\dagger \begin{pmatrix} -s_W^2 Y^p + c_W^2 T_3^p & 0_{p\times n} \\ 0_{n\times p} & -s_W^2 Y^n + c_W^2 T_3^n \end{pmatrix} R_R.$$

Fermion mediators: Gauge interactions III

The interactions between the W boson and $\tilde{\psi}^a$ are

$$\mathcal{L}_g \supset -\frac{g}{\sqrt{2}}\bar{\tilde{\psi}}\gamma^\mu\left(\hat{A}_L P_L W_\mu^+ + \hat{A}_R P_R W_\mu^+\right)\tilde{\psi} + \text{h.c.},$$

where

$$\hat{A}_L = R_L^\dagger \begin{pmatrix} T_+^p & 0_{p\times n} \\ 0_{n\times p} & T_+^n \end{pmatrix} R_L, \quad \hat{A}_R = R_R^\dagger \begin{pmatrix} T_+^p & 0_{p\times n} \\ 0_{n\times p} & T_+^n \end{pmatrix} R_R,$$

with $(T_+^p)_{ab} = \sqrt{a(p-a)}\delta_{a,b-1}$ and similarly for T_+^n.

Fermion mediators: Unitarity I

The top loop has a mass of ~ 173 GeV, $y_t \sim 1$ and $N_c = 3$, yet gives a tiny contribution to $h \to AA$.

Mediators must have either large Yukawa couplings or dark electric charge.

An amplitude can be expanded as

$$\mathcal{M} = 16\pi \sum_l (2l+1) a_l P_l(\cos\theta).$$

Consider the basis of $\bar{\psi}_1^a \psi_2^b$ pairs given by

$$\bar{\psi}_1^1\psi_2^1,\ \bar{\psi}_1^1\psi_2^2,\ \ldots,\ \bar{\psi}_1^1\psi_2^n,\ \bar{\psi}_1^2\psi_2^1,\ \bar{\psi}_1^2\psi_2^2,\ \ldots,\ \bar{\psi}_1^2\psi_2^n,\ \ldots,\ \bar{\psi}_1^p\psi_2^1,\ \bar{\psi}_1^p\psi_2^2,\ \ldots,\ \bar{\psi}_1^p\psi_2^n.$$

Fermion mediators: Unitarity II

Then, the matrix of a_0 for the scattering $\bar{\psi}_1^a \psi_2^b \to \bar{\psi}_1^c \psi_2^d$ is

$$a_0^{\mathrm{mat}} = \begin{pmatrix}
F_{11}^{11} & F_{11}^{12} & \dots & F_{11}^{1n} & F_{11}^{21} & F_{11}^{22} & \dots & F_{11}^{2n} & \dots & F_{11}^{p1} & F_{11}^{p2} & \dots & F_{11}^{pn} \\
F_{12}^{11} & F_{12}^{12} & \dots & F_{12}^{1n} & F_{12}^{21} & F_{12}^{22} & \dots & F_{12}^{2n} & \dots & F_{12}^{p1} & F_{12}^{p2} & \dots & F_{12}^{pn} \\
\dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
F_{1n}^{11} & F_{1n}^{12} & \dots & F_{1n}^{1n} & F_{1n}^{21} & F_{1n}^{22} & \dots & F_{1n}^{2n} & \dots & F_{1n}^{p1} & F_{1n}^{p2} & \dots & F_{1n}^{pn} \\
F_{21}^{11} & F_{21}^{12} & \dots & F_{21}^{1n} & F_{21}^{21} & F_{21}^{22} & \dots & F_{21}^{2n} & \dots & F_{21}^{p1} & F_{21}^{p2} & \dots & F_{21}^{pn} \\
F_{22}^{11} & F_{22}^{12} & \dots & F_{22}^{1n} & F_{22}^{21} & F_{22}^{22} & \dots & F_{22}^{2n} & \dots & F_{22}^{p1} & F_{22}^{p2} & \dots & F_{22}^{pn} \\
\dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
F_{2n}^{11} & F_{2n}^{12} & \dots & F_{2n}^{1n} & F_{2n}^{21} & F_{2n}^{22} & \dots & F_{2n}^{2n} & \dots & F_{2n}^{p1} & F_{2n}^{p2} & \dots & F_{2n}^{pn} \\
\dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
F_{p1}^{11} & F_{p1}^{12} & \dots & F_{p1}^{1n} & F_{p1}^{21} & F_{p1}^{22} & \dots & F_{p1}^{2n} & \dots & F_{p1}^{p1} & F_{p1}^{p2} & \dots & F_{p1}^{pn} \\
F_{p2}^{11} & F_{p2}^{12} & \dots & F_{p2}^{1n} & F_{p2}^{21} & F_{p2}^{22} & \dots & F_{p2}^{2n} & \dots & F_{p2}^{p1} & F_{p2}^{p2} & \dots & F_{p2}^{pn} \\
\dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
F_{pn}^{11} & F_{pn}^{12} & \dots & F_{pn}^{1n} & F_{pn}^{21} & F_{pn}^{22} & \dots & F_{pn}^{2n} & \dots & F_{pn}^{p1} & F_{pn}^{p2} & \dots & F_{pn}^{pn}
\end{pmatrix}$$

Fermion mediators: Unitarity III

Where F_{ab}^{cd} is

$$F_{ab}^{cd} = \frac{d_{ab2}^{pn} d_{cd2}^{pn}}{32\pi} \begin{pmatrix} -|A_R|^2 & A_R A_L^* \\ A_L A_R^* & -|A_L|^2 \end{pmatrix}$$

and corresponds to a_0 for different combinations of helicity in the basis $(\uparrow\uparrow, \downarrow\downarrow)$. Call a_0^{eig} the set of eigenvalues of a_0^{mat}. Unitarity requires

$$\max\left(\left|\text{Re}\left(a_0^{\text{eig}}\right)\right|\right) < \frac{1}{2}.$$

This simplifies to

$$|A_R|^2 + |A_L|^2 < \frac{32\pi}{p}.$$

Scalar mediators: Scalar case I A

Consider:

- A complex scalar ϕ_1 with the properties:
  - Representation of $SU(2)_L$ of dimension $p = n \pm 1$
  - Weak hypercharge of $Y^p = Y^n + 1/2$
  - Charge of Q' under $U(1)'$
- A complex scalar ϕ_2 with the properties:
  - Representation of $SU(2)_L$ of dimension n
  - Weak hypercharge of Y^n
  - Charge of Q' under $U(1)'$

Scalar mediators: Scalar case I B

This allows the Lagrangian

$$\mathcal{L}_m^1 = -\left[\sum_{a,b,c} \mu \hat{d}_{abc}^{pn} \phi_1^{a\dagger} \phi_2^b H^c + \text{h.c.}\right] - m_1^2|\phi_1|^2 - m_2^2|\phi_2|^2,$$

where

$$\hat{d}_{abc}^{pn} = C_{j_1 m_1 j_2 m_2}^{JM},$$

with

$$J = \frac{p-1}{2}, \qquad j_1 = \frac{n-1}{2}, \qquad j_2 = \frac{1}{2},$$
$$M = \frac{p+1-2a}{2}, \quad m_1 = \frac{n+1-2b}{2}, \quad m_2 = \frac{3-2c}{2}.$$

μ can be made real.

Scalar mediators: Scalar case II A

Consider:

- A complex scalar ϕ with the properties:
  - Representation of $SU(2)_L$ of dimension n
  - Weak hypercharge of Y^n
  - Charge of Q' under $U(1)'$

Scalar mediators: Scalar case II B

This allows the Lagrangian

$$\mathcal{L}_m^2 = -\sum_{r\in\{n-1,n+1\}}\sum_{a,b,c,d}\lambda^r \hat{d}^{nr}_{abcd} H^{a\dagger} H^b \phi^{c\dagger}\phi^d - m^2|\phi|^2.$$

The $SU(2)_L$ tensor $\hat{d}^{nr}_{abcd}$ is given by

$$\hat{d}^{nr}_{abcd} = \sum_M C^{JM}_{j_1 m_1 j_2 m_2} C^{JM}_{j_3 m_3 j_4 m_4},$$

where M is summed over $\{-J, -J+1, -J+2, ..., +J\}$ and

$$j_1 = \frac{1}{2}, \quad j_2 = \frac{n-1}{2}, \quad j_3 = \frac{1}{2}, \quad j_4 = \frac{n-1}{2}, \quad J = \frac{r-1}{2},$$
$$m_1 = \frac{3-2a}{2}, \quad m_2 = \frac{n+1-2c}{2}, \quad m_3 = \frac{3-2b}{2}, \quad m_4 = \frac{n+1-2d}{2}.$$

Two possible coefficients if $n \neq 1$.

Scalar mediators: Scalar case III A

Consider:

- A complex scalar ϕ_1 with the properties:
  - Representation of $SU(2)_L$ of dimension $p \in \{n-2, n, n+2\}$
  - Weak hypercharge of $Y^p = Y^n$
  - Charge of Q' under $U(1)'$
- A complex scalar ϕ_2 with the properties:
  - Representation of $SU(2)_L$ of dimension n
  - Weak hypercharge of Y^n
  - Charge of Q' under $U(1)'$

Scalar mediators: Scalar case III B

This allows the Lagrangian:

$$\mathcal{L}_m^3 = -\left[\sum_{r\in\mathcal{R}}\sum_{a,b,c,d}\lambda^r \hat{d}^{pnr}_{abcd} H^{a\dagger} H^b \phi_1^{c\dagger}\phi_2^d + \text{h.c.}\right] - m_1^2|\phi_1|^2 - m_2^2|\phi_2|^2,$$

where $\mathcal{R} = \{n-1, n+1\}\cap\{p-1, p+1\}$ and $\hat{d}^{pnr}_{abcd}$

$$\hat{d}^{pnr}_{abcd} = \sum_M C^{JM}_{j_1 m_1 j_2 m_2} C^{JM}_{j_3 m_3 j_4 m_4},$$

where M is summed over $\{-J, -J+1, -J+2, ..., +J\}$ and

$$j_1 = \frac{1}{2}, \quad j_2 = \frac{p-1}{2}, \quad j_3 = \frac{1}{2}, \quad j_4 = \frac{n-1}{2}, \quad J = \frac{r-1}{2},$$
$$m_1 = \frac{3-2a}{2}, \quad m_2 = \frac{p+1-2c}{2} \quad m_3 = \frac{3-2b}{2}, \quad m_4 = \frac{n+1-2d}{2}.$$

If $p = n \neq 1$, there are two coefficients and 1 otherwise.

Scalar mediators: Scalar case IV A

Consider:

- A complex scalar ϕ_1 with the properties:
  - Representation of $SU(2)_L$ of dimension $p \in \{n-2, n, n+2\}$
  - Weak hypercharge of $Y^p = Y^n + 1$
  - Charge of Q' under $U(1)'$
- A complex scalar ϕ_2 with the properties:
  - Representation of $SU(2)_L$ of dimension n
  - Weak hypercharge of Y^n
  - Charge of Q' under $U(1)'$
- p and n are not both 1.

Scalar mediators: Scalar case IV B

Consider:

$$\mathcal{L}_m^4 = -\left[\lambda \hat{d}^{pn}_{abcd} H^a H^b \phi_1^{c\dagger} \phi_2^d + \text{h.c.}\right] - m_1^2|\phi_1|^2 - m_2^2|\phi_2|^2,$$

where

$$\hat{d}^{pn}_{abcd} = \sum_{M_1} C^{J_1 M_1}_{j_1 m_1 j_2 m_2} C^{J_2 M_2}_{J_1 M_1 j_3 m_3},$$

where M_1 is summed over $\{-1, 0, 1\}$ and

$$j_1 = \frac{1}{2}, \quad j_2 = \frac{1}{2}, \quad j_3 = \frac{n-1}{2}, \quad J_1 = 1, \quad J_2 = \frac{p-1}{2},$$
$$m_1 = \frac{3-2a}{2}, \quad m_2 = \frac{3-2b}{2} \quad m_3 = \frac{n+1-2d}{2}, \quad M_2 = \frac{p+1-2c}{2}.$$

λ can be made real.

Scalar mediators: Constraints I

- Treatment of gauge and Higgs interactions is similar to the fermion case.
- Higgs signal strengths is similar to the fermion case.
- We computed the oblique parameters ourselves.
- EDM constraints are not needed (no γ^5).

Scalar mediators: Constraints III

For case IV, this can be simplified to:

$$|\lambda| < 8\pi\sqrt{\frac{6}{p}}.$$

Also,

$$|Q'e'| < \frac{\sqrt{4\pi}}{q^{1/4}},$$

where $q = n + p$ for cases I, III, IV and $q = n$ for case II.